

1 **The Physics of Space Weather/Solar Terrestrial Physics (STP): What We Know Now and**
2 **What Are the Current and Future Challenges?**

3

4 Bruce T. Tsurutani¹, Gurbax S. Lakhina², Rajkumar Hajra³

5

6 ¹Pasadena, Calif, USA

7 ²Indian Institute for Geomagnetism, Navi Mumbai, India

8 ³Indian Institute of Technology Indore, Simrol, Indore, India

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10 **ABSTRACT**

11 Major geomagnetic storms are caused by unusually intense solar wind southward magnetic fields
12 that impinge upon the Earth's magnetosphere (Dungey, 1961). How can we predict the occurrence
13 of future interplanetary events? Do we currently know enough of the underlying physics and do
14 we have sufficient observations of solar wind phenomena that will impinge upon the Earth's
15 magnetosphere? We view this as the most important challenge in Space Weather. We discuss the
16 case for magnetic clouds (MCs), interplanetary sheaths upstream of interplanetary coronal mass
17 ejections (ICMEs), corotating interaction regions (CIRs) and solar wind high-speed streams
18 (HSSs). The sheath- and CIR-related magnetic storms will be difficult to predict and will require
19 better knowledge of the slow solar wind and modeling to solve. For interplanetary space weather,
20 there are challenges for understanding the fluences and spectra of solar energetic particles (SEPs).
21 This will require better knowledge of interplanetary shock properties as they propagate and evolve
22 going from the Sun to 1 AU (and beyond), the upstream slow solar wind and energetic "seed"
23 particles. Dayside aurora, triggering of nightside substorms, and formation of new radiation belts
24 can all be caused by shock and interplanetary ram pressure impingements onto the Earth's
25 magnetosphere. The acceleration and loss of relativistic magnetospheric "killer" electrons and
26 prompt penetrating electric fields in terms of causing positive and negative ionospheric storms are
27 reasonably well understood, but refinements are still needed. The forecasting of extreme events
28 (extreme shocks, extreme solar energetic particle events, and extreme geomagnetic storms
29 ("Carrington" events or greater)) are also discussed. Energetic particle precipitation into the
30 atmosphere and ozone destruction is briefly discussed. For many of the studies, the Parker Solar

31 Probe, Solar Orbiter, Magnetospheric Multiscale Mission (MMS), Arase, and SWARM data will
32 be useful.

33

34

1. INTRODUCTION

35 **1.1. Some Comments on the History of the Physics of Space Weather/Solar Terrestrial** 36 **Physics**

37

38 Space Weather is a new term for a topic/science that actually began over a century and a half ago.
39 Since everything in Solar-Terrestrial Physics (STP) is interconnected we think of STP as the same
40 as Space Weather. It is just that with the space age beginning in 1957 (with the launch of Sputnik)
41 and soon thereafter, many scientifically instrumented satellites led to an explosion of knowledge
42 of the physics of Space Weather. However it is useful to review some of the early scientific studies
43 that occurred prior to 1957. Prior to the space age (where we have satellites orbiting the Earth,
44 probing interplanetary space and viewing the Sun in UV, EUV and X-ray wavelengths), it was
45 clearly realized that solar phenomena caused geomagnetic activity at the Earth. For example,
46 Carrington (1859) noted that there was a magnetic storm that followed ~17 h 40 min after the well-
47 documented optical solar flare which he reported. This storm (Chapman and Bartels, 1940) was
48 only more recently studied in detail by Tsurutani et al. (2003) and Lakhina et al. (2012), but the
49 hints of a causal relationship was there in 1859. After Carrington (1959) published his seminal
50 paper, Hale (1931), Newton (1943) and others showed that magnetic storms were delayed by
51 several days from intense solar flares. These types of magnetic storms are now known to be caused
52 by either their associated interplanetary coronal mass ejections (ICMEs) or their upstream sheaths.
53 Details will be discussed later in this review.

54

55 Maunder (1904) showed that geomagnetic activity often had a ~27 day recurrence. This periodicity
56 was associated with some mysteriously unseen (by visible light) feature on the Sun. Chree (1905,
57 1913) showed that these data were statistically significant, thus inventing the Chree “superposed
58 epoch analysis”, a scientific data analysis technique which is still used today. The mysteriously
59 unseen solar features responsible for the geomagnetic activity were called “M-regions” by Bartels
60 (1934) where the “M” stood for “magnetically active”. It is now known that M-regions are coronal
61 holes (Krieger et al., 1973), solar regions from which solar wind high-speed streams (HSSs)

62 emanate, causing geomagnetic activity at the Earth (Sheeley et al., 1976, 1977; Tsurutani et al.
63 1995). The current status of geomagnetic activity associated with HSSs and the future work
64 needed to better understand and to predict the various facets of Space Weather events will be
65 discussed later.

66
67 With the advent of rockets and satellites, the near-Earth interplanetary medium has been probed
68 by magnetic field, plasma, and energetic particle detectors. The Sun has been viewed in many
69 different wavelengths. The Earth's auroral regions have recently been viewed by UV imagers
70 giving a global view of auroras including the dayside. The ionosphere has been probed by global
71 positioning system (GPS) dual frequency radio signals, allowing a global map of the ionospheric
72 total electron content (TEC) in relatively high spatial and temporal resolution. The purpose of this
73 review article will be to give a reasonably thorough review of some of the major Space Weather
74 effects in the magnetosphere, ionosphere and atmosphere and in interplanetary space, in order to
75 explain what the solar and interplanetary causes are or are expected to be. The most useful part of
76 this review will be to focus on what future advances in Space Weather might be in the next 10 to
77 25 years. In particular we will mention what outstanding problems the Parker Solar Probe, Solar
78 Orbiter, MMS, Arase, ICON, GOLD, and SWARM data might be useful in solving.

79
80 Our discussion will first start with phenomena that occur most frequently during solar maxima
81 (flares, CMEs and ICME-induced magnetic storms). We will explain to the reader what is meant
82 by an ICME and why we distinguish this from a CME. Next, phenomena associated with the
83 declining phase of the solar cycle will be addressed. These include corotating interaction regions
84 (CIRs) and HSSs, which cause high-intensity long-duration continuous AE activity (HILDCAA)
85 events, and the acceleration and loss of magnetospheric relativistic electrons. We will then return
86 to the topic of interplanetary shocks and their acceleration of energetic particles in interplanetary
87 space and also their creating new radiation belts inside the magnetosphere. Interplanetary shock
88 impingement onto the magnetosphere create dayside auroras and also trigger nightside substorms.
89 Prompt penetration electric fields during magnetic storm main phases will be discussed in terms
90 of the consequences of positive and negative ionospheric storms, depending on the local time of
91 the observation and the phase of the magnetic storm. Two relatively new topics, that of
92 supersubstorms (SSSs) and the possibility of precipitating magnetospheric relativistic electrons

93 affecting atmospheric weather will be discussed. A glossary will be provided to give definition of
94 the terms used in this review article.

95
96 There have been some recent books/articles that touch on the many topics of the physics of Space
97 Weather, however not in the same way that we will attempt to do here. We recommend the
98 interested reader: “*From the Sun: Auroras, Magnetic Storms, Solar Flares, Cosmic Rays*” by Suess
99 and Tsurutani (1989), “*Magnetic Storms*” by Tsurutani, Kamide, Gonzalez, Arballo (1997a),
100 “*Storm-Substorm Relationship*” by Sharma, Kamide, Lakhina (2004), “*Recurrent Magnetic*
101 *Storms: Corotating Solar Wind Streams*” by Tsurutani, McPherron, Gonzalez, Lu, Sobral,
102 Gopalswamy (2006a), “*The Sun and Space Weather*” by Hanslmeier (2007), “*Physics of Space*
103 *Storms: From the Solar Surface to the Earth*” by Koskinen (2011), and “*Extreme Events in*
104 *Geospace: Origins, Predictability and Consequences*” by Buzulukova (2018). Because Space
105 Weather is an enormous field/topic, not all facets of it have ever been covered in one book. The
106 present authors are active researchers in the field and will attempt to introduce new viewpoints and
107 topics not covered in the above works.

108 109 **1.2. Organization of Paper**

110
111 The concept of magnetic reconnection is introduced first for the non-space plasma reader.
112 Magnetic reconnection is the physical process responsible for transferring solar wind energy into
113 the magnetosphere during magnetic storms. We have organized the rest of the paper by discussing
114 Space Weather phenomena by solar cycle intervals. However it should be mentioned that this is
115 not totally successful since some phenomena span all parts of the solar cycle.

116
117 Solar maximum phenomena such as CMEs, ICMEs, fast shocks, sheaths, and the forecasting of
118 geomagnetic storms associated with the above are covered in Subsections 2.1 to 2.4. The Space
119 Weather phenomena associated with the declining phase of the solar cycle are discussed in Section
120 3.0. Topics such as CIRs, CIR storms, HSSs, embedded Alfvén wave trains within HSSs,
121 HILDCAA events, relativistic magnetospheric electron acceleration and loss, and electron
122 precipitation and ozone depletion are discussed in Subsections 3.1 to 3.6. Although interplanetary
123 shocks are primarily features associated with fast ICMEs and thus primarily a solar maximum

124 phenomenon, shocks can also bound CIRs (~20% of the time) at 1 AU during the solar cycle
125 declining phase as well. Shocks and the high density plasmas that they create can input ram energy
126 into the magnetosphere. Topics such as solar cosmic ray particle acceleration, dayside auroras,
127 triggering of nightside substorms and the creation of new magnetospheric radiation belts are
128 covered in Subsections 4.1 to 4.4. Solar flares and ionospheric TEC increases is another Space
129 Weather effect causing direct solar-ionospheric coupling not involving interplanetary space nor
130 the magnetosphere. This is briefly discussed in Section 5.0. Prompt penetration electric fields
131 (PPEFs) and ionospheric TEC increases (and decreases) occur during magnetic storms. Although
132 the biggest effects are observed during ICME magnetic storms (solar maximum), effects have been
133 noted in CIR magnetic storms as well. This is discussed in Section 6.0. The “Carrington” magnetic
134 storm is the most intense magnetic storm in recorded history. The aurora associated with the storm
135 reached 23° from the geomagnetic equator (Kimball, 1960), the lowest in recorded history. Since
136 this event has been used as an example for extreme Space Weather and events of this type are a
137 problem for the U.S. Homeland Security, we felt that there should be a separate section on this
138 topic, Section 7.0. We also discuss the possibility of events even larger than the Carrington storm
139 occurring. In Section 8.0 auroral SSSs are discussed. Why is this topic covered in this paper? It
140 is possible that SSSs which occur within superstorms are the actual causes for the extreme
141 ionospheric currents, geomagnetically induced currents (GICs), that are responsible for potential
142 power grid failures and not the geomagnetic storms themselves. Section 9.0 gives our
143 summary/conclusions for the physics and the possibility of forecasting Space Weather events.
144 Section 10.0 is a glossary of Space Weather terms used by researchers in the field. Most of the
145 definitions were carefully constructed in a previous book (Suess and Tsurutani, 1998). These
146 should be useful for an ionospheric person looking up solar terms, etc. It could be particularly
147 useful for the non-space plasma readership as well.

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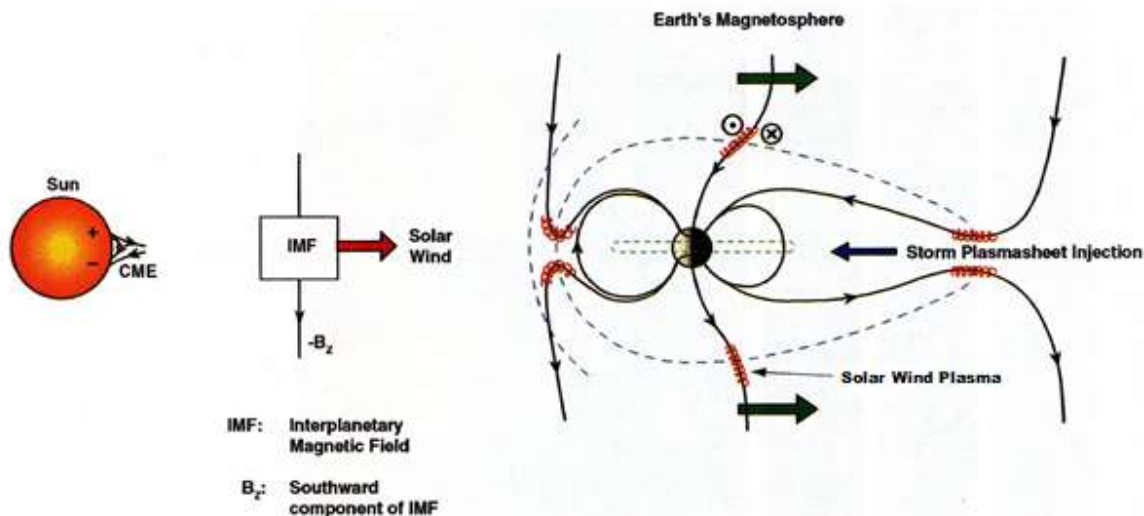
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2. RESULTS: Solar Maximum

2.1. Southward Interplanetary Magnetic Fields, Magnetic Reconnection and Magnetic Storms

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152



153
 154 Figure 1. Magnetic reconnection powering geomagnetic storms and substorms. Adapted from
 155 Dungey (1961).

156
 157 Figure 1 shows the Dungey (1961) scenario of magnetic reconnection. A one-to-one relationship
 158 between southward interplanetary magnetic fields (IMFs) and magnetic storms has been shown by
 159 Echer et al. (2008a) for 90 intense ($Dst < -100$ nT) magnetic storms that occurred during solar
 160 cycle 23. If the IMF is directed southward, it will interconnect with the Earth's magnetopause
 161 northward magnetic fields (the Earth's north magnetic pole is located in the southern hemisphere
 162 near the south rotational pole). The solar wind drags the interconnected magnetic fields and plasma
 163 downstream (in the antisunward direction). The open magnetic fields then reconnect in the tail.
 164 Reconnection leads to strong convection of the plasmashield into the nightside magnetosphere.

165
 166 What is known by theory and verified by observations is that the stronger the southward
 167 component of the IMF and the stronger the solar wind velocity convecting the magnetic field, the
 168 stronger the solar wind-magnetospheric system is driven (e.g., Gonzalez et al., 1994). Intense IMF
 169 B_{south} in MCs (and sheaths) drive intense magnetic reconnection at the dayside magnetopause
 170 and intense reconnection on the nightside. Strong nightside magnetic reconnection leads to strong
 171 inward convection of the plasmashield. The stronger the magnetotail reconnection, the stronger the
 172 inward convection. Via conservation of the first two adiabatic invariants (Alfvén, 1950), the
 173 greater the convection, the greater the energization of the radiation belt particles.

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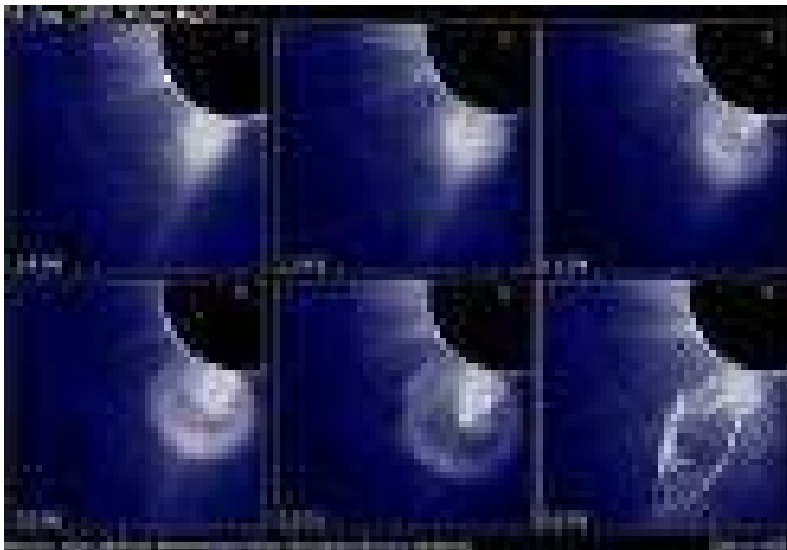
175 As the midnight sector plasmashet is convected inward to lower L, the initially ~100 eV to 1 keV
176 plasmashet electrons and protons are adiabatically compressed (kinetically energized) so that the
177 perpendicular (to the ambient magnetic field) energy becomes greater than the parallel energy.
178 This leads to plasma instabilities, wave growth and wave-particle interactions (Kennel and
179 Petschek, 1966). The resultant effect is the “diffuse aurora” caused by the precipitation of the ~10
180 to 100 keV electrons and protons into the upper atmosphere/lower ionosphere. At the same time
181 double layers are formed just above the ionosphere, giving rise to ~1 to 10 keV “monoenergetic”
182 electron acceleration and precipitation in the formation of “discrete auroras” (Carlson et al., 1998).

183
184 If the IMF southward component is particularly intense, this can lead to a magnetic storm with Dst
185 < -100 nT. The Dst decrease is caused by strong convection of the plasmashet into the inner part
186 of the magnetosphere and the formation of an intensified ring current. This ring current produces
187 a diamagnetic field which causes the reduced field strength at surface of the Earth. This is the
188 magnetic storm main phase.

189
190 After the southward field decreases or changes orientation to northward fields, the magnetic storm
191 recovers. The recovery is associated with a multitude of physical processes associated with the
192 loss of the energetic ring current particles: charge exchange, Coulomb collisions, wave-particle
193 interactions and convection out the dayside magnetopause (West et al., 1972; Kozyra et al. 1997,
194 2006a; Jordanova et al., 1998; Daglis et al. 1999). A typical time for storm recovery is ~10 to 24
195 h (Burton et al., 1975; Hamilton et al., 1988; Ebihara and Ejiri, 1998; O’Brien and McPherron,
196 2000; Dasso et al., 2002; Kozyra et al., 2002; Wang et al., 2003; Weygand and McPherron, 2006;
197 Monreal MacMahon and Llop, 2008).

198
199 **2.2. Coronal Mass Ejections (CMEs), Interplanetary Coronal Mass Ejections (ICMEs) and**
200 **Magnetic Storms**

201



202
 203 Figure 2. A sequence of images showing the emergence of parts of a CME coming from the Sun.
 204 The time sequence starts at the upper left and ends at the lower right. Taken from Illing and
 205 Hundhausen (1986).

206
 207 What are the solar and interplanetary sources of intense IMFs that lead to magnetic reconnection
 208 at Earth and intense magnetic storms? What we know from space age observations is that these
 209 magnetic fields come from parts of a CME, a giant blob of plasma and magnetic fields which are
 210 released from the Sun associated with solar flares and disappearing filaments (Tang et al., 1989).
 211 Figure 2 shows the emergence of a CME from behind a solar occulting disc. The time sequence
 212 starts at the upper left, goes to the right and then to the bottom left, and ends at the bottom right.
 213 The three parts of a CME are best noted in the image on the bottom left. There is a bright outer
 214 loop most distant from the Sun, followed by a “dark region”, and then closest to the Sun is the
 215 solar filament.

216 217 **2.3. Forecasting Magnetic Storms and Extreme Storms Associated with ICMEs**

218
 219 We will precede ourselves and state here that for the limited number of cases studied to date, the
 220 most geoeffective part of the CME is the “dark region”. Interplanetary scientists (Burlaga et al.,
 221 1981; Choe et al., 1982; Tsurutani and Gonzalez, 1994) have identified this as the low plasma beta
 222 region called a magnetic cloud (MC), first identified by Burlaga et al.(1981) and Klein and Burlaga
 223 (1982) in interplanetary space by magnetic field and plasma measurements. When there are

224 southward component magnetic fields within the MC (thought to typically be a giant fluxrope), a
225 magnetic storm results (Gonzalez and Tsurutani, 1987; Gonzalez et al. 1994; Tsurutani et al.,
226 1997b; Zhang et al., 2007; Echer et al. 2008a).

227

228 It should be noted that fast CMEs and intense MC fields are relatively rare. The SOHO LASCO
229 instrument has observed $> 10,000$ CMEs but only $\sim 5\%$ have speeds faster than ~ 700 km/s. Only
230 very few have speeds $> 2,000$ km/s and these are coming from coronal regions associated with
231 Active Regions (ARs) (Yashiro et al. 2004).

232

233 Interplanetary and magnetospheric scientists have developed the term ICME or interplanetary
234 CME because it is not currently known (for individual events) how the CME evolves as it
235 propagates from the Sun to the Earth and beyond. Leamon et al. (2004) in comparing interplanetary
236 MCs to associated solar active regions found that there was little or no relationship, compelling
237 the authors to conclude that “MCs are formed during magnetic reconnection and are not simple
238 eruptions of preexisting coronal structures”. Yurchyshyn et al. (2007) in a similar study found that
239 “for the majority of interplanetary MCs, the fluxrope axis orientation changed less than 45° going
240 from the Sun to 1 AU”. Palmerio et al. (2018) found “for the majority of cases, the flux rope tilt
241 angles rotated several tens of degrees (between the Sun and the Earth) while 35% changed by more
242 than 90° ”. 3D MHD simulations have shown that CMEs can be severely distorted as they interact
243 with different types of interplanetary structures as they propagate through interplanetary space
244 (Odstrcil and Pizzo, 1999a,b). The latter authors have shown that the CME distortion is
245 substantially different when it interacts with the streamer belt (heliospheric plasma sheet/HPS)
246 than with an HSS. The distortion of the CME can make the ICME unrecognizable at a distance
247 further away from the Sun.

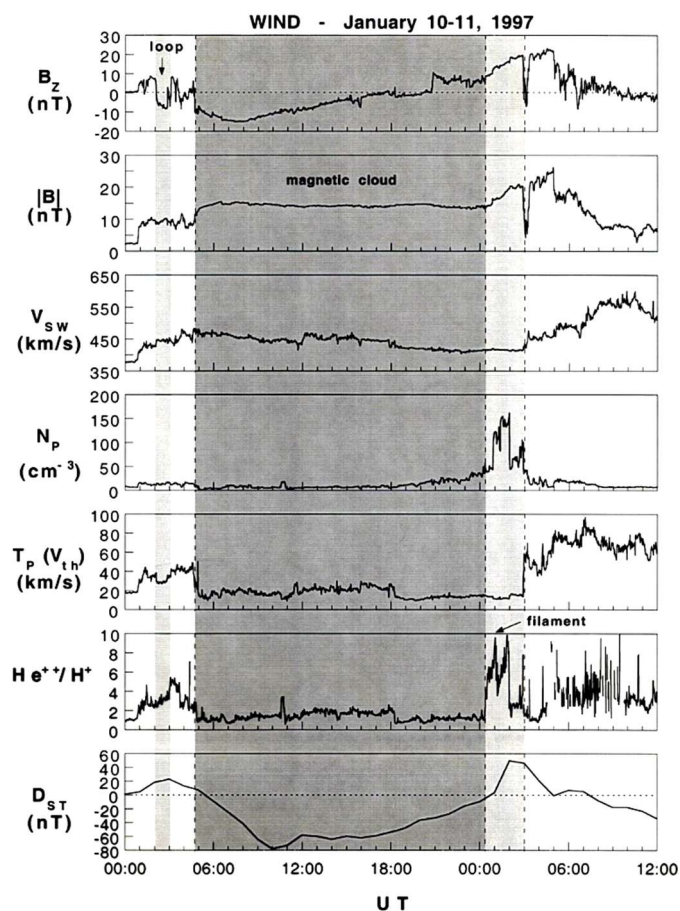
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249 More detailed topics not covered in Palmerio et al. (2018) or in Odstrcil and Pizzo (1999a,b) are
250 the topics of the fate of the principal features of CMEs as discussed by Illing and Hundhausen
251 (1986). For example, the bright outer loops are seldomly identified at 1 AU (one rare case was
252 identified by Tsurutani et al., 1998) and the filaments are typically not found within the ICME at
253 1 AU. The first filament detection at 1 AU was not reported until 1998 (Burlaga et al., 1998). For
254 more recent observations of filaments at 1 AU, we direct the reader to Lepri and Zurbuchen (2010).

255 Where have the bright outer loops and filaments gone to? Have they simply detached only to
256 impinge onto the magnetosphere at a later time, or do they go back into the Sun? Or is it possible
257 that many CMEs do not have filaments at their bases? Remote imaging observations from
258 STEREO should be able to answer these questions. New in situ results from Parker Solar Probe,
259 Solar Orbiter and ACE plus ground-based solar observations could perhaps help address the
260 plasma physics of why typical ICMEs do not have attached filaments.

261
262 It should be remarked that the high-density solar filaments could be extremely geoeffective if they
263 collided with the Earth's magnetosphere (this is covered later in Section 3.2.5). Is it possible for
264 the MC to rotate so that initially southward magnetic fields become northward components? Can
265 the MC fields be compressed or expanded by interplanetary interactions? Can magnetic
266 reconnection be taking place within the ICME between the solar corona and 1 AU as suggested by
267 Manchester et al. (2006) and Kozyra et al., (2013)? If so, how often does this occur and can it be
268 predicted? Modeling and examining the Parker Solar Probe and Solar Orbiter data (for studies on
269 the same ICME) could help us understand whether the MCs evolve as they propagate through
270 interplanetary space.

271
272 Of course, the most important goal for Space Weather is predicting the southward magnetic fields
273 within the ICME. This extremely difficult task is the holy grail of Space Weather. It is more
274 important than predicting the time of the release of a CME, its speed and its direction.



275

276

277 Figure 3. An ICME detected at 1 AU just upstream of the Earth.

278

279 Figure 3 shows a rare case of an ICME at 1 AU where all three parts of a CME are detected. The
 280 MC is indicated by the shaded region in the figure. The outer loop was identified by Tsurutani et
 281 al. (1998) and the filament by Burlaga et al. (1998).

282

283 From top to bottom are the IMF B_z component (in geocentric solar magnetospheric/GSM
 284 coordinates), the field magnitude, the solar wind velocity, density, temperature and the $\text{He}^{++}/\text{H}^+$
 285 ratio. The bottom panel gives the ground based Dst index whose amplitude is used as an indicator
 286 of the occurrence of a magnetic storm. Dst becomes negative when the Earth's magnetosphere is
 287 filled with storm-time energetic $\sim 10\text{-}300$ keV electrons and ions (Williams et al., 1990). Dessler
 288 and Parker (1959) and Sckopke (1966) have shown that the amount of magnetic decrease is linearly
 289 related to the total kinetic energy of the enhanced radiation belt particles. This is because the

290 energetic particles which comprise the storm-time ring current, through gradient drift of the
291 charged particles, form a diamagnetic current which decreases the Earth's magnetic field inside
292 the current. We refer the reader to Sugiura (1964) and Davis and Sugiura (1966) for further
293 discussions of the Dst index. The Dst index is a one hr index. More recently a 1 min SYM-H index
294 (Iyemori, 1990; Wanliss and Showalter, 2006) has been developed. This is more useful for high
295 time resolution studies. Both indices are produced by the Kyoto Data Center.

296
297 In this example (top panel of Figure 3) the MC fields start with a strong southward ($B_z < 0$ nT)
298 component and then later turns northward. In the bottom panel, the magnetic storm Dst index
299 becomes negative with very little delay from the southward magnetic fields. The energy transfer
300 mechanism is magnetic reconnection, as discussed earlier in Section 2.1. The high-density
301 filament (fourth panel from the top) is present after the MC passage. Values as high as ~ 160 cm⁻³
302 have been detected. These values are extreme values (the nominal solar wind density is ~ 3 to 5
303 cm⁻³: Tsurutani et al., 2018a). The high densities impinging on the magnetosphere in this case
304 caused compression of the magnetosphere and the Dst index to reach $\sim +55$ nT.

305
306 The stronger the southward component of the MC fields, the more intense the magnetic storm at
307 the Earth. In extreme cases storms with intensities of $Dst < -250$ nT can occur (Tsurutani et al.
308 1992a; Echer et al. 2008b). An empirical relationship between the speed of the MC at 1 AU and
309 its magnetic intensity has been shown by Gonzalez et al. (1998). A hypothetical explanation is the
310 "melon seed model": squeezing a melon seed will cause it to squirt out, squeezing it harder will
311 make it come out fast. A larger magnetic field will require greater pressure to release it. However
312 a substantial MHD or plasma kinetic model is needed to explain the physics of this empirical
313 relationship in more detail.

314
315 Because extremely strong MC magnetic fields are needed to produce extreme magnetic storms
316 like the "Carrington" event (Tsurutani et al., 2003; Lakhina and Tsurutani, 2017), one should focus
317 on extremely fast events for forecasting purposes. The geoeffective interplanetary dawn-to-dusk
318 electric field is $V_{sw} \times B_{south}$. Because Gonzalez et al (1998) have shown that $|B|$ is empirically
319 proportional to V_{sw} , the dawn-to-dusk interplanetary electric field has a V_{sw}^2 dependence. The
320 Carrington ICME took ~ 17 hr 40 min to go from the Sun to Earth (Carrington, 1859), causing the

321 largest magnetic storm in history. The minimum Dst has been estimated to be -1760 nT. However
322 the August 1972 event was even faster, taking only ~14 h 40 min to go from the Sun to Earth
323 (Vaisberg and Zastenker 1976; Zastenker et al. 1978). Although the 1972 MC was indeed extreme
324 in speed and magnetic field intensity, the direction of the magnetic field was northward and thus
325 there was geomagnetic quiet following the MC impingement onto the magnetosphere (Tsurutani
326 et al. 1992b). So again, predicting the ICME magnetic field direction is paramount in importance
327 for Space Weather applications.

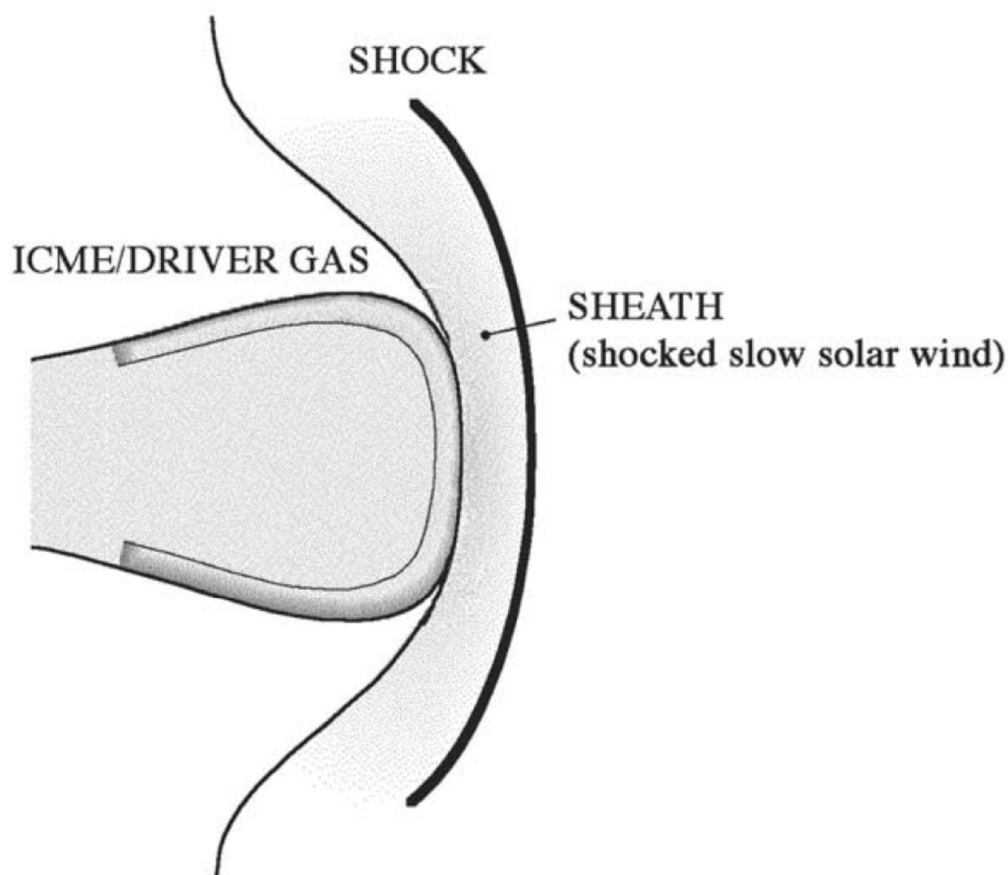
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329 Modeling ICME propagation in interplanetary space during disturbed AR periods has met only
330 limited success (Echer et al., 2009; Mostl et al., 2015; Hajra et al., 2019). Sometimes it is difficult
331 to even identify to which flare or disappearing filament a detected ICME is related (see Tang et
332 al., 1989). The propagation times from the Sun to 1 AU has often been in error by days (Zhao and
333 Dryer, 2014). The additional information provided by the Parker Solar Probe and Solar Orbiter
334 and examination of present ICME propagation codes could help improve the ability to make more
335 accurate forecasts.

336

337 **2.4. Fast Shocks, Sheaths and Magnetic Storms**

338



339
 340 Figure 4. A schematic of an interplanetary sheath antisunward of an ICME. In this diagram the
 341 Sun is on the left (not shown).

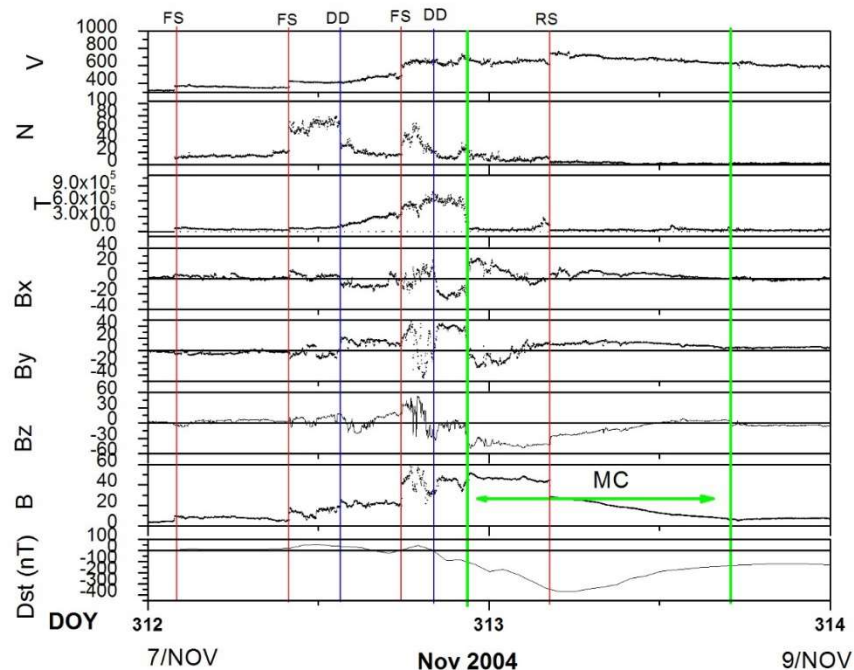
342
 343 Figure 4 shows a schematic of a shock and sheath upstream of an ICME. “Fast” CMEs/ICMEs can
 344 create upstream fast forward shocks (Tsurutani et al., 1988). By “fast” it is meant that the
 345 CME/ICME is moving at a speed higher than the upstream magnetosonic (fast wave mode) speed
 346 relative to the upstream plasma and by “forward” we mean that the shock is propagating in the
 347 same direction as the “driver gas” or the CME/ICME, antisunward. When a shock is formed, it
 348 compresses the upstream plasma and magnetic fields. In this terminology, the upstream direction
 349 is the direction in which the shock is propagating (antisunward in this case) and the downstream
 350 direction is towards the Sun (see Kennel et al., 1985 and Tsurutani et al., 2011 for details on
 351 shocks). The compressed plasma and magnetic fields downstream of the shock is the “sheath”.
 352 The shock and sheath are not part of the CME/ICME. The origin of this plasma and magnetic
 353 fields is the slow solar wind altered by shock compression. This is important to understand if one

354 wishes to predict magnetic storms caused by interplanetary sheath southward magnetic fields. It
 355 should be noted that “slow” ICMEs have been detected at 1 AU (Tsurutani et al., 1994a). These
 356 phenomena do not have upstream shocks and sheaths, as expected. However the southward MC
 357 magnetic fields still cause magnetic storms.

358

359 Kennel et al. (1985) used MHD simulations to show that the plasma densities and magnetic field
 360 magnitudes downstream of shocks are roughly related to the shock magnetosonic Mach numbers.
 361 This theoretical relationship holds up to a Mach number of ~ 4 . For higher Mach numbers MHD
 362 predicts that the compression will remain at a factor of ~ 4 . Since interplanetary shocks detected
 363 at 1 AU typically have Mach numbers only of 1 to 3 (Tsurutani and Lin, 1985; Echer et al., 2011;
 364 Meng et al. 2019), 1 to 3 are the typical shock magnetic field and density compression ratios
 365 detected at 1 AU. One question for future studies is “does the MHD relationships of magnetic
 366 field magnitude and density jumps hold for extreme shocks?” If not, there will be important
 367 consequences for extreme Space Weather.

368



369

370 Figure 5. An example of three fast forward shocks pumping up the interplanetary magnetic field
 371 intensity. Taken from Tsurutani et al. (2008a).

372
373 Figure 5 shows a complex interplanetary event that was selected by the CAWSES II team to study
374 in detail. The full information on this event from the Sun to the atmosphere can be found in the
375 special issue: Large Geomagnetic Storms of Solar Cycle 23
376 ([https://agupubs.onlinelibrary.wiley.com/doi/toc/10.1002/\(ISSN\)1944-8007.CYCLE231](https://agupubs.onlinelibrary.wiley.com/doi/toc/10.1002/(ISSN)1944-8007.CYCLE231)). What
377 is important is that this event was associated with a solar active region (AR) and the results are
378 quite important in terms not only for interplanetary disturbance phenomena but also for
379 geomagnetic activity at the Earth.

380
381 From top to bottom in Figure 5 are the solar wind speed, density, and temperature, the IMF B_x,
382 B_y and B_z components and the magnetic field magnitude in solar magnetospheric (GSM)
383 coordinates. In this coordinate system, **x** points in the direction of the Sun, the **y** direction is given
384 by $(\Omega \times \mathbf{x})/|\Omega \times \mathbf{x}|$ where Ω is the Earth's south magnetic pole (the south magnetic pole is near the
385 north geographic pole) and **z** axis, which is in the plane containing both the Earth-Sun line and the
386 dipole axis, completes the right hand system. The magnetic storm Dst index is given at the bottom.
387 Fast forward shocks are denoted by the three vertical red lines on 7 November 2004. There are
388 sudden increases in the velocity, density, temperature and magnetic field magnitude at all three
389 events. The Rankine-Hugoniot relationships have been applied to the plasma and magnetic field
390 data and the analysis did determine that they are indeed fast shocks.

391
392 The point of showing this interplanetary event is to indicate that each shock pumps up the
393 interplanetary sheath magnetic field by factors of ~2 to 3. The initial magnetic field magnitude
394 started with a value of ~4 nT and at the peak value after the three shocks, it reached a value of ~60
395 nT. This final value was higher than the MC magnetic field, which was ~45 nT. Details
396 concerning the shocks and compressions can be found in the original paper for readers who are
397 interested. What is important here is how intense interplanetary magnetic fields are created. They
398 can come from the MCs themselves or the sheaths, as shown here. However, in this case the
399 southward magnetic fields that caused the magnetic storm came from the MC and not the sheath.

400
401 In the above example it is believed that three fast forward shocks were associated with three ICMEs
402 released from the AR. The longitudinal extent of shocks are, however, wider than the MCs, so

403 only one MC was detected in the event. A similar situation was found for the August 1972 event
404 discussed earlier.

405
406 It should be noted that a fast reverse wave (here by “reverse” we mean that the wave is propagating
407 in the solar direction) was detected during the Figure 5 event. It is identified as the red vertical line
408 on 8 November. In detailed examination of the Rankine-Hugoniot conservation equations, this
409 wave was found to propagate at a speed below the upstream magnetosonic speed and thus was a
410 magnetosonic wave and not a shock. This reverse wave caused a decrease in the MC magnetic
411 field (and the southward component) and thus the start of the recovery phase of the magnetic storm.
412 The reader should note that fast reverse waves and shocks are also important for geomagnetic
413 activity. A detailed discussion of shock and discontinuity effects on geomagnetic activity can be
414 found in Tsurutani et al. (2011).

415

416 **2.4.1. Forecasting ICME sheath magnetic storms**

417

418 Determination of the IMF Bz component in the sheaths will be a difficult task. To do this, more
419 effort in understanding the slow solar wind plasma, magnetic fields and their variations will be
420 required. To date, there has been little effort expended in this area. This is, however easy for us to
421 hope for, but in practice is far more difficult to do. Use of data from Solar Probe, Solar Orbiter and
422 a 1 AU spacecraft such as ACE could help in these analyses.

423

424 This problem has recently been emphasized by results from Meng et al. (2019). Meng et al. have
425 shown that superstorms ($Dst < -250$ nT) that occurred during the space age (1957 to present) are
426 mostly driven by sheath fields or a combination of sheath plus a following magnetic cloud (MC).

427

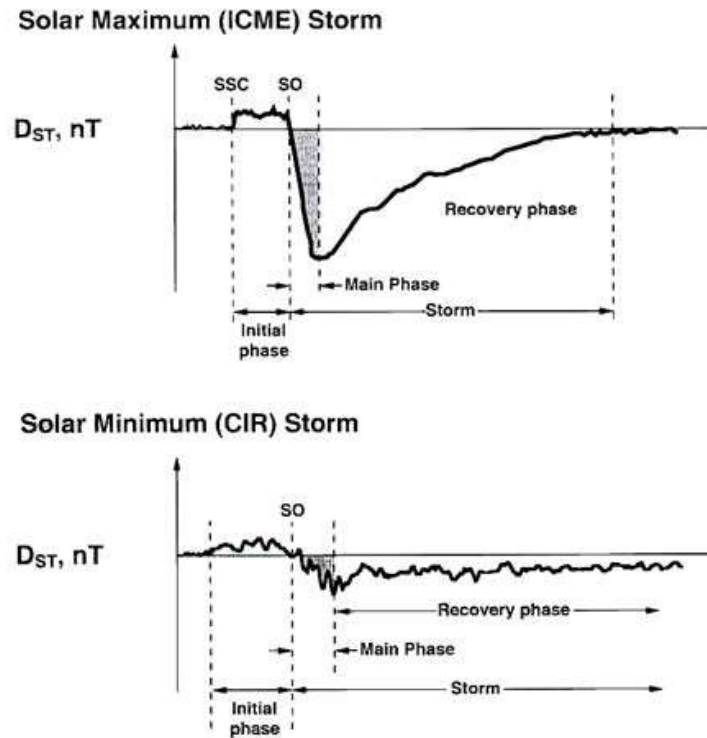
428 Substorms are generated by lower intensity southward magnetic fields with the process of
429 magnetic reconnection being the same as above. However substorm plasmashet injections only
430 go in to $L \sim 4$, the outer part of the magnetosphere (Soraas et al., 2004). The auroras associated
431 with substorms appear in the “auroral zone”, 60° to 70° magnetic latitude (MLAT). Magnetic
432 storms associated with much larger IMF B_{south} are detected at subauroral zone latitudes.

433

3. RESULTS: Declining Phase of the Solar Cycle

434
435
436
437

3.1. Corotating Interaction Region (CIR) Magnetic Storms



438

439 Figure 6. The magnetic Dst profiles of a CIR magnetic storm (bottom) and an ICME magnetic
440 storm (top). Taken from Tsurutani (2000).

441

442 During the declining phase of the solar cycle a different type of solar and interplanetary activity
443 dominates the physical cause of magnetic storms, that of Corotating Interaction Regions (CIRs).
444 HSSs emanating from coronal holes (CHs) interact with the slow solar wind and form CIRs at their
445 interaction interfaces. The magnetic storms caused by CIRs are quite different from storms caused
446 by ICMEs and/or their sheaths. Figure 6 shows the difference in profiles of two different types of
447 magnetic storms. The profile of a CIR magnetic storm is shown on the bottom and that of a shock
448 sheath ahead of an ICME MC magnetic storm on top.

449

450 The ICME MC magnetic storm Dst profile, discussed briefly earlier (see Figure 3), is reasonably
451 easy to identify (top panel). There is a sudden, ~tens of second duration positive increase in Dst

452 which is caused by the sudden increase in solar wind ram pressure due to the passage of the sheath
453 high density jump downstream of the shock. This compresses the magnetosphere, creating the
454 sudden impulse (SI^+ : see Joselyn and Tsurutani, 1990) detected everywhere on the ground (Araki
455 et al., 2009). Later, in either the sheath or the MC there may be a southward IMF which causes
456 the magnetic storm. If there is a southward component in the MC, it is usually smoothly varying
457 in intensity and direction. This leads to a smooth monochromatic storm main phase as seen in the
458 Dst index (and illustrated in Figures 3 and 6). The loss of the ring current particles is the cause of
459 the storm recovery phase. The details of storm recovery phase durations and causative mechanisms
460 will be an interesting topic for magnetospheric scientists to study in the near future. The Arase
461 mission data will be quite useful for these studies.

462
463 The bottom panel of Figure 6 shows the typical profile of a CIR magnetic storm. It is quite different
464 from a sheath-MC magnetic storm profile. There is no SI^+ associated with the beginning of the
465 geomagnetic disturbance. This is because CIRs detected at 1 AU typically are not led by fast
466 forward shocks (Smith and Wolf, 1976; Tsurutani et al. 1995). The positive increase in Dst is
467 associated with the impact of a high density region near the heliospheric current sheet (HCS)
468 (Smith et al., 1978; Tsurutani et al. 2006b) called the heliospheric plasmashet (HPS; Winterhalter
469 et al., 1994) and/or associated with the compressed plasma at the leading edge of the CIR. These
470 are slow solar wind plasma densities. The most distinguishing feature of the CIR storm main phase
471 is the lack of smoothness, in sharp contrast to the MC magnetic storm. This irregular Dst storm
472 main phase is caused by large B_z fluctuations within the CIR.

473
474 CIR magnetic fields have magnitudes of ~ 20 to 30 nT and typically do not reach the much higher
475 intensities that MC fields typically do. For this reason and also because of the IMF B_z fluctuations,
476 CIR magnetic storms usually have intensities $Dst \geq -100$ nT (small or no magnetic storms).
477 Extreme magnetic storms with $Dst < -250$ nT caused by CIRs are rare, if they occur at all (none
478 were found in the Meng et al. 2019 study). However it is clear that compound events involving
479 both CIRs, sheaths ahead of ICMEs and ICMEs could certainly cause extreme magnetic storm
480 events.

481

482 CIR related magnetic storms occur most frequently during the declining phase of the solar cycle
483 and ICME magnetic storms typically occur near the maximum phase of the solar cycle. However,
484 it should be noted that both CIR storms and sheath and/or ICME MC magnetic storms can occur
485 during any phase of the solar cycle. We have simply ordered things by solar cycle so that it will
486 be easier to give the reader the general picture of Space Weather.

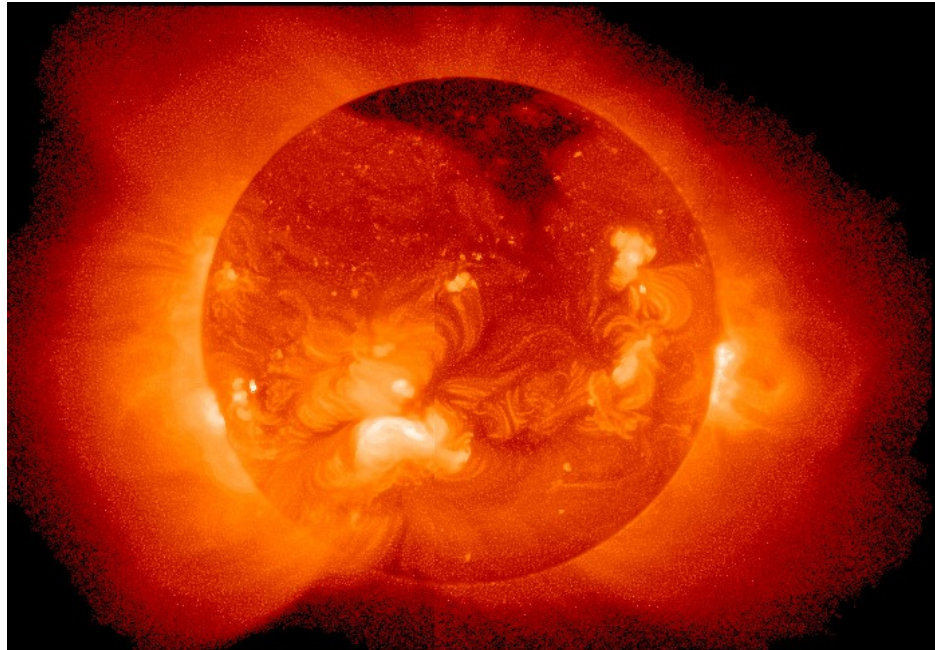
487

488 **3.2. Coronal Holes, High Speed Solar Wind Streams and Geomagnetic Activity**

489

490 **3.2.1. Coronal holes and high-speed solar wind streams**

491



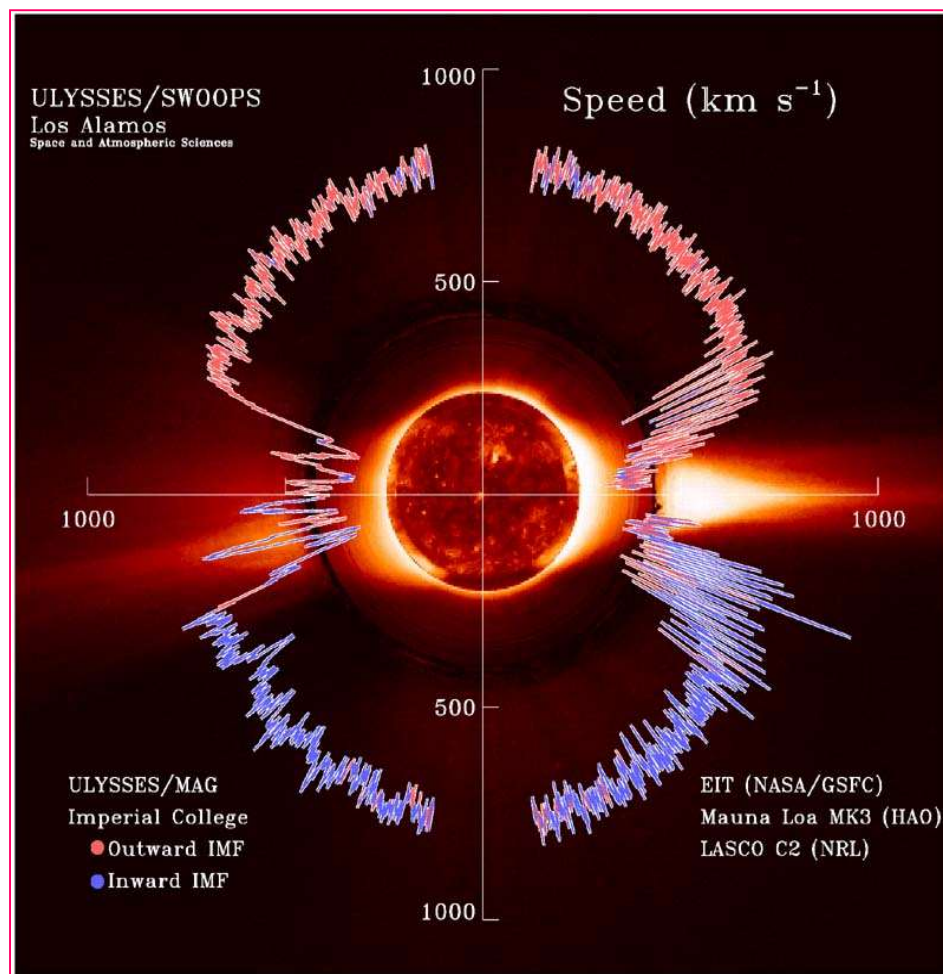
492

493 Figure 7. A large coronal hole (the dark region) near the north pole of the Sun. The figure was
494 taken by the soft X-ray telescope (SXT) onboard the Yokoh satellite in 1992.

495

496 Figure 7 shows a polar coronal hole at the north pole of the Sun. This image was taken by the soft
497 x-ray telescope (SXT) onboard the Yokoh satellite
498 (http://www.spaceweathercenter.org/swop/Gallery/Solar_pics/yohkoh_060892.html). The dark
499 (low temperature) region at the pole is the coronal hole. Large polar coronal holes occur typically
500 in the declining phase of the solar cycle (Bravo and Otaola, 1989; Bravo and Stewart, 1997; Zhang
501 et al., 2005).

502



503

504 Figure 8. High speed solar wind streams emanating from coronal holes in the north and south
 505 solar poles. The figure was taken from Phillips et al. (1995) and McComas et al. (2002).

506

507 Figure 8 gives a "dial plot" of the solar wind speed for the first traversal of the Ulysses spacecraft
 508 over the Sun's poles. The radius from the center of the Sun to the trace indicates the solar wind
 509 speed. The magnetic field polarity is indicated by the color of the trace, red for outward IMFs and
 510 blue for inward IMFs. A SOHO EIT soft x-ray image of the Sun is placed at the center of the figure
 511 and a High Altitude Observatory Mauna Loa coronagraph image shows the inner corona at that
 512 time. The outer corona is an image taken by the SOHO C2 coronagraph.

513

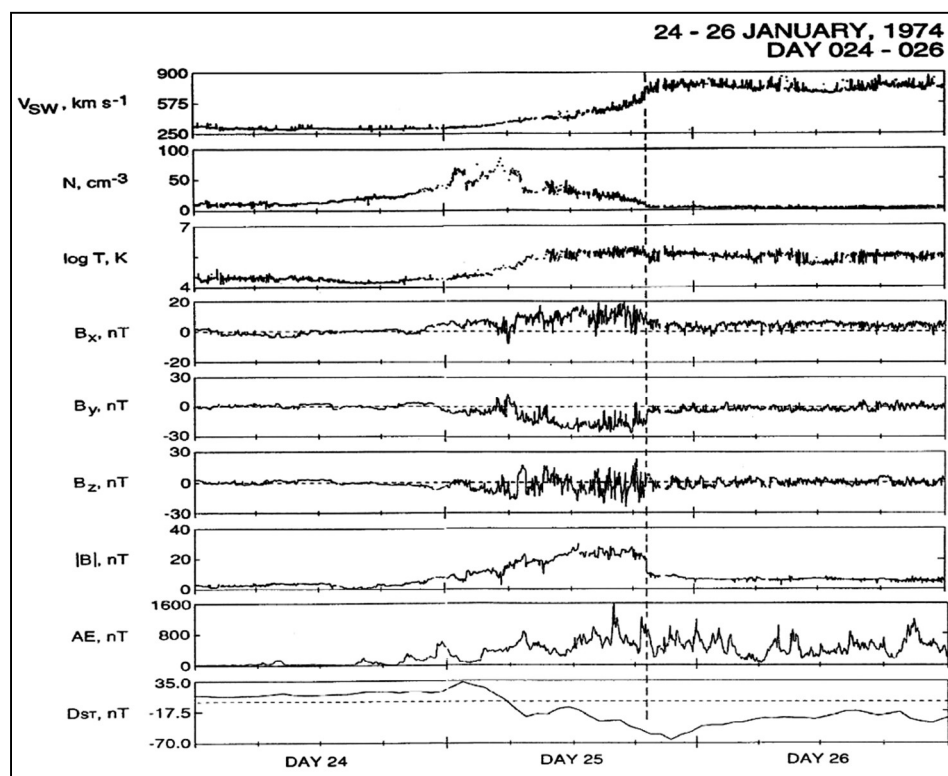
514 Two large polar coronal holes are detected at the Sun, one at the north pole and the other at the
 515 south pole. It is noted that HSSs of ~750 to 800 km/s are detected at Ulysses when over the polar

516 coronal hole regions. When Ulysses was near the solar equatorial region where helmet streamers
 517 are present, the solar wind speeds are of the slow solar wind variety, $V_{sw} \sim 400$ km/s. The reader
 518 should note that it took years for Ulysses to make this polar orbit while the solar and coronal
 519 images were taken at one point in time. However, this composite figure is useful to illustrate the
 520 main points about the origins of HSSs.

521

522 3.2.2. High speed solar wind streams and the formation of CIRs

523



524

525 Figure 9. A high-speed solar wind stream-slow solar wind interaction and the formation of a CIR
 526 during January 1974. The format is the same as in Figure 4 except that the AE index is given in
 527 the next to bottom panel. The figure is taken from Tsurutani et al. (2006b).

528

529 Figure 9 shows a HSS-slow speed stream interaction during January 1974. The right portion of the
 530 top panel on day 26 shows a HSS with speeds of 750-800 km/s at 1 AU. On day 24, the top panel
 531 left indicates a solar wind speed of ~ 300 km/s, or the slow solar wind. The effects of the stream-
 532 stream interaction occur on day 25. This is best seen in the IMF magnitude panel, 7th from the top.
 533 The stream-stream interaction creates intense magnetic fields of ~ 25 nT. The 6th from the top

534 panel is the IMF Bz component (in GSM coordinates). The Bz is highly fluctuating. Magnetic
535 reconnection between the IMF southward components and the magnetopause magnetic fields leads
536 to the irregularly shaped storm main phase shown in the bottom (Dst) panel.

537

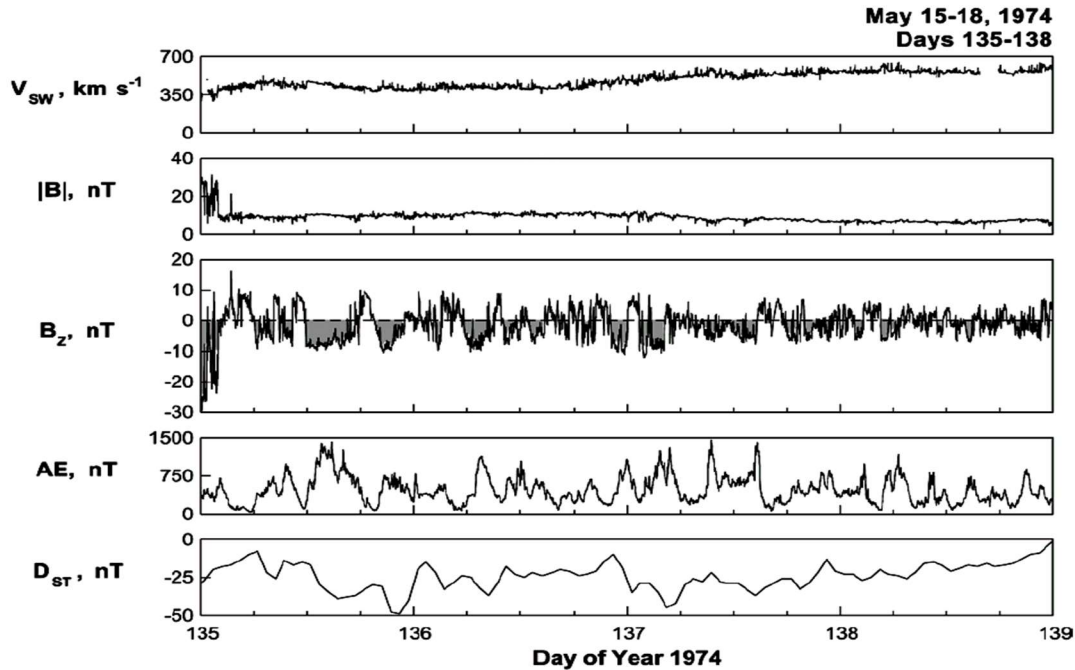
538 To be able to forecast a CIR magnetic storm, one would have to first understand the sources of the
539 IMF Bz fields. For example, are they compressed upstream Alfvén waves (Tsurutani et al. 1995,
540 2006c)? Or could they be waves generated by the shock interaction with upstream waves in the
541 slow solar wind? That would be only the first step for forecasting, of course. Then with knowledge
542 of the properties of the slow speed stream, the details of the wave compression/interaction would
543 then have to be calculated/modeled.

544

545 Another approach would be to determine if there is an underlying southward component of the
546 IMF within the CIR. This would most likely be caused by the geometry of the HSS-slow speed
547 stream interaction and may be predictable from MHD modeling. If this is correct, then the wave
548 fluctuations can be modeled as being superposed on top of these DC magnetic fields. In (rare)
549 cases of radial alignment, Solar Probe closest to the Sun could characterize sheath fields. The
550 evolution of those fields would be detected by Solar Orbiter. Simulation of further evolution could
551 be applied and predictions of the fields at 1 AU could be tested by ACE data. If there are waves
552 generated by the shock, then the above scenario would not work as well as expected, or at least
553 would be more complicated to apply in a useful manner.

554

555 **3.2.3. High speed solar wind streams, Alfvén waves and HILDCAAs**



556
 557 Figure 10. A high-intensity, long-duration continuous AE activity (HILDCAA) event during 1974.
 558 Taken from Tsurutani et al. (2006c).

559
 560 The schematic in Figure 6 showed a long “recovery phase” that trails the CIR magnetic storm main
 561 phase (see Tsurutani and Gonzalez, 1987). However, we now know that the storm wasn’t
 562 “recovering” as in the case of an MC magnetic storm recovery but that something else was
 563 occurring. This “recovery” can last from days to weeks. Thus, processes of charge exchange,
 564 Coulomb collisions, etc. for ring current particle losses are not tenable to explain such long
 565 “recoveries”.

566
 567 Figure 10 shows the interplanetary cause of this extended geomagnetic activity. It occurs primarily
 568 during HSSs independent of whether a CIR magnetic storm occurred prior to it or not (Tsurutani
 569 and Gonzalez, 1987; Tsurutani et al., 1995, 2006b; Kozyra et al. 2006b; Turner et al. 2006; Hajra
 570 et al. 2013, 2014a, 2014b, 2014c, 2017). From top to bottom are the solar wind speed, the IMF
 571 magnitude, the IMF Bz component (in GSM coordinates) and the auroral electrojet (AE) index.
 572 The bottom panel is the Dst index.

573

574 The interplanetary data were taken from the IMP-8 spacecraft, an Earth orbiting satellite that was
575 located upstream of the magnetosphere in the solar wind at this time. The location was inside 40
576 Re, where an Re is an Earth radius. The magnetic Bz fluctuations have been shown to be Alfvén
577 waves which are of large nonlinear amplitudes in HSSs (Belcher and Davis, 1971; Tsurutani and
578 Gonzalez, 1987; Tsurutani et al., 2018b). What is apparent from this figure is that every time the
579 IMF Bz is negative (southward), there is an AE increase and a Dst decrease. This has been
580 interpreted as being due to magnetic reconnection between the southward components of the
581 Alfvén waves and the Earth magnetopause. The AE is enhanced by the same magnetic
582 reconnection process that occurs during substorms, and a small parcel of plasmashet plasma is
583 injected into the nightside magnetosphere causing the Dst index to decrease slightly. It is noted
584 that there are many southward IMF Bz dips in this four day interval of data shown in Figure 10.
585 There are also many corresponding AE increases and Dst decreases. Thus, the interpretation of the
586 constant/average Dst value of ~ -25 nT for four days is that continuous plasma injection and decay
587 is occurring. This is clearly not a “recovery phase” where the ring current particles are simply
588 lost, it only appears as a recovery from the Dst trace. Soraas et al. (2004) have shown that particles
589 are injected during these events but only to L values of 4 and greater (the L = 4 magnetic field line
590 is the dipole magnetic field that crosses the magnetic equator a distance 4 Earth radii from the
591 center of the Earth). These are shallow injections as suggested above.

592
593 These geomagnetic activity events have been named High-Intensity, Long-Duration Continuous
594 AE events or HILDCAAs (Tsurutani and Gonzalez, 1987). This name is simply a description of
595 the events without an interpretation. In 2004 when a detailed examination using Polar EUV auroral
596 imaging was applied, it was found that many phenomena besides simple isolated substorms
597 occurred (Guarnieri, 2006; Guarnieri et al., 2006). Although substorms occur during HILDCAA
598 events, there are AE increases (injection events?) that are not well-correlated with substorm onsets
599 (Tsurutani et al., 2004b). The full extent of HILCAAs is not well understood (see also Souza et
600 al., 2016, 2018; Mendes et al., 2017). By using IMAGE auroral observations and geomagnetic
601 indices to identify convection events which are not classical Akasofu (1964) substorms, the fields
602 and particle data from SWARM, MMS and Arase could be used to characterize the physics
603 properties of these “convection” events.

604

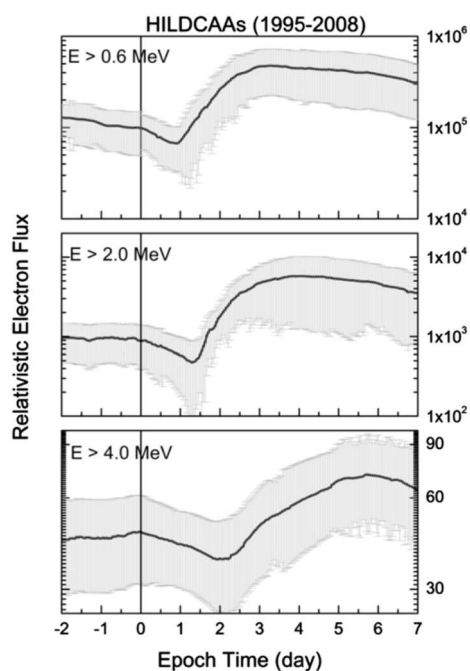
605 There is also the question of the origin of the interplanetary Alfvén waves? Do they originate at
 606 the Sun caused by supergranular circulation, or is that mechanism untenable as argued by Hollweg
 607 (2006)? Could the waves be generated locally between the Sun and Earth as speculated by Matteini
 608 et al. (2006, 2007) and Hellinger and Travnicek (2008)? Parker Solar Probe could identify Alfvén
 609 waves within high speed streams and Solar Orbiter (when radially aligned) could determine the
 610 wave evolution.

611
 612 The original requirement for identifying a HILCAA event was quite strict. The event had to occur
 613 outside of a magnetic storm main phase (Dst was required to be > -50 nT: Gonzalez et al. 1994),
 614 the peak AE intensity had to be greater than 1,000 nT (high-intensity), the event had to last longer
 615 than 2 days (long-duration), and there could not be any dips in AE less than 200 nT for longer than
 616 two hours (continuous). Clearly there are events with the same interplanetary causes and
 617 geomagnetic effects as for the strict definition. However, the strict definition is useful for further
 618 studies using different data sets.

619

620 3.2.4. HILDCAAs and the Acceleration of Relativistic Magnetospheric Electrons

621



622
 623 Figure 11. The relationship between HILDCAAs and relativistic electron acceleration. The figure
 624 is taken from Hajra et al. (2015a).

625
626 One of the consequences of HSSs and HILDCAAs is the acceleration of relativistic (\sim MeV)
627 electrons. These energetic particles can damage orbiting satellite electronic components (Wrenn,
628 1995), and thus are known as “killer electrons”. Figure 11 shows the relationship between the onset
629 of HILCAA events (vertical line) and relativistic electron fluxes. From top to bottom are the $E >$
630 0.6 MeV, the $E > 2.0$ MeV and the $E > 4.0$ MeV electron fluxes detected by the GOES-8 and
631 GOES-12 satellites located at $L = 6.6$. This figure is a superposed epoch analysis (Chree, 1913)
632 result of 35 HILDCAA events in solar cycle 23, from 1995 to 2008, which are not preceded by
633 magnetic storms. This was done to avoid contamination by storm-time particle acceleration (by
634 intense convection/compression). The zero-epoch time (vertical line) corresponds to the
635 HILDCAA onset time. Here the “strict” definition of HILDCAAs was used to define the onset
636 times.

637
638 The figure shows that the flux enhancement of $E > 0.6$ MeV electrons is statistically delayed by
639 ~ 1.0 day from the onset of the HILDCAAs. The $E > 4.0$ MeV electrons are statistically delayed
640 by ~ 2.0 days from the HILDCAA onset. It is thus possible that HILCAAs may be used to forecast
641 relativistic electron flux enhancements in the magnetosphere (see Hajra et al., 2015b; Tsurutani et
642 al., 2016a; Hajra and Tsurutani, 2018a; Guarnieri et al., 2018). This however has not been done
643 yet and could be implemented by scientists today.

644
645 The physics for electron acceleration to relativistic (\sim MeV) energies has been well-developed by
646 magnetospheric scientists. Two competing acceleration mechanisms have been developed. In one
647 mechanism, with each injection of plasmashet particles on the nightside magnetosphere, the
648 anisotropic ~ 10 to 100 keV electrons generate electromagnetic whistler mode chorus waves
649 (Tsurutani and Smith, 1974; Meredith et al. 2002) by the loss cone/temperature anisotropy
650 instability (Brice, 1964; Kennel and Petschek, 1966; Tsurutani et al., 1979; Tsurutani and Lakhina,
651 1997). The chorus then interacts with the ~ 100 keV injected electrons to energize them to ~ 0.6
652 MeV energies (Inan et al., 1978; Horne and Thorne, 1998; Thorne et al., 2005, 2013; Summers et
653 al., 2007; Tsurutani et al., 2010; Reeves et al., 2013; Boyd et al., 2014). The lower-frequency part
654 of the chorus in turn interact with the ~ 0.6 MeV electrons to accelerate them to ~ 2.0 MeV energies,
655 etc. This bootstrapping mechanism has been suggested by several authors (Baker et al., 1979,

656 1998; Li et al., 2005; Turner and Li, 2008; Boyd et al., 2014, 2016; Reeves et al., 2016) and has
657 been confirmed by Hajra et al. (2015a) during HILDCAA events.

658

659 An alternative scenario is that relativistic electrons are created through particle radial diffusion
660 driven by micropulsations (Elkington et al., 1999, 2003; Hudson et al., 1999; Li et al., 2001,
661 O'Brien et al., 2001; Mann et al., 2004; Miyoshi et al., 2004). However the same general scenario
662 would hold as for chorus acceleration. The substorms and convection events within HILDCAAs
663 would be the sources for the micropulsations and the micropulsations would last from days to
664 weeks in duration. Bootstrapping of energy would still take place.

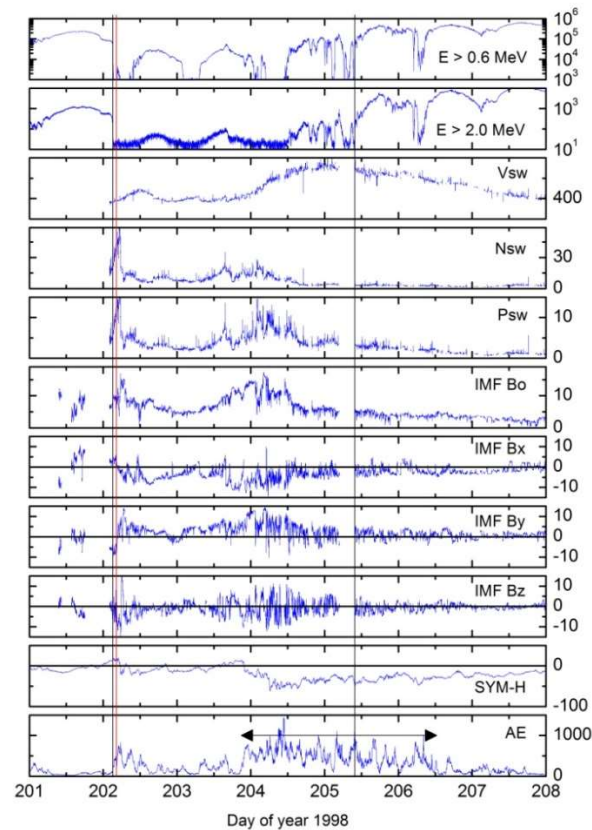
665

666 A few important questions for researchers to ask are: "How high can the relativistic
667 magnetospheric electron energy get?". If there are two HSSs, one from the south pole and another
668 from the north pole so that Earth's magnetosphere is bathed in HSSs for years, as happened during
669 1973-1975 (Sheeley et al., 1976, 1977; Gosling et al. 1976; Tsurutani et al. 1995), will the energies
670 go above ~10 MeV? What will physically limit the energy range? This answer is important for
671 keeping Earth-orbiting satellites safe during such events.

672

673 **3.2.5. Solar wind ram pressure pulses and the loss of relativistic electrons**

674



675
 676 Figure 12. A relativistic electron decrease (RED) event and later acceleration. Taken from
 677 Tsurutani et al. (2016b).

678
 679 Figure 12 shows a relativistic electron decrease (RED) event occurring during 1998. From top to
 680 bottom are the $E > 0.6$ MeV electron fluxes, the $E > 2.0$ MeV electron fluxes, the solar wind speed,
 681 density and ram pressure, and the IMF magnitude, Bx, By and Bz component in the GSM
 682 coordinate system. The bottom two panels are the 1 min SYM-H index (a high time resolution
 683 Dst index) and the AE index. The relativistic electron measurements were taken at $L = 6.6$.

684
 685 At the beginning of day 202, a vertical black line indicates the onset of a high density HPS crossing
 686 (Winterhalter et al., 1994) that is identified in the fourth panel from the top. The HPS is by
 687 definition located adjacent to the HCS (Smith et al. 1978). The HCS is noted by the reversal in the
 688 signs of the IMF Bx and By components (seventh and eighth panels from the top). The onset of
 689 the HPS is followed within one hour by the vertical red line, the sudden disappearance of the $E >$
 690 0.6 MeV (first panel) and $E > 2.0$ MeV (second panel) relativistic electron fluxes. Tsurutani et al.

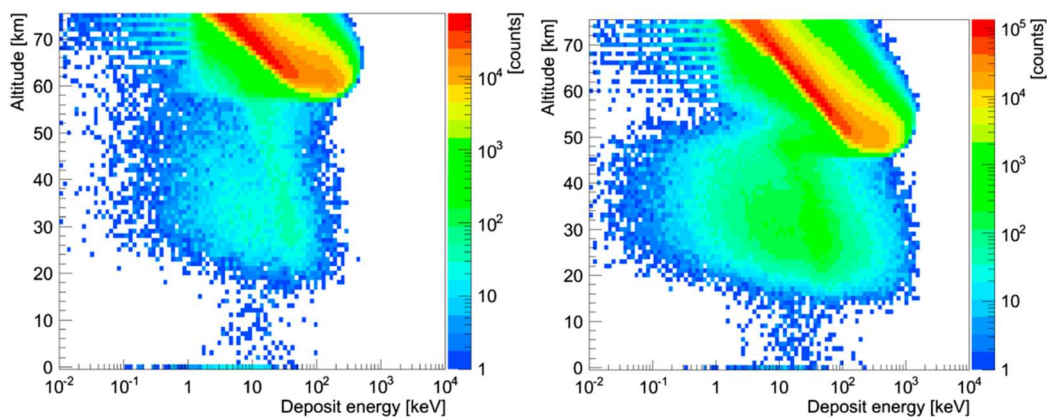
691 (2016b) has shown that for 8 relativistic electron flux disappearance events during solar cycle 23
692 all of the disappearances were associated with HPS impingements onto the magnetosphere.

693

694 Where have the relativistic electrons gone? There are two primary possibilities. One is that the
695 energetic electrons have gradient drifted out of the magnetosphere through the dayside
696 magnetopause, a feature that has been called “magnetopause shadowing” by West et al. (1972).
697 However, a second possible mechanism is electron pitch angle scattering by electromagnetic ion
698 cyclotron (EMIC) waves. We think that this second possibility is more intriguing and has far more
699 interesting consequences, if correct. One might ask where the EMIC waves come from and why is
700 pitch angle scattering particularly important? It has been shown by Remya et al. (2015) that when
701 the magnetosphere is compressed, both electromagnetic chorus (electron) waves (Thorne et al.,
702 1974; Tsurutani and Smith, 1974; Meredith et al. 2002) and EMIC (ion) waves (Cornwall, 1965;
703 Kennel and Petschek, 1966; Olsen and Lee, 1983; Anderson and Hamilton, 1993; Engebretson et
704 al., 2002; Halford et al. 2010; Usanova, 2012; Saikin, 2016) are generated. The compression of
705 the magnetosphere causes betatron acceleration of remnant ~ 10 to 100 keV electrons and protons,
706 and thus plasma instabilities associated with both particle populations occur. What is particularly
707 important is that the EMIC waves are coherent (Remya et al., 2015), leading to extremely rapid
708 pitch angle scattering of ~ 1 MeV electrons by the waves. The scattering rate has been shown to
709 be three orders of magnitude faster than that with incoherent waves (Tsurutani et al., 2016b).

710

711 Another possible loss mechanism is associated with possible generation of PC waves by the HPS
712 impingement followed by radial diffusion of the relativistic electrons. Wygant et al. (1998) and
713 Halford et al. (2015) have mentioned that larger loss cone sizes at lower L could be a source of
714 loss to the ionosphere. Rae et al. (2018) has shown that superposition of compressional PC waves
715 and the conservation of the first two adiabatic invariants could enhance particle losses. However
716 one should mention that there are not observations of PC wave generation during HPS
717 impingements and this needs to be tested. It is also uncertain how rapidly the relativistic electrons
718 would be lost by the above processes. It has been shown that the total loss of $L > 6.6$ relativistic
719 electrons occurs in ~ 1 hour (Tsurutani et al., 2016b).



720
 721 Figure 13. The GEANT4 code run results for the precipitation of $E > 0.6$ MeV electrons (left
 722 panel) and $E > 2.0$ MeV electrons (right panel). The vertical scale is altitude above the ground and
 723 the horizontal scale is energy deposition. The color scheme (legend on the right) gives the amount
 724 of counts. Taken from Tsurutani et al. (2016b).

725
 726 Why can the loss of relativistic electrons to the atmosphere be important? Figure 13 shows the
 727 results of the GEometry ANd Tracking 4 (GEANT4) code developed by the European
 728 Organization for Nuclear Research (Agostinelli et al., 2003) applied to the relativistic electron
 729 disappearance problem. The GEANT4 code takes into account Rayleigh scattering, Compton
 730 scattering, photon absorption, gamma ray pair production, multiple scattering, ionization,
 731 bremsstrahlung for electrons and positrons and annihilation of positrons (positron formation is not
 732 germane for these “low energy” relativistic particles, but the code includes it anyway). A standard
 733 atmosphere was used.

734
 735 Figure 13 shows the GEANT4 Monte Carlo results for the electron shower for $E > 0.6$ MeV
 736 electrons on the left and for $E > 2.0$ MeV electrons on the right. Two important features should be
 737 noticed. First the bulk of energy deposition (the red areas) descends down to ~ 60 km for the $E >$
 738 0.6 MeV electron simulation and down to ~ 50 km for the $E > 2.0$ MeV electron simulation. This
 739 portion of the energy from the incident electrons is due to direct ionization and particle energy
 740 cascading. However, there is a second region which might be extremely important. That is the
 741 blue-green area that goes down to ~ 20 km for the $E > 0.6$ MeV simulation and ~ 16 km for the $E >$
 742 2.0 MeV simulation. There are also “hits” seen on the ground. This lower altitude energy
 743 deposition is due to the relativistic electrons interacting with atmospheric atomic and molecular

744 nuclei creating bremsstrahlung X-rays and γ -rays. X-rays and γ -rays have very large mean free
745 paths and thus can freely propagate through the dense atmosphere without interactions. They
746 propagate to much lower altitudes where they interact and continue the energy cascading process
747 further.

748
749 The reason why this process may be quite an important Space Weather topic is that it might relate
750 to atmospheric weather as well. Wilcox et al. (1973) discovered a correlation between
751 interplanetary HCS crossings and high atmospheric vorticity winds at 300 mb altitude. Over the
752 years a number of different explanations for the physics of the trigger has been offered (Tinsley
753 and Deen, 1991; Lam et al., 2013). Tsurutani et al. (2016b) presented the above relativistic
754 electron precipitation scenario (instead of HCS crossings) for the possible triggers of high
755 atmospheric vorticity winds. Quantitative estimates of potential energy deposition at different
756 atmospheric altitudes were provided in the original paper.

757
758 It is noted that the energy deposition should occur in a limited spatial region of the globe (just
759 inside the auroral zone and a small region of the dayside atmosphere) which is more geoeffective
760 than either cosmic ray energy or solar flare particle deposition. The fact that it is relativistic
761 electron precipitation gives an additional advantage that substantial energy is deposited at quite
762 low altitudes.

763
764 Advances to this problem can be made in a number of different ways. Simultaneous ground-
765 detected EMIC waves, γ -rays and atmospheric heating/cooling could be sought. Correlation with
766 such events with solar wind pressure pulses like the HPSs or interplanetary shocks (see Hajra and
767 Tsurutani, 2018b) would advance our knowledge of the details of such events.

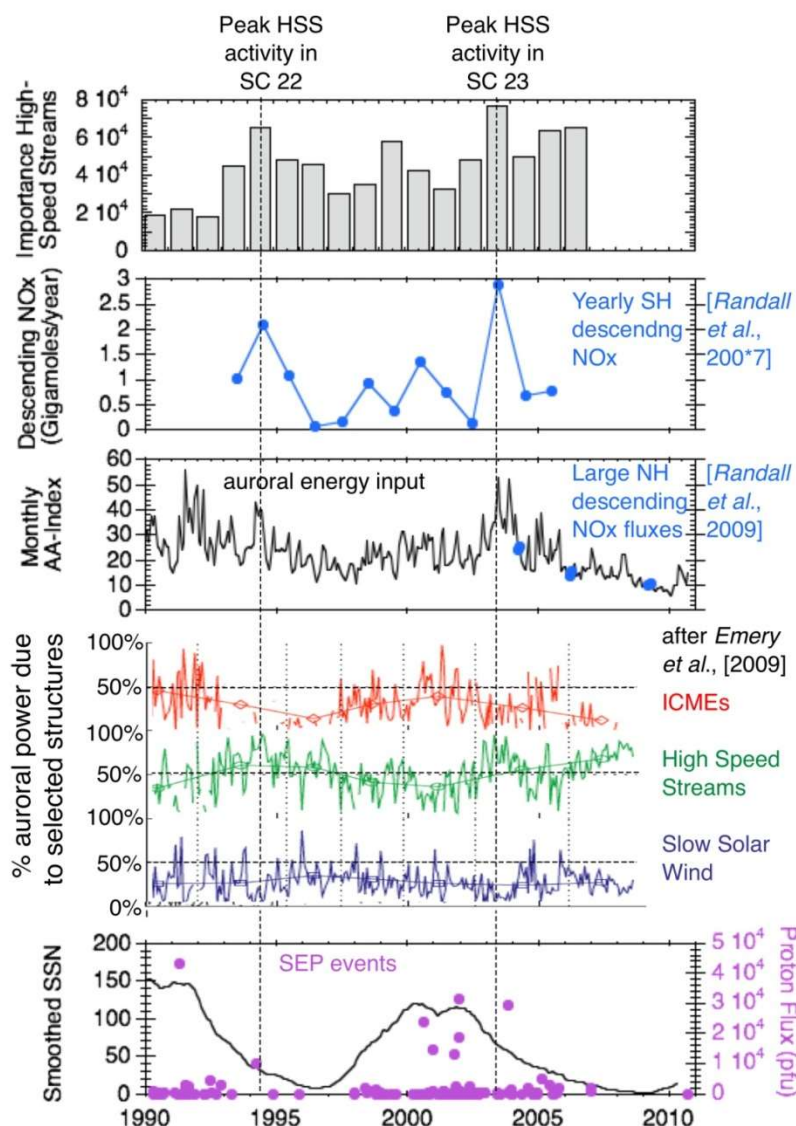
768
769 Maliniemi, Asikainen and Mursula (2014) studied the Earth's winter surface temperatures and the
770 North Atlantic Oscillation (NAO) during all 4 phases of the solar cycle using 13 solar cycles of
771 data (1869-2009). The authors found that the clearest pattern for temperature anomalies is not
772 during sunspot maximum or minimum but during the declining phase when the temperature pattern
773 closely resembles that found during positive NAO. This feature could be due to energetic 10-100
774 keV electron precipitation discussed earlier.

775
776 Atmospheric heating events known as Sudden Stratospheric Warmings (SSWs) (Scherhag, 1960;
777 Harada et al., 2010) occur at subauroral latitudes by unknown causes. They are known to be related
778 to atmospheric wind system changes, perhaps the same phenomenon as the Wilcox et al. (1973)
779 effect. Atmospheric scientists generally assume that SSWs are created by gravity waves
780 propagating from lower atmosphere upward, but so far no one-to-one correlated case has been
781 found. Thus, it would be quite interesting to see if Space Weather can have a major impact on
782 atmospheric weather. The connection between these two disciplines could be quite interesting for
783 the next generation of Space Weather scientists.

784

785 **3.2.6. Energetic particle precipitation and ozone depletion**

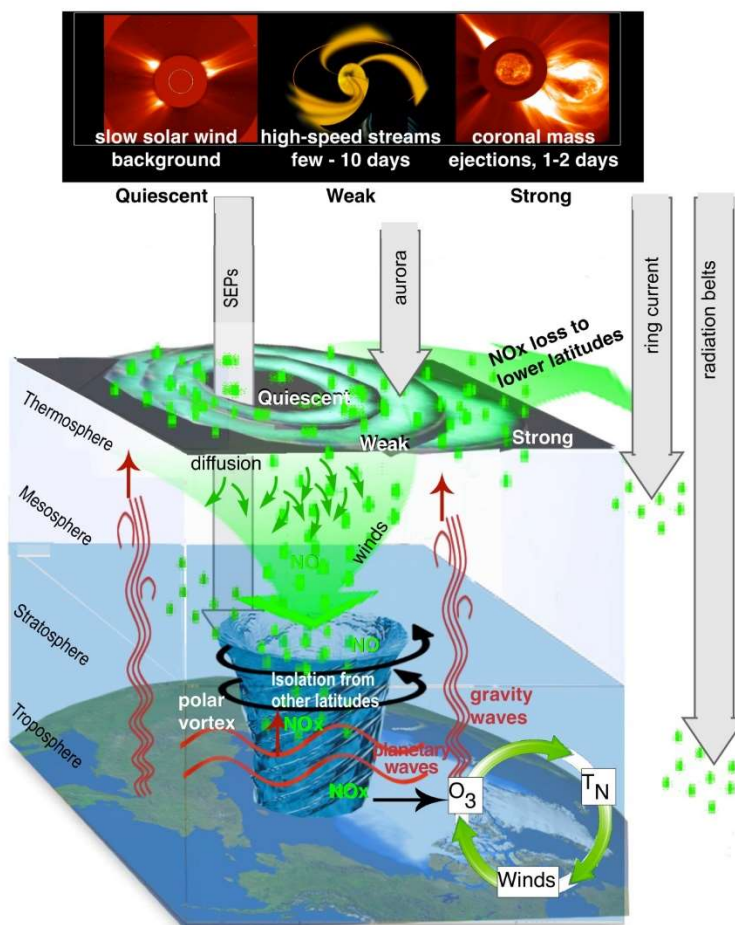
786



787
 788 Figure 14. The dashed vertical lines show the peaks in solar wind high speed streams during SC
 789 22 and SC23. These are coincident with the peaks in auroral energy input and the peaks in yearly
 790 NOx descent. The authors thank J.U. Kozyra for providing this unpublished figure.

791
 792 Figure 14 shows two solar cycles of data, SC22 and SC23. From top to bottom are the
 793 “importance” of high-speed streams, the descending NOx, the monthly AA index, the percent
 794 auroral power due to three types of solar wind phenomena (ICMEs, HSSs and slow solar wind),
 795 and the bottom panel solid line trace is the sunspot number (SSN). Also shown in the bottom panel
 796 is the solar energetic particle (SEP) flux.

797
 798 There are two vertical dashed lines. They correspond to the peaks in HSS activity for SC22 and
 799 SC23 (top panel), peaks in auroral energy input (third panel from the top), and peaks in the yearly
 800 descending NO_x (second panel from the top). It is noted that all three peaks are aligned in time.
 801 The bottom panel shows that both dashed vertical lines correspond to times in the descending
 802 phase of the solar cycle.
 803



804
 805 Figure 15. The scenario for polar cap ozone destruction using the observations shown in Figure
 806 14. The authors thank J.U. Kozyra and her colleagues (personal communication, 2019) for this
 807 unpublished figure.
 808
 809 Figure 15 shows the Kozyra et al. (2019) scenario for ozone destruction over the polar cap. The
 810 top of the Figure shows the various types of solar wind (and associated energetic particles) that

811 can affect atmospheric ozone. The quiet solar wind will lead to quiescence. HSSs lasting a few to
812 ten days have weak effects and ICMEs (and of course shock acceleration of energy particles) can
813 have much stronger effects.

814
815 Energetic particles from different sources will precipitate in different regions of the ionosphere.
816 The energetic particles associated with interplanetary CME shock acceleration will be deposited
817 in the polar regions of the both the north and south ionospheres. If the particles are energetic
818 enough with sufficient gyroradii, they can reach to as low latitudes as $\sim 50^\circ$ magnetic latitude.
819 Precipitating substorm/HILDCAA ~ 10 - 100 keV magnetospheric charged particles will deposit
820 their energy on closed auroral zone ($\sim 60^\circ$ to 70°) magnetic field lines.

821
822 The energetic particle entering the atmosphere lose a portion of their energy in the dissociation of
823 N^2 into $N + N$. The nitrogen atoms will attach to oxygen atoms to form NO_x . Auroral HILDCAA
824 ~ 10 - 100 keV energy particles will only penetrate to depths of ~ 75 km above the surface of the
825 Earth. Solar energetic particles with greater kinetic energies can penetrate lower into the
826 atmosphere to ~ 50 to 60 km. If there is a polar vortex, this vortex can “entrain” the NO_x molecules
827 and atmospheric diffusion can bring them down to lower altitudes over months time duration. The
828 NO_x can act as a catalyst in the destruction of ozone.

829
830 One interesting consequence of extreme ICME shocks is that one would expect extreme Mach
831 numbers to lead to both extreme SEP fluences and also extremely high energies. The former will
832 lead to greater production of NO_x at the polar regions and the latter to deeper penetration and thus
833 less loss of NO_x as they diffuse downward. Alternatively there is a scenario where radiation belt
834 “killer” relativistic electrons can play an important role. If there are large solar polar coronal holes
835 like in 1973-1975, HSSs could produce extremely intense and energetic relativistic electrons.
836 Shocks and HPS impingements on the magnetosphere could cause loss of the electrons to the lower
837 atmosphere. This magnetospheric energy pumping and dumping may have important
838 consequences for NO_x production. The topic of shock acceleration of energetic particles will be
839 discussed in more details in Section 4.1.

840

841

4. RESULTS: Interplanetary Shocks

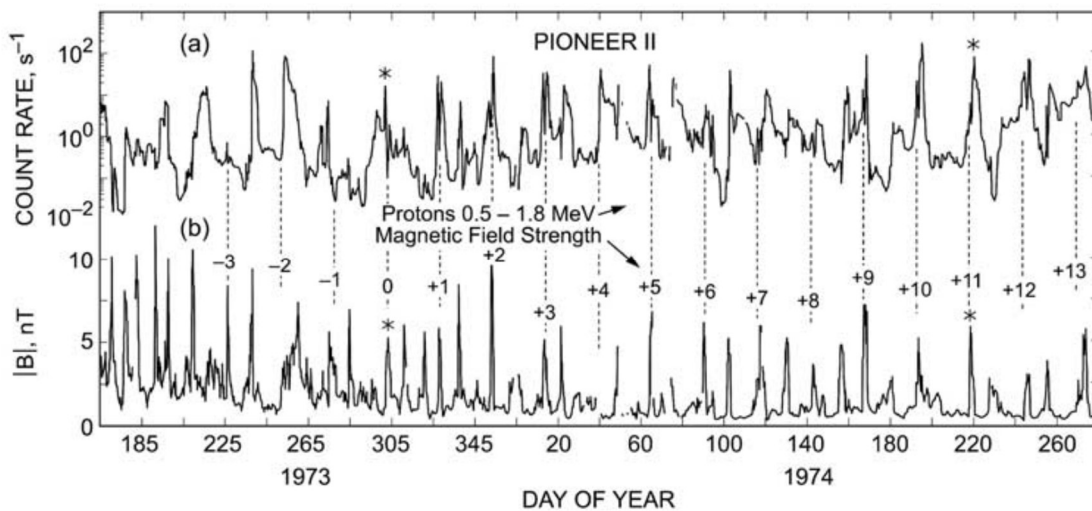
842

843 **4.1. Interplanetary Shocks and Energetic Charged Particle Acceleration**

844

845 Interplanetary shocks have a variety of effects both in interplanetary space and to the Earth's
 846 magnetosphere. It is important for the reader to note that these Space Weather phenomena can
 847 occur with or without the occurrence of magnetic storms. Shock and magnetic storm intensities
 848 are related but only in a loose sense. The physical mechanism for energy transfer for different
 849 phenomena is different. As one example, interplanetary shock acceleration of energetic charged
 850 particles (called "solar cosmic rays") are due to an ICME ram energy driving the fast shocks which
 851 then transfers energy to the charged particles. Solar cosmic ray events can occur with or without
 852 magnetic storms (Halford et al. 2015, 2016; Mays et al., 2015; Foster et al. 2015). Some of the
 853 major extreme Space Weather topics will be addressed below.

854



855

856 Figure 16. Energetic ~ 0.5 to 1.8 MeV protons accelerated by interplanetary fast forward and fast
 857 reverse shocks. Taken from Tsurutani et al. (1982).

858

859 Acceleration of energetic particles in deep space was discovered by Pioneer 11 energetic particle
 860 scientists (McDonald et al., 1976; Barnes and Simpson, 1976; Pesses et al., 1978, 1979; Van
 861 Hollebeke et al., 1978; Christon and Simpson, 1979). As the Pioneer 11 spacecraft traveled away
 862 from the Sun, it was found that the particle fluences kept increasing, contrary to the concept of
 863 adiabatic deceleration. The interplanetary magnetic field magnitude decreases with increasing
 864 distance from the Sun, so one would expect energetic particle deceleration with distance. Thus it

865 was clear to scientists that something must be accelerating these particles in the interplanetary
866 medium. Figure 16 shows one channel of the Pioneer 11 energetic proton count rate, ~ 0.5 to 1.8
867 MeV (see Simpson et al., 1974). The bottom panel is the Pioneer 11 magnetic field (Smith et al.,
868 1975). Some of the peak magnetic fields are numbered, corresponding to a ~ 25 day recurrence of
869 these magnetic structures. The magnetic magnitude structures are identified as well-developed
870 CIRs (see Smith and Wolfe, 1976), bounded by fast forward and fast reverse shocks.

871
872 Tsurutani et al. (1982) identified the shocks and showed statistically that both forward and reverse
873 shocks were related to proton peak count rates. One of the results, which still remains to be solved,
874 is that the proton peaks were generally higher at the reverse shocks. What is the mechanism for
875 greater particle acceleration at fast reverse shocks? This has received little attention and should be
876 addressed in the future.

877
878 Reames (1999) has argued that fast forward shocks upstream (anti-solarward) of ICMEs are the
879 most important phenomenon for the acceleration of “solar flare” particle events. Particle
880 acceleration occurs throughout interplanetary space from near the Sun (where the shocks first
881 form) to 1 AU and beyond as the shocks propagate through the heliosphere. Studies of this
882 acceleration as a function of longitudinal distance away from magnetic connection to the flare site
883 (this gives the variations in the shock normal angle and thus dominant mechanism for acceleration
884 – see Lee (2017) and references therein) have been done by Lario (2012). The features of the
885 energetic particles in space have different characteristics depending on these distances and the
886 portion and characteristics of the shock that the particles are being accelerated from.

887
888 Forecasting the solar flare/interplanetary shock features such as the fluence, energy, spectra and
889 composition will require knowledge of the upstream seed population, upstream (and downstream)
890 waves, and shock properties such as the magnetosonic Mach number and shock normal angle.
891 This is a very difficult task since knowledge of the entire slow solar wind plasma from the Sun to
892 1 AU will be required for accurate forecasting. But again, the Parker Solar Probe and Solar Orbiter
893 may help in developing two points of measurements for modeling of specific events.

894

895 A more fundamental problem is why measured interplanetary fast forward shock Mach numbers
896 at 1 AU are so low? As previously mentioned, Tsurutani and Lin (1985) from ISEE-3
897 measurements have found that at 1 AU, the measured magnetosonic Mach numbers were typically
898 only 1 to 3. Tsurutani et al (2014) have identified a shock with Mach number ~ 9 and Riley et al.
899 (2016) has identified an event with magnetosonic Mach number ~ 28 . The latter event was
900 associated with the SOHO 2012 extreme ICME which did not impact the Earth's magnetosphere.
901 The above are extreme events and little or no events have been detected with intermediate values.
902 A study that is needed is to determine shock Mach numbers at different distances from the Sun.
903 These will give clues as to why 1 AU shock Mach numbers are so low. Is the acceleration of
904 energetic particles causing the dissipation of shock energy as they propagate from the Sun to 1
905 AU? Data from Parker Solar Probe, Solar Orbiter and ACE could be useful in this regard.

906

907 In a related issue, the use of STEREO imaging and MHD modeling could be useful to determine
908 the mass loading of ICME sheaths in causing the deceleration of the ICMEs. This deceleration
909 will also lower the Mach number of the shocks.

910

911 **4.2. Extreme Interplanetary Shocks and Extreme Interplanetary Energetic Particle** 912 **Acceleration**

913

914 Tsurutani and Lakhina (2014) have shown from simple calculations that for CMEs have extreme
915 speeds of 3,000 km/s (Yashiro et al., 2004; Gopalswamy, 2011), shock Mach numbers of ~ 45 are
916 possible. These Mach numbers are getting close to expected supernova shock values. Why haven't
917 such strong shocks been observed at 1 AU? If such events are possible, what would the energetic
918 particle fluences be? Experts on shock particle acceleration will hopefully answer this complex
919 question. It is well known that such solar flare particles enter the polar regions of the Earth's
920 atmosphere and cause radio blackouts. Will extreme solar flare particle fluence precipitation cause
921 different ionospheric effects other than those known today? This latter question might be addressed
922 by ionospheric modelers.

923

924 It should be noted that although Space Weather is a chain of events/phenomena going from the
925 Sun to interplanetary space to the magnetosphere, ionosphere and atmosphere, there is often not a

926 direct link between different facets of Space Weather. Each feature of Space Weather should be
927 examined separately and it should not be assumed that an extreme flare will cause extreme
928 cascading Space Weather phenomena. We use solar flare particles as an example for the reader.
929 The largest solar flare particle event in the space age occurred in August 1972 (Dryer et al., 1976
930 and references therein). However, there was no magnetic storm caused by the MC impact onto
931 the Earth's magnetosphere (the MC field was directed almost entirely northward, leading to
932 geomagnetic quiet: Tsurutani et al. 1992b). On the other hand, the largest magnetic storm on
933 record is the "Carrington" storm. The storm intensity will be discussed further in Section 7.0.
934 There is little or no evidence of large solar flare particle fluences in Greenland ice core data from
935 that event (Wolff et al., 2012; Schrijver et al., 2012). Usoskin and Kovfaltsov (2012) examining
936 historical proxy data (^{14}C and ^{10}Be) also find a lack of any signature associated with the Carrington
937 flare. Although this is an extreme example, it is useful to mention it to illustrate the point: different
938 facets of Space Weather may have only loose correlations with other facets.

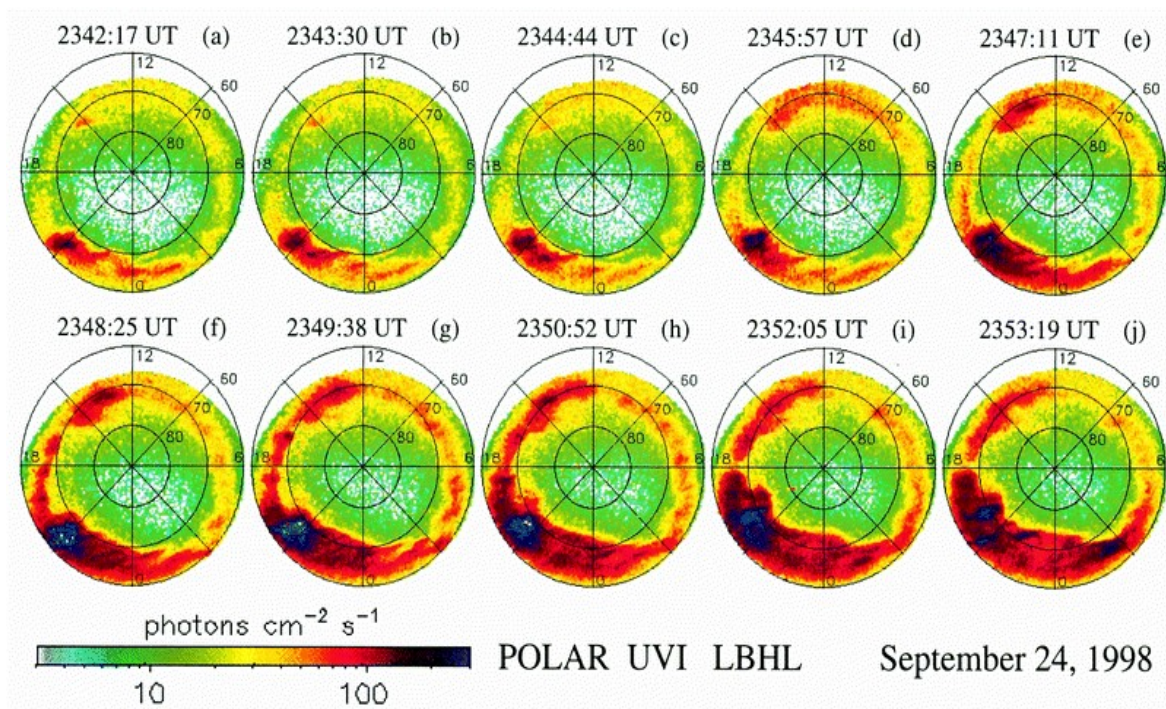
939

940 An area that has received a lot of attention lately is ancient solar flares. Miyake et al. (2012)
941 discovered an anomalous 12% rapid increase in ^{14}C content from 774 to 775 AD in Japanese cedar
942 tree rings. Usoskin et al. (2013) have argued that such an extreme radiation event could be
943 associated with an extreme solar energetic particle event (or a sequence of events). The latter
944 authors estimated that the fluence of > 30 MeV particles was $\sim 4.5 \times 10^{10} \text{ cm}^{-2}$. Could such an
945 extreme particle event be associated with an extremely strong interplanetary shock or instead series
946 of strong shocks? Space Weather scientists are currently working on this problem.

947

948 **4.3. Interplanetary shocks, dayside aurora and nightside substorms**

949



950
 951 Figure 17. Interplanetary shocks cause dayside auroras and trigger nightside substorms. The
 952 images show the northern polar views of polar cap and auroral zones taken in UV wavelengths.
 953 Local noon is at the top in each image. The Figure is taken from Zhou and Tsurutani (2001).

954
 955 Interplanetary shocks can trigger the precipitation of energetic ~ 10 to 100 keV electrons into the
 956 auroral ionosphere (Halford et al. 2015). In fact, low energy ($E < 10$ keV) electron precipitation
 957 can occur as well. Figure 17 shows interplanetary shock impingement auroral UV effects for an
 958 event on September 23, 1998. Each image has the north pole at the center and 60° magnetic
 959 latitude (MLAT) shown at the outer edge. Noon is at the top and dawn is at the right. The cadence
 960 between images is ~ 1 min 13 s. From ACE measurements and propagation calculations it is known
 961 that the fast forward shock arrived the magnetosphere between the images c), 2344:44 UT and d),
 962 2345:47 UT. What is apparent in panel d) is the sudden appearance of aurora on the dayside (Zhou
 963 and Tsurutani et al., 1999). From further analyses of these shock auroral events, Zhou et al. (2003)
 964 have shown that magnetospheric compression of preexisting ~ 10 to 100 keV electrons and protons
 965 will generate both electromagnetic electron and proton plasma waves and diffuse auroras (as
 966 discussed previously). Also noted were the generation of field-aligned dayside currents.
 967 Compression of the magnetosphere will generate Alfvén waves (Haerendel, 1994) which will

968 propagate along the magnetic field lines down to the ionosphere. Wave damping could provide
969 substantial ionospheric heating.

970
971 The mechanism for energy transfer from the solar wind to the magnetosphere is the absorption of
972 the solar wind ram energy. Dayside auroras occur with shock impingement irrespective of the
973 interplanetary magnetic field Bz direction. Another possible mechanism for the dayside aurora
974 not mentioned above are double layers above the ionosphere (Carlson et al., 1998) with the
975 acceleration of ~1 to 10 keV electrons and the formation of discrete dayside auroras. What is the
976 relative importance of these three different auroral energy mechanisms? This would be an
977 excellent topic for the SWARM and Arase satellite missions. Coordinated ground measurements
978 would be useful.

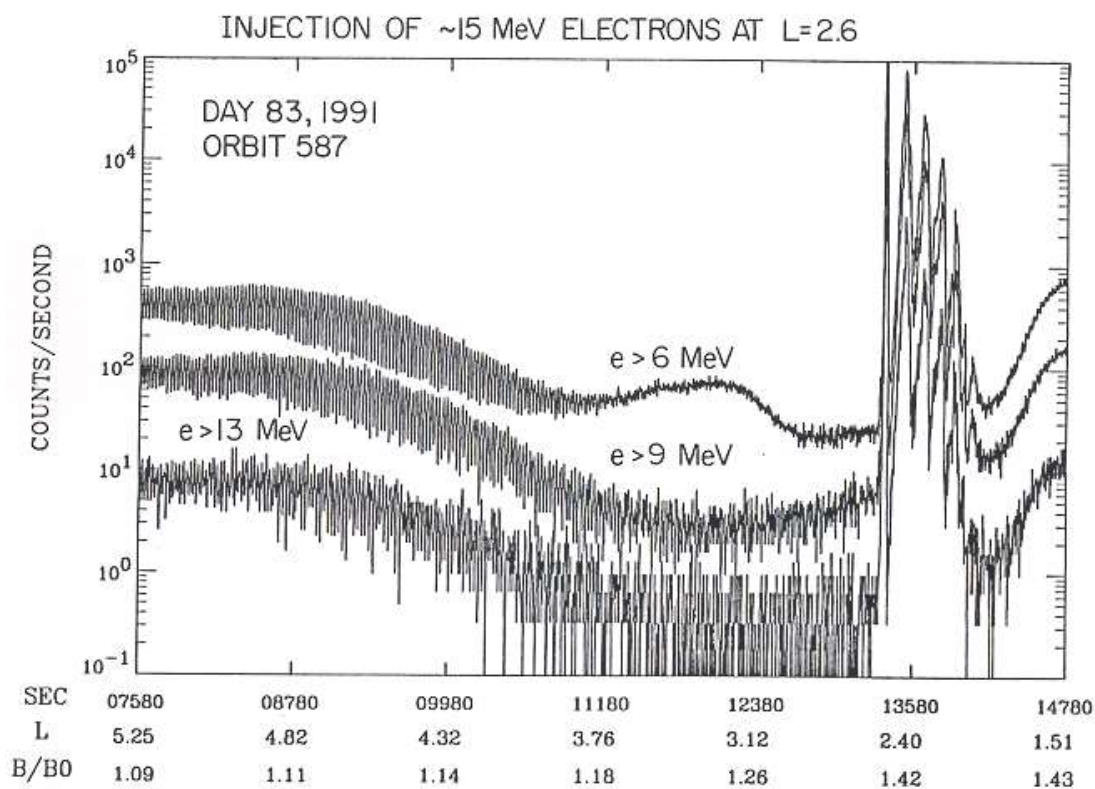
979
980 Returning back to Figure 17 panel e) 2347:11UT, there is a substorm intensification centered at
981 ~2100 magnetic local time (MLT). The substorm further intensification and expansion can be
982 noted in the sequence of images. Interplanetary shock triggering of substorms has been known to
983 occur before the advent of imaging polar orbiting spacecraft (Heppner, 1955; Akasofu and Chao,
984 1980). The AE index had been used to identify these events.

985
986 An important fundamental question for substorm physics that has existed for a long time, is where
987 in the tail/magnetosphere does the substorm get initiated and by what physical mechanism? Is it
988 reconnection or plasma instabilities (Akasofu, 1972; Hones, 1979; Lui et al., 1991; Lui, 1996;
989 Baker et al., 1996; Lakhina, 2000)? Where does the energy come from, recent precursor solar
990 wind inputs as suggested by Zhou and Tsurutani (1999), or stored tail energy or even possibly
991 solar wind ram energy (see Hajra and Tsurutani, 2018b)? The rapid response of the magnetosphere
992 to the shock should limit the downstream location of the substorm initiation point. It should be
993 noted that there are probably several different mechanisms for causing substorms. Although this
994 is only the shock triggering case, knowledge of this may help understand other cases, if they are
995 indeed different. The MMS mission will be ideally suited for addressing this question in the tail
996 phase of the mission.

997

998 **4.4. Interplanetary shocks and the formation of new radiation belts**

999



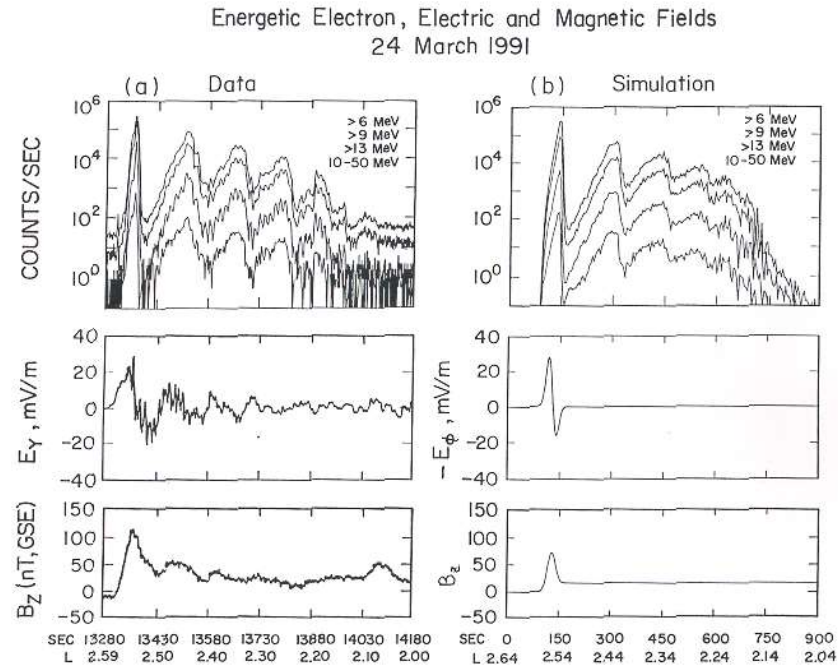
1000
1001 Figure 18. Shock creation of a new relativistic electron radiation belt in the magnetosphere. The
1002 three energy channel plots show an abrupt increase in flux at the same time. Recurrence of the flux
1003 with decreasing amplitude occurs at least 4 more times. Figure taken from Blake et al. (1992).

1004
1005 Figure 18 shows evidence of a new “radiation belt” triggered by a strong interplanetary shock. The
1006 Figure shows three traces, $E > 6$ MeV, > 9 MeV and > 13 MeV fluences. At the time of the strong
1007 and sudden increase in all energy fluxes, the spacecraft was at $L = 2.6$. This is time-coincident
1008 with the shock impingement upon the magnetosphere (not shown). With increasing time, a second,
1009 then third, etc., electron flux pulse appears. These are “drift echoes” where the energetic electron
1010 “cloud” has gradient drifted around the magnetosphere to return to the satellite location once again.

1011

1012 **4.4.1. What is the mechanism to create this new radiation belt?**

1013



1014

1015 Figure 19. An expanded version of the relativistic electron pulse and measured magnetospheric
 1016 electric field and magnetic field B_z on the left and simulation results on the right. Taken from Li
 1017 et al. (1993).

1018

1019 The left-hand column of Figure 19 shows an expanded version of Figure 16 on the top with the
 1020 addition of the ~ 10 to 50 MeV count rate channel included. Next is the d.c. electric field in the Y
 1021 direction, and magnetospheric B_z on the bottom. The right-hand column bottom shows a magnetic
 1022 pulse input into the system. This generates a time varying azimuthal electric field (right middle)
 1023 and the relativistic electron flux at the top right.

1024

1025 Using the input of a single magnetospheric magnetic pulse into the magnetosphere, Li et al. (1993)
 1026 simulated the acceleration and injection of $E > 40$ MeV electrons. What is interesting is that the
 1027 origin of the electrons was $L > 6$ with energies of only a few MeV. The reader should read Li et
 1028 al. (1993) for more details concerning the simulation and results. Related works on acceleration of
 1029 magnetospheric electrons by shock impact on the magnetosphere can be found in Wygant et al.
 1030 (1994), Kellerman and Shprits, 2012; Kellerman et al., 2014; Foster et al. (2015).

1031

1032 How strong was the interplanetary shock? There was not any spacecraft upstream of the Earth at
 1033 the time of the event, so no measurements of shock strength can be made. However, Araki (2014)
 1034 has noted that this shock caused a SI^+ of magnitude 202 nT. This is the second largest SI^+ in
 1035 recorded history. In Tsurutani and Lakhina (2014) with the assumption of a 3,000 km/s CME and
 1036 only a 10% deceleration from the Sun to 1 AU, they estimated a maximum SI^+ of 234 nT under
 1037 normal conditions. Could this 1991 shock strength have been close to the $M = 45$ estimate
 1038 mentioned earlier? One cannot really tell for sure because the shock Mach number strongly
 1039 depends on the upstream plasma conditions, which can only be estimated in this case.

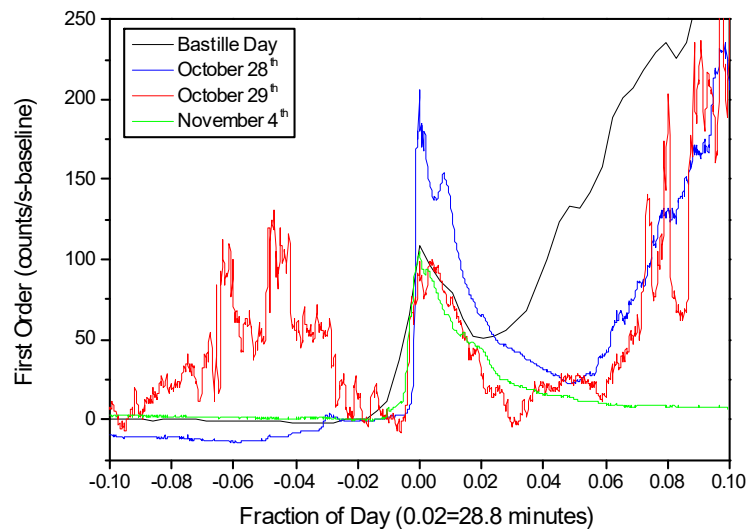
1040

1041 Tsurutani and Lakhina (2014) estimated a dB/dt six times larger than the one used in the Li et al.
 1042 (1993) modeling. What would a maximum dB/dt cause in a new radiation belt formation? How
 1043 much greater could the relativistic electron energy and flux become?

1044

1045 5. RESULTS: Solar Flares and Ionospheric Total Electron Content

1046



1047

1048 Figure 20. The largest solar EUV flare in recorded history, October 28, 2003. Taken from
 1049 Tsurutani et al. (2005b).

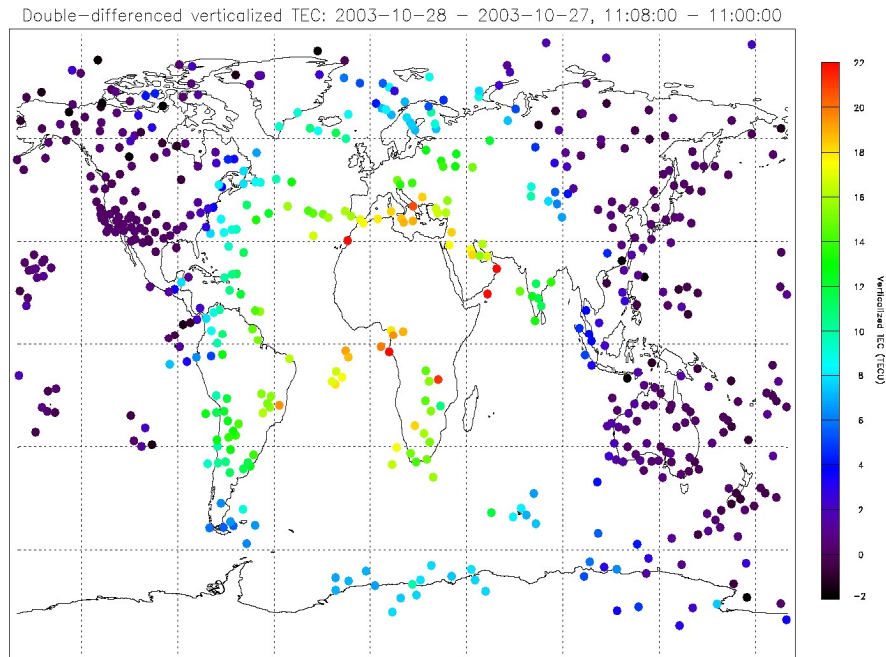
1050

1051 Figure 20 shows four well-known solar X-ray flare events taken in a narrow band 26-34 nm EUV
1052 spectrum. The four flare events are the Bastille day (July 14, 2000) flare and three “Halloween”
1053 flares occurring on October 28, 29 and November 4, 2003. The narrow band EUV spectrum is
1054 shown because some of the flare X-ray and EUV fluxes were so intense that most spacecraft
1055 detectors became saturated (all except the SOHO SEM narrowband EUV detector). The X-ray
1056 flare intensities could only be estimated from fitting techniques for the saturated data. Here we
1057 use the narrow band channel of the SOHO SEM detector where the four above mentioned flares
1058 were not saturated. The four flare count rate profiles were aligned so that they start at time zero.
1059 What is particularly remarkable is that the October 28, 2003 flare has the highest EUV peak
1060 intensity of all four events and was greater by a factor of ~ 2 . This is the most intense EUV solar
1061 flare in recorded history.

1062
1063 After each flare reached a peak intensity and then decreased in count rate, there was often a
1064 following increase in count rate. This is particularly notable in the Bastille day (black trace) flare.
1065 This increase is contamination due to delayed energetic electrons propagating through space along
1066 interplanetary magnetic field lines reaching the spacecraft later in time. The November 4 flare
1067 (green) did not have such contamination because it was a limb flare and presumably (magnetic)
1068 connection from the flare site to the spacecraft did not occur.

1069
1070 NOAA personnel have estimated the November 4 flare had an intensity of $\sim X28$. This event
1071 saturated the detector so this is a conservative estimate. Thomson et al. (2004) using a different
1072 technique estimated a value of $X45$ for this event. NOAA has estimated that the October 28 flare
1073 as $\sim X17$. However, in EUV fluxes, the October 28 flare was the most intense by far.

1074



1075

1076 Figure 21. The global TEC during the October 28, 2003 solar flare. The scale is given on the right.

1077 The figure is taken from Tsurutani et al. (2005b).

1078

1079 Figure 21 shows the global total electron content (TEC) in the ionosphere after the October 28,

1080 2003 solar flare. The map has been adjusted so Africa, the subsolar point, is in the center of the

1081 Figure. The top and bottom of the plot correspond to the Earth's polar regions and the left side

1082 and right-side edges local midnight. The enhanced TEC area corresponds to the sunlit hemisphere.

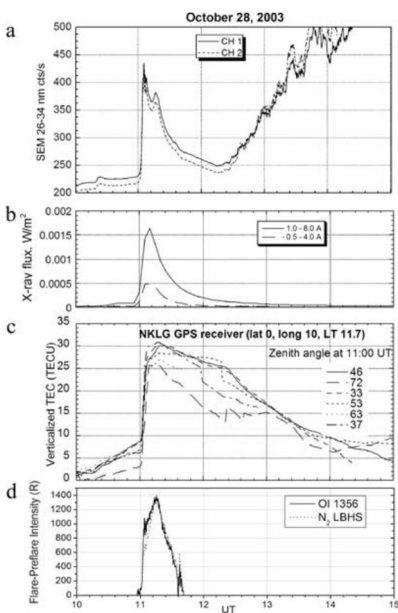
1083 At the subsolar point the TEC enhancement was $\sim 30\%$. This is the record for flare-induced

1084 ionospheric TEC (Tsurutani et al., 2005b). The nightside hemisphere shows no TEC enhancement,

1085 as expected. The TEC enhancement is due to ionization by X-rays, EUV photons and UV photons,

1086 all part of the solar flare spectrum.

1087



1088
 1089 Figure 22. The ionospheric and atmospheric effects of the October 28, 2003 solar flare. Taken
 1090 from Tsurutani et al. (2005b).

1091
 1092 Figure 22 shows the effects of the October 28 solar flare. From top to bottom are the SOHO SEM
 1093 EUV count rate, the GOES X-ray flux, the Libreville, Gabon TEC data and the GUVI O and N²
 1094 dayglow data. It is noted that the flare profiles in EUV and X-rays last ~tens of mins and are similar
 1095 in profile to each other. However, the TEC over Libreville last hours. This is due to the EUV
 1096 portion of the solar flare. These photons deposit their energy at ~170 to 220 km altitude where the
 1097 recombination time scales are ~ 3 to 4 hours. Thus, EUV photon ionization has longer lasting
 1098 ionospheric TEC effects. The X-ray portion of the solar flare spectrum deposit their energy in the
 1099 ~80 to 100 km altitude range where the recombination time scale is tens of min (Thomson et al.,
 1100 2005, and references therein). This solar flare example is one where solar energy (photons) goes
 1101 directly from the Sun to the Earth's ionosphere (previously shown examples such as with ICMEs
 1102 and sheaths with magnetic storms have solar plasma and magnetic field energy transfer from the
 1103 Sun to interplanetary space to the magnetosphere).

1104
 1105 Some future Space Weather problems are to be able to predict the solar flare energy spectrum
 1106 given the underlying solar flare surrounding geometry. We have indicated that the 28 October
 1107 2003 and the 4 November 2003 flares were significantly different spectra-wise. The question is
 1108 why and how often does this happen? Ionospheric satellites like the Constellation Observing

1109 System for Meteorology, Ionosphere and Climate-2 (COSMIC II) and SWARM can probe for
 1110 detailed altitude dependence of ionization to work backwards to attempt to identify what energy
 1111 spectrum would cause the layered ionization detected. Solar flare data taken by instrumentation
 1112 onboard the RHESSI and EVE/SDO spacecraft would be useful to understand the details of flare
 1113 spectral differences. Other questions are how large can X-ray and EUV flares become? What will
 1114 their ionospheric effects be?

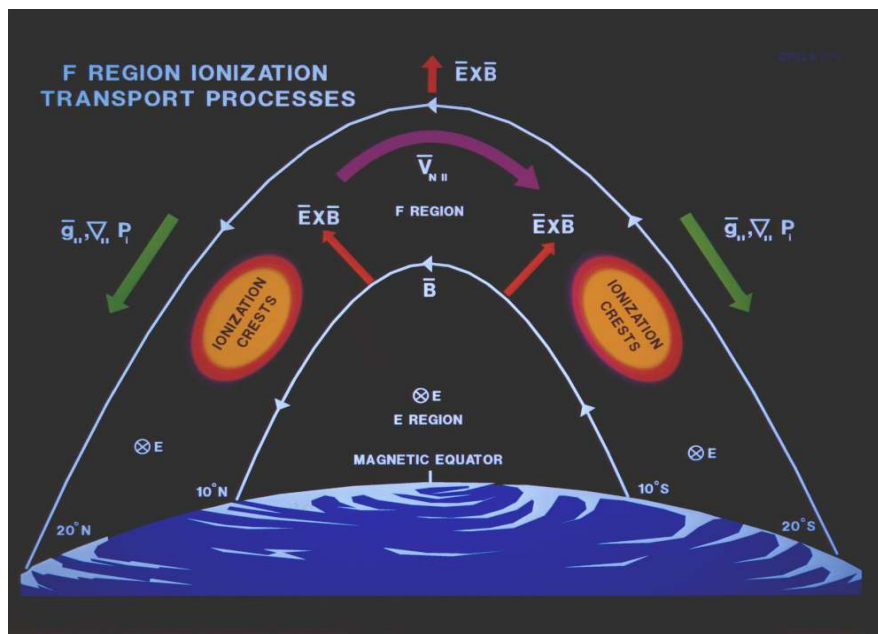
1115

1116 6. RESULTS: Magnetic Storms and Prompt Penetrating Electric Fields (PPEFs)

1117

1118 For substorms, PPEFs occurring in the ionosphere have been known for a long time, since the
 1119 beginning of the space age (Nishida and Jacobs, 1962; Obayashi, 1967; Nishida, 1968; Kelley et
 1120 al. 1979, 2003). In the last 10 years lots of work has been done on PPEFs during magnetic storms.
 1121 Why didn't people look at storms earlier? Because it was theoretically predicted that the PPEFs
 1122 would be shielded out. Why doesn't shielding happen? This is a very good question for workers
 1123 in the field. Right now we don't know the answer.

1124



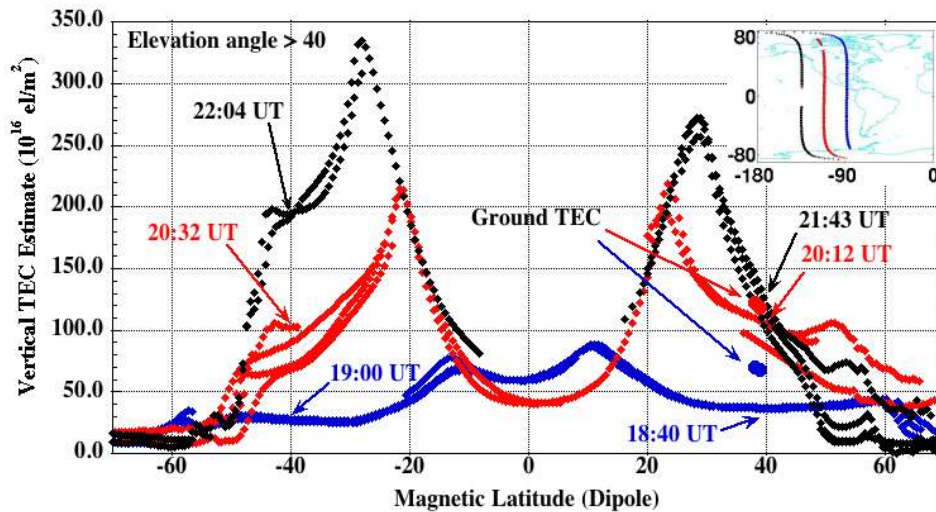
1125

1126 Figure 23. Dayside (near) equatorial ionization anomalies (EIAs) located $\sim \pm 10^\circ$ on both sides of
 1127 the magnetic equator. The local Earth magnetic field is shown in this schematic. The figure is
 1128 taken from Anderson et al. (1996).

1129

1130 Figure 23 show the geometry of the Earth's magnetic field near the magnetic equator. It is parallel
 1131 to the Earth's surface at the equator but where the equatorial ionization anomalies (EIAs) are
 1132 located, the magnetic field is slanted. The EIAs are standardly located at $\sim\pm 10^\circ$ MLAT in the
 1133 dayside ionosphere. With red arrows, the figure also shows the direction of $E \times B$ convection. At
 1134 exactly the magnetic equator, $E \times B$ is in a purely upward direction. At the positions of the EIAs,
 1135 the $E \times B$ direction is both upward and to higher absolute magnetic latitudes.

1136



1137

1138 Figure 24. Three passes of the CHAMP satellite measuring the near equatorial and midlatitude
 1139 TEC during October 30, 2003. CHAMP was at an altitude of ~ 430 km, so the TEC measured was
 1140 the total thermal electron column density above that altitude. The figure is taken from Mannucci
 1141 et al. (2005).

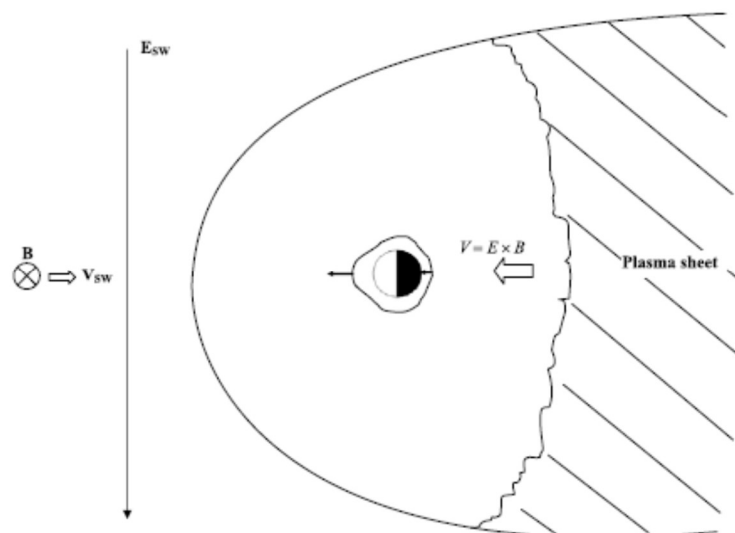
1142

1143 Figure 24 shows three passes of the CHAMP satellite in polar orbit with an altitude of ~ 430 km at
 1144 the near equatorial crossings. The three orbits are given in the upper right-hand portion of the
 1145 figure. The first TEC trace shown in blue is before the onset of the October 30-31 magnetic storm.
 1146 The two EIAs are identified by the TEC enhancements at $\sim\pm 10^\circ$ with peak intensities of ~ 80 TEC
 1147 units. In the next pass (red trace), the EIAs are located at $\sim\pm 21^\circ$ MLAT and the peak intensities
 1148 are ~ 210 TEC units. During the next satellite pass, the EIAs are located near $\pm 30^\circ$ and the TEC

1149 values become as high as ~ 330 TEC units. This “movement” of the EIAs to higher magnetic
 1150 latitudes can be explained by a convective electric field (PPEF) in the east-west direction causing
 1151 an uplift to both EIAs by $E \times B$ convection as explained earlier associated with Figure 23. One
 1152 might ask why does the TEC increase to such high values?

1153
 1154 The answer is as the PPEF removes the plasma from the ionospheric lower F region and brings it
 1155 to higher altitudes where the recombination time scale is longer (hours), the Sun’s EUV photons
 1156 replace the plasma by photoionization of the upper atmosphere, replacing the lost plasma and thus
 1157 increasing the “total electron content” of the ionosphere. This is one cause of a “positive
 1158 ionospheric storm”.

1159



1160
 1161 Figure 25. The interplanetary, magnetospheric and equatorial ionospheric electric fields during a
 1162 PPEF event. The figure is taken from Tsurutani et al. (2004c; 2008b).

1163
 1164 Figure 25 shows the interplanetary motional electric field for southward interplanetary B_z . The
 1165 electric field will be in the dawn-to-dusk direction. When magnetic reconnection takes place in
 1166 the nightside plasmasheet, the convective electric field will be in the same direction but with a
 1167 reduced amplitude. This electric field brings the plasmasheet plasma into the nightside low L
 1168 region magnetosphere during magnetic storms. The PPEFs penetrate into the dayside equatorial
 1169 ionosphere (shown in Figure 24) and also the nightside equatorial ionosphere. However
 1170 significantly different from the dayside case, the $E \times B$ convection on the nightside will bring the

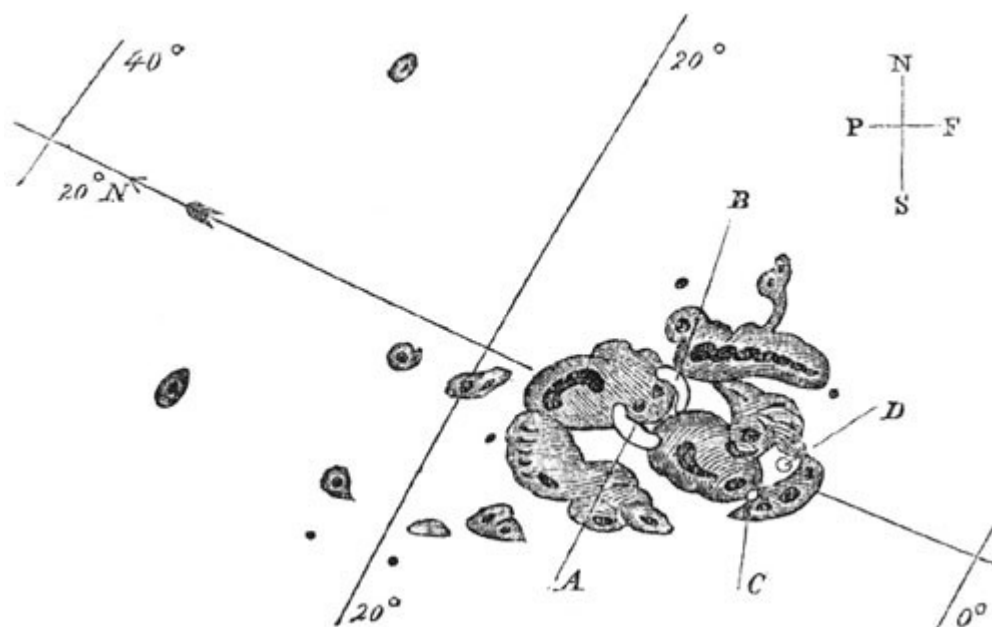
1171 ionospheric plasma to lower altitudes, leading to recombination and reduction in TEC. This is one
 1172 form of a “negative ionospheric storm”. See Mannucci et al. (2005, 2008) for discussions of
 1173 positive and negative ionospheric storms.

1174
 1175 There are many important questions about PPEFs which are almost always present during major
 1176 magnetic storms. As previously mentioned, “why aren’t the electric fields shielded out?” What is
 1177 the mechanism for generating PPEFs, wave propagation from the polar ionosphere as suggested
 1178 by Kikuchi and Hashimoto (2016) or a more global picture as Figure 25 and Nishida and Jacobs
 1179 (1962) suggest? Figure 25 is a simple schematic. What are the real local time dependences of the
 1180 PPEF? Does this vary from storm to storm, and if so, why? Why does the relative PPEF magnitude
 1181 vary from one storm to the next? Again, future spacecraft and ground-based studies will be able to
 1182 help answer these questions.

1183

1184 7. RESULTS: The Carrington Storm

1185

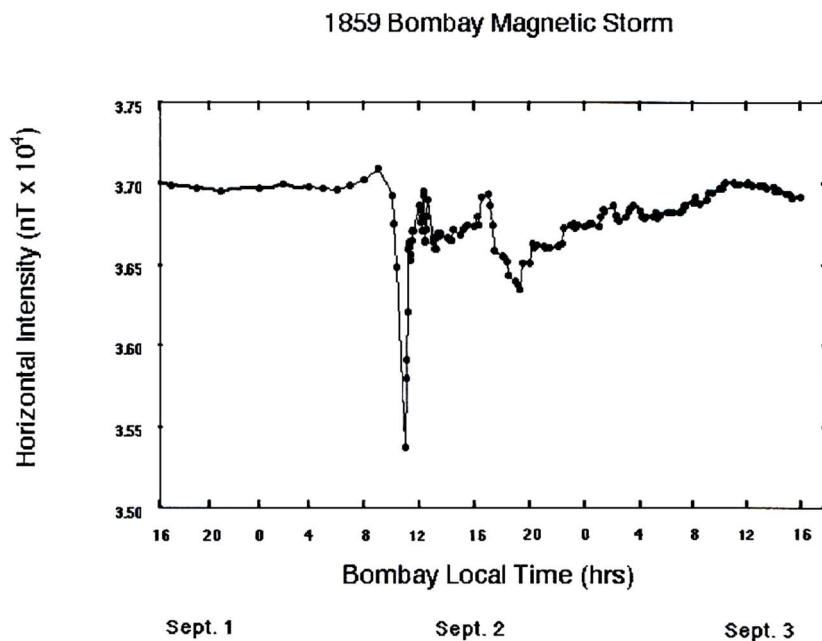


1186
 1187 Figure 26. The solar active region during the Carrington 1 September 1859 optical solar flare. The
 1188 figure is taken from Carrington (1859).

1189
 1190 Figure 26 is the active region (AR) that was hand-drawn by Richard Carrington. This was the
 1191 source of the optical solar flare that he and Hodgson (1859) saw and reported on 1 September

1192 1859. See Cliver (2006) for a nice accounting of the observational activity taken during 1859 flare
 1193 interval and Kimball (1960) for an accounting of the aurora during the storm. The optical part of
 1194 the flare lasted only ~ 5 min. Some ~ 17 hr 40 min later a magnetic storm occurred at Earth
 1195 (Carrington, 1859).

1196



1197

1198 Figure 27. The Carrington storm detected in the Colaba, India magnetometer. The Figure is taken
 1199 from Tsurutani et al. 2003 and Lakhina et al. 2012.

1200

1201 Figure 27 shows the H-component magnetic field taken by the Colaba magnetic observatory during
 1202 the “Carrington” magnetic storm. The SI^+ is estimated to be ~ 110 nT and the magnetic decrease
 1203 ~ 1600 nT at Colaba (Mumbai, India). The SI^+ and storm main phase has been recently shown to
 1204 be most likely caused by an upstream solar wind density of 5 particles cm^{-3} and a MC with intensity
 1205 ~ 90 nT (pointed totally southward) by Tsurutani et al. (2018a). No particularly unusual solar wind
 1206 conditions are believed to have been necessary (in contrast to the original conclusions of Ngwira
 1207 et al., 2014). Ngwira et al. (2018) is now in accord with this more recent assessment of a normal
 1208 upstream solar wind.

1209

1210 The intensity of the “Carrington” storm was estimated as $Dst = -1760$ nT (Tsurutani et al., 2003)
 1211 based on observations of the lowest latitude of red auroras being at $\pm 23^\circ$ (Kimball, 1960). The

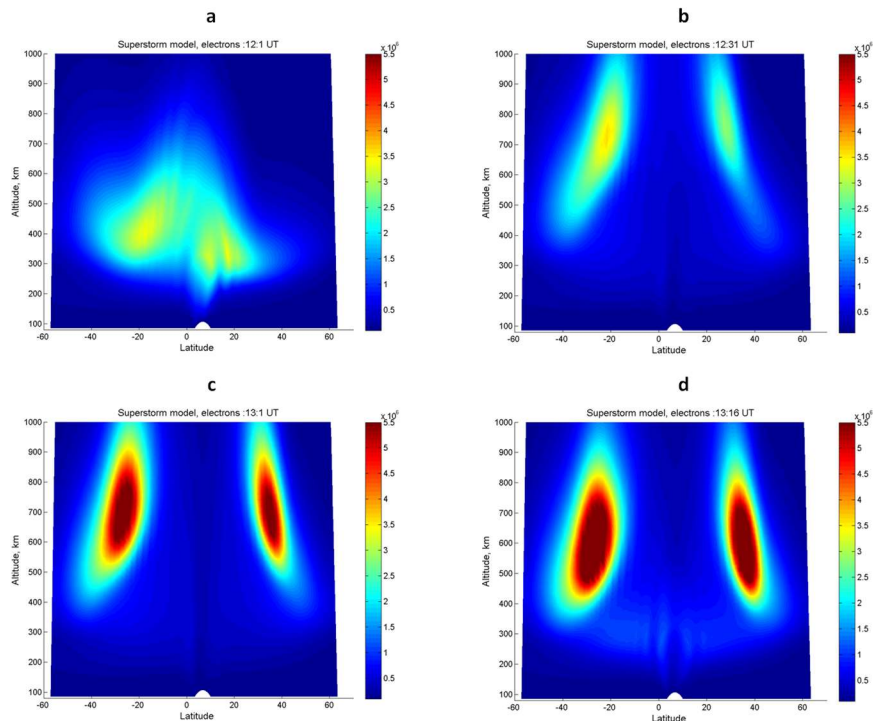
1212 storm intensity was calculated using recent theoretical expressions of magnetospheric potentials
1213 needed to convect plasma into such low latitudes. Siscoe (1979) basing his estimate on a model
1214 that treats the pressure as a constant along the magnetic flux tube came up with a value of Dst = -
1215 2000 nT.

1216
1217 It should be mentioned that some researchers have taken exception with the Colaba magnetogram
1218 as an indication of ring current effects (see Comment by Akasofu and Kamide (2005) and Reply
1219 by Tsurutani et al. (2005a)). The Colaba magnetic profile is unlike those of ICME magnetic storms
1220 discussed in Sections 2.3, 2.4 and 3.1 of this paper. Several researchers have estimated the storm
1221 intensity based on the Colaba magnetogram (see articles in a special journal edited by Clauer and
1222 Siscoe, 2006; Acero et al. 2018). The Colaba data clearly show that the storm had exceptionally
1223 large geomagnetic effects, irregardless of the interpretation of the Colaba data. Possible
1224 interpretations of the Colaba profile will be discussed later in the paper.

1225
1226 The most accurate method of estimating a magnetic storm intensity is by using the latitude of the
1227 aurora. Red auroras (Stable Auroral Red or SAR arcs) are presumably an indication of the location
1228 of the plasmopause (R.M. Thorne, private communication, 2002). Kimball (1960) noted that “red
1229 glows” were detected at $\pm 23^\circ$ from the geomagnetic equator during the Carrington event. In 1960
1230 the term “SAR arc” was not in use, but we can assume that this was what he was reporting. At the
1231 present time, this is the most equatorward SAR arcs that have been observed (thus the most intense
1232 magnetic storm). That is until researchers find records of even lower latitude red auroras!

1233
1234 Comments on the short duration of the recovery phase has been made by Li et al. (2006). A high-
1235 density filament was used to explain this unusual feature of the magnetic storm profile. Tsurutani
1236 et al. (2018a) have recently proposed another possibility. During extreme events when the storm
1237 time convection brings the plasmashet into very low L, all of the standard ring current loss process
1238 rates will be enhanced. There will be greater Coulomb scattering, greater charge exchange loss
1239 rates and greater plasma wave growth with consequential greater wave-particle pitch angle
1240 scattering and losses to the atmosphere. In Tsurutani et al. (2018a) the authors focused particularly
1241 on wave-particle interactions because the size of the loss cone will increase dramatically with
1242 decreasing L. This, plus greater energetic particle compression due to the extreme inward

1243 convection, will lead to stronger loss cone/temperature anisotropy instabilities, greater wave
 1244 growth and thus greater losses. This hypothesis can be easily tested by magnetospheric spacecraft
 1245 observations during large magnetic storms and by magnetospheric modeling perhaps bringing
 1246 some light to the unusual Colaba magnetic signature.
 1247



1248
 1249 Figure 28. A model of the PPEF effects of the Carrington 1859 storm on the dayside ionosphere.
 1250 The input electric field was taken from Tsurutani et al. (2003) and the simulation was performed
 1251 using the Huba et al. (2000, 2002) SAMI2 code. The figure is taken from Tsurutani et al. (2012).

1252

1253 7.1. The Carrington PPEF

1254

1255 One of the concerns for extreme Space Weather in the ionosphere are extremely intense PPEFs
 1256 and the daytime superfountain effect on the uplift of O^+ ions (positive ionospheric storms). Higher
 1257 ion densities in the exosphere will lead to the possibility of enhanced low altitude satellite drag.
 1258 In Tsurutani et al. (2003), the authors used modern theories of the electric magnetospheric potential
 1259 given by Volland (1973), Stern (1975) and Nishida (1978) to determine the electric field during

1260 the Carrington storm main phase. The former authors obtained an estimate of ~ 20 mV/m. They
1261 then applied this electric field in the SAMI2 model with the results shown in Figure 28.

1262
1263 Figure 28 shows the SAMI2 results of the modeled dayside ionosphere with a ~ 20 mV/m added
1264 to the diurnal variation electric field. The quiet ionosphere is shown at the upper left. The uplift
1265 of the O^+ ions both in altitude and MLAT after ~ 30 min is given on the upper right panel. The
1266 maximum time that the electric field was applied was 1 hr. The ionosphere at that time is shown
1267 on the lower left. The storm time equatorial ionospheric anomalies (EIAs) are located at $|\text{MLAT}|$
1268 $\sim 30^\circ$ to 40° and an altitude of ~ 550 to 900 km for the most dense portion of the EIAs. The bottom
1269 right panel shows that the EIAs have come down in altitude but to higher latitudes ~ 15 min after
1270 the termination of the PPEF application. Parts of the still intense EIAs are now beyond $|\text{MLAT}| >$
1271 40° and now the bulk of the maximum density portion is at ~ 400 to 800 km altitude.

1272
1273 It was found that at altitudes of ~ 700 to $1,000$ km, the O^+ densities are predicted to be ~ 300 times
1274 that of the quiet time neutral densities. It has been also been shown by Tsurutani and Lakhina
1275 (2014) that in extreme cases, the magnetospheric/ionospheric electric field can be twice as large
1276 as the Carrington storm and six times as large as the 1991 event. Even if the magnetospheric
1277 radiation belt is saturated (there are other scientific papers that state that magnetospheric beta can
1278 be greater than one: Chan et al. 1994; Saitoh et al. 2014; Nishiura et al., 2015), this is a different
1279 facet of Space Weather and the electric field may not be saturated. What will be the ionospheric
1280 effects of these even larger electric fields?

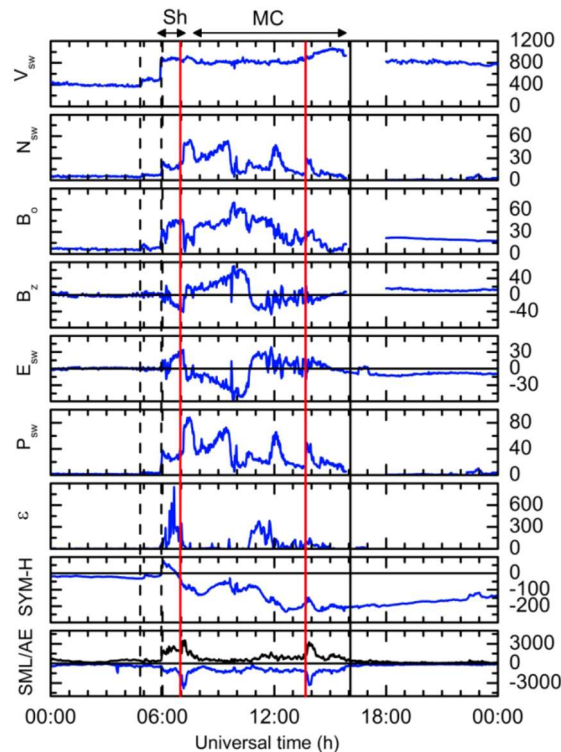
1281
1282 A fundamental question for the future is “can the upward O^+ ion flow drag sufficient numbers of
1283 oxygen neutrals upward so that the oxygen ions plus neutral densities are even higher still?” A
1284 short-time interval analytic calculation done by Lakhina and Tsurutani (2017) and a mini-
1285 Carrington event modeled by Deng et al. (2018) have indicated that the answer is “yes”. However
1286 a full code needs to be developed and run to answer this question quantitatively. This is an
1287 interesting future problem for computer modelers.

1288

1289 **8. RESULTS: Supersubstorms**

1290

1291 Super intense substorms (supersubstorms: SSSs) appear to be externally (solar wind) triggered.
 1292 Why are they important? They might be the feature within extreme magnetic storms that cause
 1293 geomagnetically induced currents (GICs)/power outages. This hypothesis needs to be tested.
 1294



1295
 1296 Figure 29. Two supersubstorms (SSSs) that occur during a two-phase magnetic storm on 20
 1297 November 2001. The onsets of the supersubstorms are indicated by the vertical red lines. The
 1298 figure is taken from Tsurutani et al. (2015).

1299
 1300 Figure 29 shows the solar wind data during an intense magnetic storm and two SSSs. From top to
 1301 bottom are the solar wind speed and density, the magnetic field magnitude and Bz component, and
 1302 the interplanetary motional electric field, ram pressure and Akasofu epsilon parameter (Perreault
 1303 and Akasofu, 1978). The bottom two parameters are the SYM-H index and the SML index (blue)
 1304 and AE index (black). An initial forward shock is indicated by a vertical dashed line at ~0500 UT,
 1305 a second shock at ~0600 UT, and the two SSS onsets by red vertical lines. The criterion for a SSS
 1306 event was a SML peak value < -2500 nT (an arbitrary number, but chosen to be an extremely high
 1307 value). At the top of the diagram, the sheath region is indicated by a “Sh” and the magnetic cloud
 1308 region by “MC”. The first storm main phase is caused by southward Bz in the sheath and the

1309 second, more intense main phase by southward Bz in the MC. The interplanetary magnetic field
1310 measurement cadence is 1 min. It has been noted that the magnetosphere typically reacts to
1311 southward Bz with durations > 10 to 15 min (Tsurutani et al., 1990), so this high rate of cadence
1312 is sufficient to identify any causes of geomagnetic response.

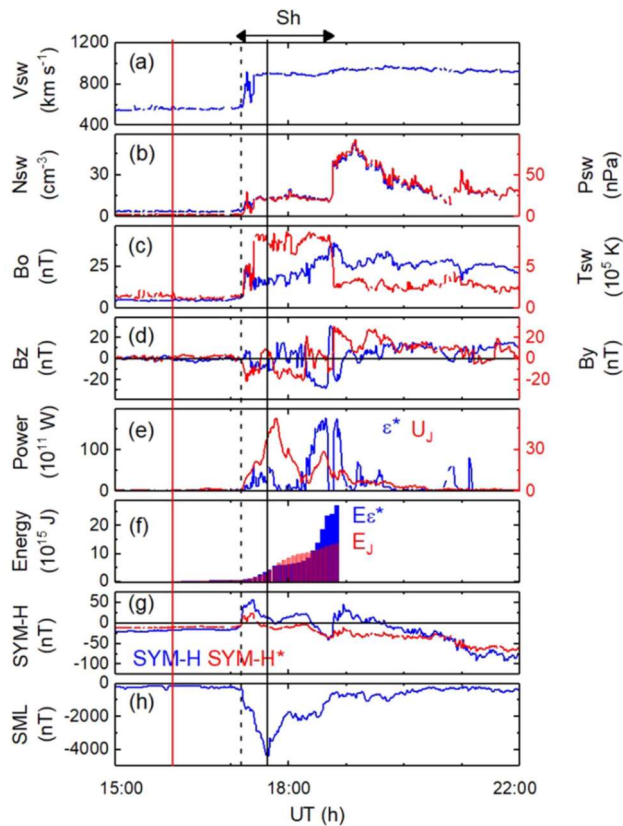
1313

1314 It is noted that the SSS events in this case are not triggered at either of the two shocks nor do they
1315 occur during the peak negative SYM-H values of the storm main phases. However, the first SSS
1316 event is collocated with a peak Esw and a peak southward Bz of the sheath plasma. The SSS event
1317 is also collocated with a large solar wind pressure pulse which is caused by an intense solar wind
1318 density feature. The second SSS event occurred in the recovery phase of the second magnetic
1319 storm. The IMF Bz was ~ 0 nT. The second SSS event was associated with a solar wind pressure
1320 pulse associated with a small density enhancement.

1321

1322 A study of SSSs from 1981 to 2012 was conducted by Hajra et al. (2016). In that study a variety
1323 of solar wind features were found to be associated with SSS onsets. In that survey it was noted
1324 that two SSS events were triggered by fast forward shocks. One of these events will be discussed
1325 below.

1326



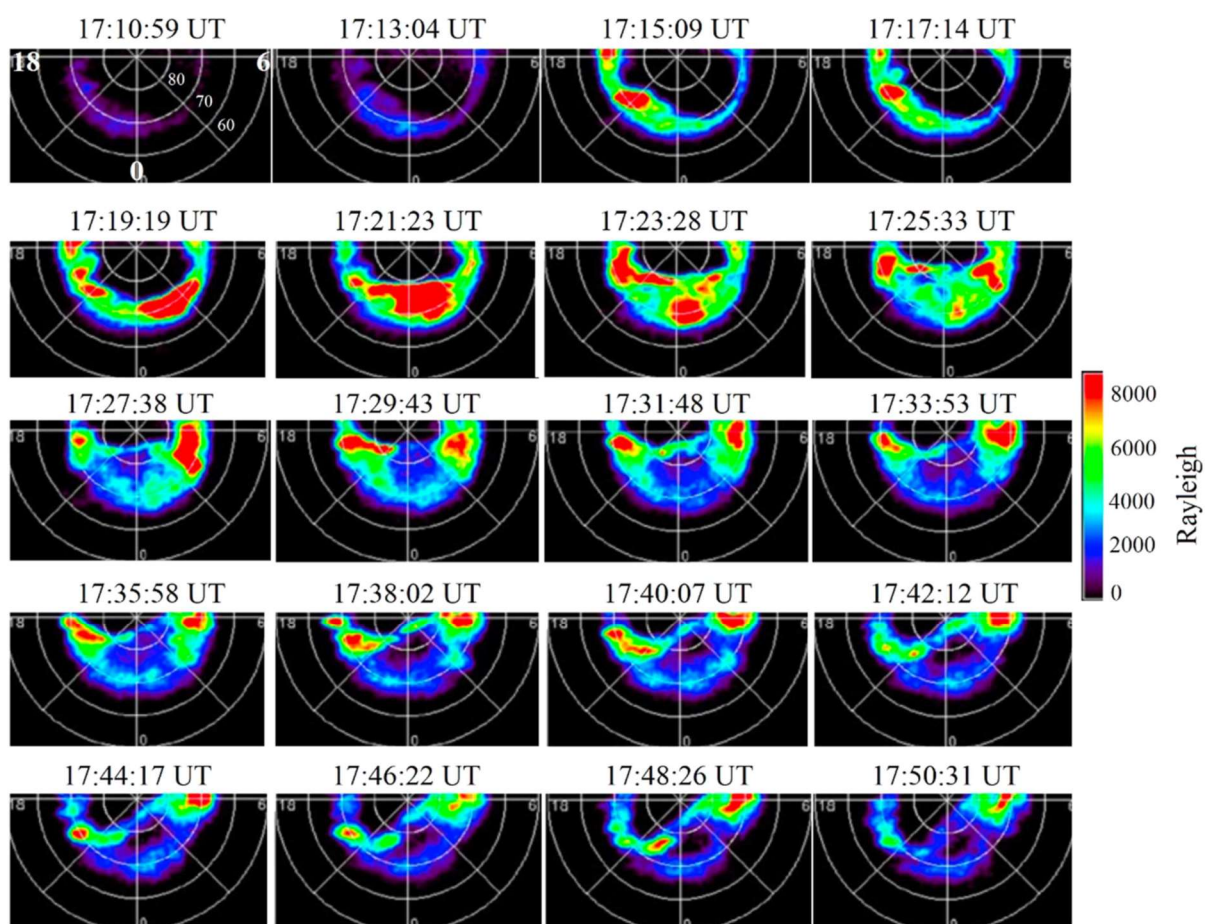
1327
 1328 Figure 30. An SSS triggered by an interplanetary shock on 21 January 2005. The dashed vertical
 1329 line indicates a fast forward shock and the solid black line the peak intensity of the SSS event. The
 1330 figure is taken from Hajra and Tsurutani (2018b).

1331
 1332 Figure 30 shows solar wind/interplanetary parameters and geomagnetic parameters during a SSS
 1333 event on 21 January 2005. From top to bottom are the solar wind speed, density and ram pressure,
 1334 the magnetic field magnitude and solar wind temperature (in the same panel), the IMF Bz and By
 1335 components (GSM coordinates), Joule energy and the Akasofu epsilon pressure corrected
 1336 parameter ϵ^* , the time-integrated energy input into the magnetosphere and time-integrated joule
 1337 energy. The next to the bottom panel contains the SYM-H index and the pressure corrected SYM-
 1338 H index (SYM-H*). The bottom panel is the SML index. A dashed vertical line denotes the
 1339 occurrence of a fast forward shock. A vertical solid line indicates the peak of the SSS event.

1340
 1341 The SSS event onset at 1711 UT coincided with a shock with magnetosonic Mach number of ~ 5.5
 1342 with a shock normal angle of 81° . The high-density sheath sunward of the shock causes a SI^+ of
 1343 ~ 57 nT. The solar feature associated with this event was an X7 class flare that occurred at ~ 0700

1344 UT January 20 (Bombardieri et al., 2008; Saldanha et al., 2008; Pérez-Peraza et al., 2009; Wang
 1345 et al., 2009; Firoz et al., 2012; Bieber et al., 2013; Tan, 2013). The IMF Bz turned abruptly
 1346 southward at the time of the shock so this is part of the energy driving the event. When the IMF
 1347 Bz turned abruptly northward at ~ 1738 UT, the SSS began a recovery phase. This was followed
 1348 by an interplanetary solar filament (Kozyra et al., 2013), but the latter was not geoeffective in this
 1349 case. This high plasma density, high magnetic field intensity feature was interpreted by Kozyra et
 1350 al. (2013) as the interplanetary manifestation of the Illing and Hundhausen (1986) most sunward
 1351 portion of the 3 parts of a CME discussed earlier.

1352



1353
 1354 Figure 31. IMAGE-FUV images taken from ~ 1711 UT to ~ 1751 UT on January 21, 2005. These
 1355 selected auroral images correspond to the SSS event in Figure 30.

1356
 1357 Figure 31 contains the Imager for Magnetopause-to-Aurora global Exploration (IMAGE) far
 1358 ultraviolet images for the SSS event in Figure 30. At ~ 1713 UT there was a small brightening at

1359 ~68° MLAT, which was a very small substorm or pseudobreakup (Elvey, 1957; Tsurutani et al.,
1360 1998; Aikio et al., 1999). At ~1715 UT, 2 min later there was a ~2100 MLT premidnight
1361 brightening of the aurora at ~68° to 75°. At ~1719 UT the most intense aurora was located at ~68°
1362 to 72° in the postmidnight/morning sector, ~0000 to 0400 MLT. The aurora moved from a
1363 dominant premidnight location to a postmidnight location in ~4 min.

1364
1365 By ~1726 UT there was almost no aurora of significant intensity at local midnight. At the peak of
1366 the SML value at ~1738 UT until ~1751, there were both intense premidnight and postmidnight
1367 auroras.

1368
1369 The SSS event did not exhibit the Akasofu (1964) standard model of a substorm with an
1370 intensification at midnight and then expansion to the west, east and north. The changes in the
1371 location of intense auroras were too rapid to track with the IMAGE cadence of ~2 min.

1372
1373 The SSS events display rapid auroral movements which may entail the appearance of sudden local
1374 field-aligned currents. Even smooth motion of auroral forms will cause strong dB/dt effects over
1375 local ground stations. SSS events may be features that can cause GIC effects that have been
1376 attributed to “magnetic storms”. Thus, it might be the SSS events within magnetic storms which
1377 are the real cause. SWARM satellites are excellently instrumented spacecraft that can study the
1378 SSS events in detail and possible resultant GIC effects. However as noted in the auroral images,
1379 there is a need for even higher time resolution global images than is present today. Therefore, it is
1380 important to development and fly auroral UV imagers that can be operated at ~1 s cadence in
1381 intense auroral substorm events.

1382
1383 **9. CONCLUSIONS: The Physics of Space Weather/ Solar Terrestrial Physics and Possible**
1384 **Forecasting**

1385
1386 We have discussed the current knowledge about various facets of the physics of Space
1387 Weather/Solar Terrestrial Physics (our thought is that since everything in solar terrestrial physics
1388 is interconnected, it is the same thing as space weather). There are others which we have not

1389 touched upon because of limited time and knowledge. The reader should know that other areas of
1390 Space Weather/Solar Terrestrial Physics exist which may be equally important.

1391
1392 The most critical area for forecasting magnetic storms, either during solar maximum or the
1393 declining phase of the solar cycle is the prediction of the magnetic field Bz and the speed of the
1394 convected fields at 1 AU. For CME/MC storms (primarily during solar maximum), this is
1395 identifying MC Bz fields near the Sun and understanding the evolution of the MC as it propagates
1396 from the Sun to Earth. This major challenge will be applicable for the prediction of extreme
1397 magnetic storms and hopefully great progress will be made in the next 5 to 10 years. It was shown
1398 that for simple MCs for extreme storms one need to focus on events where the transit time from
1399 the Sun to the Earth is less than ~24 hours.

1400
1401 For sheaths upstream of ICMEs during solar maximum and CIRs during the declining phase (CIRs
1402 are double sheath structures), the problem is different. Detailed knowledge of the slow solar wind
1403 in the space between the Sun and Earth are needed to accurately describe and predict the IMF Bz
1404 that impacts the Earth. So far little work has been applied towards predicting the slow solar wind
1405 (plus verification). Effort needs to be placed in this area to be able to forecast intense to moderate
1406 magnetic storms. It was shown that sheath magnetic fields are extremely important for the
1407 generation of super intense ($Dst < -250$ nT) magnetic storms (Meng et al., 2019).

1408
1409 A great deal of knowledge presently exists for establishing SEP events, those energetic particles
1410 associated with acceleration at ICME shock fronts (see Luhmann et al., 2017). What is needed for
1411 better forecasting is to understand the Mach number of the shocks, the shock normal angles and
1412 possibly upstream “seed” particles. The upstream seed particle population is similar to the sheath
1413 Bz problem in that this component of the slow solar wind needs to be modeled carefully and
1414 accurately. Three spacecraft in the solar wind at different distances from the Sun should help a
1415 lot.

1416
1417 The appearance of HSSs at 1 AU is a very tractable problem. That is if the coronal hole boundaries
1418 in the photosphere can be established firmly and the HSS propagation to 1 AU can be done
1419 accurately. However, the most difficult task again is the IMF Bz. If Alfvén waves are generated

1420 in the interplanetary medium, this will make the task even more difficult. One solution is to
1421 measure the interplanetary magnetic field at 1 AU and use filtering techniques (Guarnieri et al.
1422 2018) or again have large apogee Earth orbiters like the IMP-8 spacecraft again. Another
1423 possibility is developing some type of statistical IMF Bz generator. Of course, this technique will
1424 only give a ~30 min to 1 hr advanced warning.

1425
1426 Predicting the interplanetary shock Mach numbers and ram pressure jumps will allow
1427 foreknowledge of new radiation belt formation, SI⁺ effects and magnetospheric and ionospheric
1428 dB/dt effects. Dayside auroral intensities and nightside substorm triggering will also be enhanced
1429 by predicting incoming shocks.

1430
1431 Several spacecraft missions have been mentioned in relationship to some forecasting problems.
1432 However, the reader should note that the missions and/or their data alone will not solve these
1433 problems. It will be the scientists either on these missions or perhaps totally independent scientists
1434 who will make the most progress on these problems. An example is magnetic storms caused by
1435 interplanetary shocks/sheaths and CIRs. How long will it take scientists to be able to accurately
1436 forecast the time of occurrence of the storm (the easiest part) and the intensity (the hardest part)?
1437 Here we will not make an estimate of how long this will take. Shock acceleration of solar flare
1438 particles is clearly a fundamental part of Space Weather. How long will scientists take to be able
1439 to predict the fluence and spectral shape at a variety of distances away from the Sun? This is a
1440 fundamental problem which space agencies are not currently directly addressing.

1441
1442 **10. FINAL COMMENTS**

1443
1444 A great amount of effort has been put into developing Space Weather models with the appropriate
1445 physics and chemistry included. Some models even use solar and solar wind data and geomagnetic
1446 indices that might be useful for short time-duration predictions (Gopalswamy et al., 2001;
1447 Srivastava, 2005; Cho et al., 2010; Kim et al., 2010; Kim et al., 2014; Schrijver et al., 2015; Savani
1448 et al., 2015). However, in most cases, the usefulness of such models for predictive purposes has
1449 not been independently and objectively tested. This needs to be done so that missing physics and
1450 chemistry can be applied. When done (testing), surprises might result. It is now being realized

1451 that not only the predictability of various models need improvement, but also the level of
1452 uncertainty of prediction needs to be assessed as well (Knipp et al., 2018; Savani et al., 2017).

1453
1454 CME propagation through the interplanetary medium using ENLIL-based codes are making good
1455 progress in estimating arrival times of ICMEs at 1 AU and have had varying success in predicting
1456 the solar wind parameters as well (Falkenberg et al., 2010; Davis et al., 2011; Pizzo et al. 2015;
1457 Jackson et al., 2015; Jian et al. 2015, 2016). However, the fundamental issue of space weather
1458 prediction for magnetic storms is the direction and intensity of the magnetic field both in the MC
1459 and upstream sheath. These topics still remain a challenge.

1460
1461 Another new approach, the application of machine learning algorithms, is quite hopeful. For this
1462 application, the physics and chemistry need not be known to be applied. Rather the reverse, finding
1463 good correlations between solar and interplanetary parameters and magnetospheric observations
1464 (for magnetic storms as an example) could lead to a better understanding of the physics, the topic
1465 of this paper. But again, one should test these approaches and carefully and objectively assess
1466 their accuracy and reliability in making predictions (see Wing et al., 2005, 2016; Reikard, 2015,
1467 2018).

1468
1469 The best test for proving that workers in Space Weather understand all of the underlying physics
1470 and/or the machine learning algorithm is robust is to use the program on a new event and see how
1471 well it does. This should be done by independent researchers like the people at CCMC at the
1472 Goddard Space Research Center, Greenbelt Maryland and other facilities.

1473
1474 We have one final comment on a third type of approach at predicting Space Weather. For
1475 atmospheric weather forecasts, the experts downselect to ~25 of their best codes, and run each of
1476 the codes with the same input data (Yun et al., 2005; Ruiz et al., 2009; Ghosh and Krishnamurti
1477 2018). The codes produce ~25 different predictions. The weather service uses the average of the
1478 values. Why this scheme works reasonably well is not understood. This may be the final path of
1479 Space Weather forecasting.

1480

1481 Our hope is that the paper is stimulating to the reader in a positive sense: that they will be energized
 1482 to attack some of the interesting problems in our field of Space Weather. On the other hand, if the
 1483 reader finds statements/topics that they disagree with, please send us email comments and we will
 1484 try to answer them the best that we can. And if you have disagreements that should see print,
 1485 Nonlinear Processes in Geophysics has a “Comment” and “Reply” format for discussions of this
 1486 type.

1487

1488

11. GLOSSARY

1489 **Partially taken from: “*From the Sun: Auroras, Magnetic Storms, Solar Flares, Cosmic Rays*”**
 1490 (Suess and Tsurutani, 1998, AGU Press)

1491

1492 **Adiabatic Invariant:** In a nearly collisionless, ionized gas, electrically charged particles orbit
 1493 around magnetic lines of force. Certain physical quantities are approximately constant for slow
 1494 (adiabatic) changes of the magnetic field in time or in space and these quantities are called
 1495 *adiabatic invariants*. For example, the magnetic moment of a charged particle, $\mu = mV_{\perp}^2 / (2B)$,
 1496 is such a constant where V_{\perp} is the velocity of the particle perpendicular to the magnetic field,
 1497 B is the magnetic field strength, and m is the particle mass. In a converging field such as in
 1498 approaching the pole of a dipole magnetic, the field strength increases and therefore V_{\perp}
 1499 increases as well because μ has to remain constant.

1500 **Aeronomy:** The science of the (upper) regions of atmospheres, those regions where dissociation
 1501 of molecules and ionization are present.

1502 **Alfvén Wave** (magnetohydrodynamic shear wave): A transverse wave in magnetized plasma
 1503 characterized by a change of direction of the magnetic field with no change in either the
 1504 intensity of the field or the plasma density.

1505 **Anisotropic Plasma:** A Plasma whose properties vary with direction relative to the ambient
 1506 magnetic field direction. This can be due, for example, to the presence of a magnetic or electric
 1507 field. See also Isotropic Plasma; Plasma.

1508 **Arase satellite, formerly called Exploration of energization and Radiation in Geospace** or
 1509 **ERG:** a scientific satellite developed by the Institute of Space and Astronautical Science (ISAS)
 1510 of the Japanese Aerospace Exploration Agency (JAXA) to study the Van Allen radiation belts.

1511 **Astronomical Unit (AU):** The mean radius of the Earth’s orbit, 1.496×10^{13} cm.

1512 **Aurora:** A visual phenomenon that occurs mainly in the high-latitude night sky. Auroras occur
 1513 within a band of latitudes known as the auroral oval, the location of which is dependent on the
 1514 intensity of geomagnetic activity. Auroras are a result of collisions between precipitating
 1515 charged particles (mostly electrons) and atmospheric atoms and molecules, exciting the
 1516 atmospheric constituents. The charged particles come from the outer parts of the magnetosphere
 1517 and guided by the geomagnetic field. Each gas (oxygen and nitrogen molecules and atoms)
 1518 emits its own characteristic radiation when bombarded by the precipitating particles. Since the
 1519 atmospheric composition varies with altitude, and the faster precipitating particles penetrate
 1520 deeper into the atmosphere, certain auroral colors originate preferentially from certain heights
 1521 in the sky. The auroral altitude range is 80 to 500 km, but typical auroras occur 90 to 250 km
 1522 above the ground. The color of the typical aurora is yellow-green, from a specific transition line
 1523 of atomic oxygen. Auroral light from lower levels in the atmosphere is dominated by blue and
 1524 red bands from molecular nitrogen and molecular oxygen. Above 250 km, auroral light is
 1525 characterized by a red spectral line of atomic oxygen. To an observer on the ground, the
 1526 combined light of these three fluctuating, primary colors produces an extraordinary visual
 1527 display. Auroras in the Northern Hemisphere are called the aurora borealis or “northern lights”.
 1528 Auroras in the Southern Hemisphere are called aurora australis. The patterns and forms of the
 1529 aurora include quiescent “arcs”, rapidly moving “rays” and “curtains”, “patches”, and “veils”.

1530 **Auroral Electrojet (AE):** See Electrojet.

1531 **Auroral Oval:** An elliptical band around each geomagnetic pole ranging from about 75 degrees
 1532 magnetic latitude at local noon to about 67 degrees magnetic latitude at midnight under average
 1533 conditions. It is the locus of those locations of the maximum occurrence of auroras, and widens
 1534 to both higher and lower latitudes during the expansion phase of a magnetic substorm.

1535 **Beta** (e.g., low-beta plasma): The ratio of the thermal pressure to the magnetic ‘pressure’ in a
 1536 plasma - $p / (B^2 / (8\pi))$ in centimeter-gram-second (c.g.s.).

1537 **Bow Shock** (Earth, heliosphere): A collisionless shock wave in front of the magnetosphere arising
 1538 from the interaction of the supersonic solar wind with the Earth's magnetic field. An analogous
 1539 shock is the heliospheric bow shock which exists in front of the heliosphere and is due to the
 1540 interaction of the interstellar wind with the solar wind and the inter planetary magnetic field.

1541 **Charge Exchange:** An interaction between a charged particle and a neutral atom wherein the
1542 charged particle becomes neutral and the neutral particle becomes charged through the
1543 exchange of an electron.

1544 **Cloud** (magnetic): see Magnetic Cloud.

1545 **Collisional (de-) Excitation:** Excitation of an atom or molecule to a higher energy state due to a
1546 collision with another atom, molecule, or ion. The higher energy state generally refers to
1547 electrons in higher energy around atoms. Deexcitation is the reduction of a higher electron
1548 energy state to a lower one, usually accomplished by a collision with another atom, molecule
1549 or ion.

1550 **Convection** (magnetospheric, plasma, thermal): The bulk transport of plasma (or gas) from one
1551 place to another, in response to mechanical forces (for example, viscous interaction with the
1552 solar wind) or electromagnetic forces. Thermal convection, due to heating from below and the
1553 gravitational field, is what drives convection inside the Sun. Magnetospheric convection is
1554 driven by the dragging of the Earth's magnetic field and plasma together by the solar wind when
1555 the magnetic field becomes attached to the magnetic field in the solar wind.

1556 **Coriolis Force:** In the frame of a rotating body (such as the Earth), a force due to the bodily
1557 rotation. All bodies that are not acted upon by some force have the tendency to remain in a state
1558 of rest or of uniform rectilinear motion (Newton's First Law) so that this force is called a
1559 "fictitious" forces. It is a consequence of the continuous acceleration which must be applied to
1560 keep a body at rest in a rotating frame of reference.

1561 **Corona:** The outermost layer of the solar atmosphere, characterized by low densities ($<10^9 \text{ cm}^{-3}$
1562 or 10^{15} m^{-3}) and high temperatures ($>10^6 \text{ K}$).

1563 **Coronal Hole:** An extended region of the solar corona characterized by exceptionally low density
1564 and in a unipolar photospheric magnetic field having "open" magnetic field topology. Coronal
1565 holes are largest and most stable at or near the solar poles, and are a source of high speed (700-
1566 800 km/s) solar wind. Coronal holes are visible in several wavelengths, most notably solar x-
1567 rays visible only from space, but also in the He 1083 nm line which is detectable from the
1568 surface of the Earth. In soft x-ray images (photon energy of $\sim 0.1-1.0 \text{ keV}$ or a wavelength of
1569 $10-100 \text{ \AA}$), these regions are dark, thus the name "holes".

1570 **Coronal Mass Ejection (CME):** A transient outflow of plasma from or through the solar corona.
1571 CMEs are often but not always associated with erupting prominences, disappearing solar
1572 filaments, and flares.

1573 **Corotation** (with the Earth): A plasma in the magnetosphere of the Earth is said to be corotating
1574 with the Earth if the magnetic field drags the plasma with it and together they have a 24 hour
1575 rotation period.

1576 **Cosmic Ray** (galactic, solar): Extremely energetic (relativistic) charged particles or
1577 electromagnetic radiation, primarily originating outside of the Earth's magnetosphere. Cosmic
1578 rays usually interact with the atoms and molecules of the atmosphere before reaching the
1579 surface of the Earth. The nuclear interactions lead to formation of daughter products, and they
1580 in turn to granddaughter products, etc. Thus, there is a chain of reactions and a "cosmic ray
1581 shower". Some cosmic rays come from outside the solar system while others are emitted from
1582 the Sun in solar flares. See also Anomalous Cosmic Ray; Energetic Particle; Solar Energetic
1583 Particle (SEP) Event.

1584 **Constellation Observing System for Meteorology, Ionosphere and Climate-2 (COSMIC II):**
1585 A joint Taiwan National Space Organization (NSPO)-U.S. National Oceanic and Atmospheric
1586 Administration (NOAA) mission of six satellites in low-inclination orbit to study the Earth's
1587 ionosphere.

1588 **Corotating Interaction Region (CIR):** An interplanetary region of high magnetic fields and
1589 plasma densities created by the interaction of a high speed solar wind stream with the upstream
1590 slow solar wind. The antisunward portion of the CIR is compressed slow solar wind plasma
1591 and magnetic fields, and the sunward portion is compressed fast solar wind plasma and
1592 magnetic fields. The two regions of the CIR are separated by a tangential discontinuity.

1593 **Cyclotron Frequency:** When a particle of charge q moves in a magnetic field B , the particle orbits,
1594 or gyrates around the magnetic field lines. The cyclotron frequency is the frequency of this
1595 gyration, and is given by $\omega_c = q|B|/mc$, where m is the mass of the particle, and c is the velocity
1596 of light (in centimeter-gram-second (c.g.s.) units).

1597 **Cyclotron Resonance:** The frequency at which a charged particle experiences a Doppler-shifted
1598 wave at the particle's cyclotron frequency. Because the particle and wave may be traveling at
1599 different speeds and in different directions, there is usually a Doppler shift involved.

1600 **D Region:** A daytime region of the Earth's ionosphere beginning at approximately 40 km,
1601 extending to 90 km altitude. Radio wave absorption in this region can be significantly increased
1602 due to increasing ionization associated with the precipitation of solar energetic particles through
1603 the magnetosphere and into the ionosphere.

1604 **Diffusion:** The slow, stochastic motion of particles.

1605 **Diffusive Shock Acceleration:** Charged particle acceleration at a collisionless shock due to
1606 stochastic scattering processes caused by waves and plasma turbulence. See also Shock Wave
1607 (collisionless).

1608 **Dipole Magnetic Field:** A magnetic field whose intensity decreases as the cube of the distance
1609 from the source. A bar magnet's field and the magnetic field originating in the Earth's core are
1610 both approximately dipole magnetic fields.

1611 **Drift** (of ions/electrons): As particles gyrate around magnetic field lines, their orbits may “drift”
1612 perpendicular to the local direction of the magnetic field. This occurs if there is a force also
1613 perpendicular to the field - e.g. an electric field, curvature in the magnetic field direction, or
1614 gravity.

1615 **Driver Gas:** A mass of plasma and entrained magnetic field that is ejected from the Sun, that has
1616 a velocity higher than the upstream plasma, and which “drives” a (usually collisionless) shock
1617 wave ahead of itself. The magnetic cloud within an ICME is the same thing as a driver gas.

1618 **Dst Index:** A measure of variation in the geomagnetic field due to the equatorial ring current. It is
1619 computed from the H-components at approximately four near-equatorial stations at hourly
1620 intervals. At a given time, the Dst index is the average of variation over all longitudes; the
1621 reference level is set so that Dst is statistically zero on internationally designated quiet days. An
1622 index of -50 nT (nanoTesla) or less indicates a storm-level disturbance, and an index of -200
1623 nT or less is associated with middle- latitude auroras. Dst is determined by the World Data
1624 Center C2 for Geomagnetism, Kyoto University, Kyoto, Japan.

1625 **Dynamo** (solar magnetospheric): The conversion of mechanical energy (rotation in the case of the
1626 Sun) into electrical current. This is the process by which magnetic fields are amplified by the
1627 induction of plasmas being forced to move perpendicular to the magnetic field lines. See also
1628 Mean Field Electro-Dynamics.

1629 **E-Region:** A daytime region of the Earth's ionosphere roughly between the altitudes of 90 and
1630 160 km. The E-region characteristics (electron density, height, etc.) depend on the solar zenith

1631 angle and the solar activity. The ionization in the E layer is caused mainly by x-rays in the range
1632 0.8 to 10.4 nm coming from the Sun.

1633 **Ecliptic Plane:** The plane of the Earth's orbit about the Sun. It is also the Sun's apparent annual
1634 path, or orbit, across the celestial sphere.

1635 **Electrically Charged Particle:** Electrons and protons, for example, or any atom from which
1636 electrons have been removed to make it into a positively charged ion. The elemental charge of
1637 particles is 4.8×10^{-10} esu. An electron and proton have this charge. Combined (a hydrogen
1638 atom), the charge is zero. Ions have multiples of this charge, depending on the number of
1639 electrons which have been removed (or added).

1640 **Electrojet:** (1) Auroral Electrojet (AE): A current that flows in the ionosphere at a height of ~100
1641 km in the auroral zone. (2) Equatorial Electrojet: A thin electric current layer in the ionosphere
1642 over the dip equator at about 100 to 115 km altitude.

1643 **Electron Plasma Frequency/Wave:** The natural frequency of oscillation of electrons in a neutral
1644 plasma (e.g., equal numbers of electrons and protons).

1645 **Electron Volt (eV):** The kinetic energy gained by an electron or proton being accelerated in a
1646 potential drop of one Volt.

1647 **ESA:** European Space Agency

1648 **Extreme Ultraviolet (EUV):** A portion of the electromagnetic spectrum from approximately 10
1649 to 100 nm.

1650 **Extremely Low Frequency (ELF):** That portion of the radio frequency spectrum from 30 to 3000
1651 Hz.

1652 **Fast Mode (wave/speed):** In magnetohydrodynamics, the fastest wave speed possible.
1653 Numerically, this is equal to the square root of the sum of the squares of the Alfvén speed and
1654 plasma sound speed.

1655 **Field Aligned Current:** A current flowing along (or opposite to) the magnetic field direction.

1656 **Filament:** A mass of gas suspended over the chromosphere by magnetic fields and seen as dark
1657 ribbons threaded over the solar disk. A filament on the limb of the Sun seen in emission against
1658 the dark sky is called a prominence. Filaments occur directly over magnetic-polarity inversion
1659 lines, unless they are active.

1660 **Flare:** A sudden eruption of energy in the solar atmosphere lasting minutes to hours, from which
 1661 radiation and energetic charged particles are emitted. Flares are classified on the basis of area
 1662 at the time of maximum brightness in H alpha.

1663 Importance 0 (Subflare): < 2.0 hemispheric square degrees

1664 Importance 1: 2.1-5.1 square degrees

1665 Importance 2: 5.2-12.4 square degrees

1666 Importance 3: 12.5-24.7 square degrees

1667 Importance 4: \geq 24.8 square degrees

1668 [One square degree is equal to $(1.214 \times 10^4 \text{ km})^2 = 48.5$ millionths of the visible
 1669 solar hemisphere.]

1670 A brightness qualifier F, N, or B is generally appended to the importance character to
 1671 indicate faint, normal, or brilliant (for example, 2B).

1672 **Flux Rope:** A magnetic phenomenon which has a force-free field configuration.

1673 **Force Free Field:** A magnetic field which exerts no force on the surrounding plasma. This can
 1674 either be a field with no flowing electrical currents or a field in which the electrical currents all
 1675 flow parallel to the field.

1676 **Free Energy** (of a plasma): When an electron or ion distribution is either non-Maxwellian or
 1677 anisotropic, they are said to have free energy" from which plasma waves can be generated via
 1678 instabilities. The waves scatter the particles so they become more isotropic, reducing the free
 1679 energy.

1680 **Frozen-in Field:** In a tenuous, collisionless plasma, the weak magnetic fields embedded in the
 1681 plasma are convected with the plasma. i.e., they are "frozen in".

1682 **Galactic Cosmic Ray (GCR):** See Cosmic Ray.

1683 **Gamma Ray:** Electromagnetic radiation at frequencies higher than x-rays.

1684 **Geomagnetic Storm:** A worldwide disturbance of the Earth's magnetic field, distinct from regular
 1685 diurnal variations. A storm is precisely defined as occurring when D_{ST} becomes less than -50
 1686 nT (See geomagnetic activity).

1687 Main Phase: Of a geomagnetic storm, that period when the horizontal magnetic field at
 1688 middle latitudes decreases, owing to the effects of an increasing magnetospheric ring
 1689 current. The main phase can last for hours, but typically lasts less than 1 day.

1690 Recovery Phase: Of a geomagnetic storm, that period when the depressed northward field
1691 component returns to normal levels. Recovery is typically complete in one to two days.

1692 **Geomagnetically Induced Currents (GICs):** Currents flowing along electric power transmission
1693 systems and other electrically conducting infrastructures are produced by naturally induce
1694 geoelectric fields during geomagnetic disturbances.

1695 **Geosynchronous Orbit:** Term applied to any equatorial satellite with an orbital velocity equal to
1696 the rotational velocity of the Earth. The geosynchronous altitude is near 6.6 Earth radii
1697 (approximately 36,000 km above the Earth's surface). To be geostationary as well, the satellite
1698 must satisfy the additional restriction that its orbital inclination be exactly zero degrees. The net
1699 effect is that a geostationary satellite is virtually motionless with respect to an observer on the
1700 ground.

1701 **GeV:** 10^9 electron Volts (Giga-electron Volt).

1702 **Global Navigation Satellite System (GNSS):** GNSS receivers use the orbiting satellite Global
1703 Positioning System (GPS) transmitted signals to obtain the geographic location of a user's
1704 receiver anywhere in the world.

1705 **Global Positioning System (GPS):** is a global navigation satellite system that provides
1706 geolocation and time information to a GPS receiver anywhere on or near the Earth where there
1707 is an unobstructed line of sight to four or more GPS satellites.

1708 **Global-scale Observations of the Limb and Disk (GOLD):** a NASA mission to "investigate the
1709 dynamic intermingling of space and Earth's uppermost atmosphere"

1710 **Heliosphere:** The magnetic cavity surrounding the Sun, carved out of the galaxy by the solar wind.

1711 **Heliospheric Current Sheet (HCS):** This is the surface dividing the northern and southern
1712 magnetic field hemispheres in the solar wind. The magnetic field is generally one polarity in
1713 the north and the opposite in the south so just one surface divides the two polarities. However,
1714 the Sun's magnetic field changes over the 11-year solar sunspot cycle and reverses polarity at
1715 solar maximum. The same thing happens in the magnetic field carried away from the Sun by
1716 the solar wind so the HCS only lies in the equator near solar minimum. It is called a "current
1717 sheet" because it carries an electrical current to balance the oppositely directed field on either
1718 side of the surface. It is very thin on the scale of the solar system - usually only a few proton
1719 gyroradii, or less than 100,000 km.

1720 **Helmet Streamer:** See Streamer.

- 1721 **High Frequency (HF):** That portion of the radio frequency spectrum between 3 and 30 MHz.
- 1722 **Heliospheric Plasma Sheet (HPS):** A high density slow solar wind region that is located adjacent
1723 to the heliospheric current sheet (HCS).
- 1724 **High-Speed Solar Wind (HSS):** A solar wind with speeds of 750 to 800 km/s emanating from
1725 solar coronal holes. The HSS is characterized by embedded, particularly large amplitude
1726 Alfvén waves. At the edges of HSSs, the velocities can be less due to superradial expansion
1727 effects.
- 1728 **Instability:** When an electron or ion distribution is sufficiently anisotropic, it becomes unstable
1729 (instability), generating plasma waves. The anisotropic distribution provides a source of free
1730 energy for the instability. A simple analog is a stick, which if stood on end is “unstable”, but
1731 which if laid on its side is “stable”. In this analog, gravity pulls on the stick and provides a
1732 source of free energy when the stick is stood on end.
- 1733 **Interplanetary Magnetic Field (IMF, Parker spiral):** The magnetic field carried with the solar
1734 wind and twisted into an Archimedean spiral by the Sun’s rotation.
- 1735 **Interplanetary Medium:** The volume of space in the solar system that lies between the Sun and
1736 the planets. The solar wind flows in the interplanetary medium.
- 1737 **Interplanetary Coronal Mass Ejection (ICME):** The evolutionary part of a CME as it propagates
1738 through interplanetary space. Typically, after the CME has propagated 1 AU from the Sun, the
1739 ICME only contains the magnetic cloud (MC) portion of the initial three parts of a CME. The
1740 MC may also have been compressed/expanded or rotated by the time it reaches 1 AU.
- 1741 **Interplanetary Shock:** A fast forward shock is characterized by a sharp increase in solar wind
1742 speed, plasma density, plasma temperature and magnetic field magnitude. The shock reduces
1743 the upstream plasma from a supermagnetosonic state to a subsonic state, much as an airplane
1744 wing sonic shock reduces the relative flow of air from a supersonic speed (relative to the
1745 airplane) to a subsonic speed. A fast (magnetosonic) forward (propagating in the direction of
1746 the “piston”, in this case the propagation of the ICME in the antisolar direction) shock is
1747 detected upstream (antisolarward) of fast ICMEs. A reverse shock propagates in the direction
1748 of the Sun. Planetary bow shocks are reverse shocks. There are other types of shocks not
1749 discussed in this paper: slow shocks and intermediate shocks.

1750 **Interstellar** (gas, neutral gas, ions, cosmic rays, wind, magnetic field, etc.) Literally, between the
1751 stars. In practical terms, it is anything beyond the outer boundary of the solar wind (the
1752 “heliopause”) yet within the Milky Way.

1753 **Ion:** (1). An electrically charged atom or molecule. (2). An atom or molecular fragment that has a
1754 positive electrical charge due to the loss of one or more electrons; the simplest ion is the
1755 hydrogen nucleus, a single proton.

1756 **Ionization State:** The number of electrons missing from an atom.

1757 **Ionosphere:** The region of the Earth's upper atmosphere containing free (not bound to an atom or
1758 molecule) electrons and ions. This ionization is produced from the neutral atmosphere by solar
1759 ultraviolet radiation at very short wavelengths (<100 nm) and also by precipitating energetic
1760 particles.

1761 **Ionospheric Storm:** A positive ionospheric storm is where the ionospheric total electron content
1762 (TEC) increases. A negative ionospheric storm is an event where the ionospheric TEC
1763 decreases.

1764 **Ionospheric Connection Explorer (ICON):** is a NASA 2-year mission that will give new views
1765 of the boundary between our atmosphere and space, where planetary weather and Space
1766 Weather meet.

1767 **Irradiance:** Radiant energy flux density on a given surface (e. g. $\text{ergs cm}^{-2}\text{s}^{-1}$).

1768 **keV:** 1000 electron Volts (kiloelectron Volt). See electron Volt. See also Anisotropic Plasma;
1769 Plasma.

1770 **L value:** For a dipole magnetic field, the field line that crosses the magnetic equator at a L value
1771 equal to the number in Earth radii.

1772 **Loop** (solar-loop prominence system): A magnetic loop is the flux tube which crosses from one
1773 polarity to another. A loop prominence bridges a magnetic inversion line across which the
1774 magnetic field changes direction. See also Magnetic Foot Point; Prominence.

1775 **Loss Cone:** A small cone angle about the ambient magnetic field direction where magnetospheric
1776 charged particles with velocity vectors within the cone will mirror at sufficiently low altitudes
1777 such that the particle will have collisions with atmospheric atoms and molecules and will be
1778 “lost” from returning to the magnetosphere.

1779 **Loss Cone Instability:** An instability generated by a plasma anisotropy where the temperature
1780 perpendicular to the magnetic field is greater than the temperature parallel to the field. This

1781 instability gets its name because this condition exists in the Earth's magnetosphere and the "loss
1782 cone" particles are those that are lost into the upper atmosphere.

1783 **Magnetic Cloud:** A region in the solar wind of about 0.25 AU or more in radial extent in which
1784 the magnetic field strength is high and the direction of one component of the magnetic field
1785 changes appreciably by means of a rotation nearly parallel to a plane. Magnetic clouds may be
1786 parts of the driver gases (coronal mass ejections) in the interplanetary medium.

1787 **Magnetic Foot Point:** For the Earth's magnetic field lines, where the magnetic field enters the
1788 surface of the Earth.

1789 **Magnetic Mirror:** Char particles moving into a region of converging magnetic flux (as at the pole
1790 of a magnet) will experience "Lorentz" forces that slow the particles and "mirror" them by
1791 eventually reversing their direction if the particles are initially moving slowly enough along the
1792 field line. See also Mirror Point.

1793 **Magnetic Reconnection:** The act of interconnection between oppositely directed magnetic field
1794 lines. Magnetic reconnection is recognized as a basic plasma process, which converts magnetic
1795 energy into plasma kinetic energy accompanied by topological changes in the magnetic field
1796 configuration. It does not allow an excessive buildup of magnetic energy in the current sheets.

1797 **Magnetic Storm:** see Geomagnetic Storm.

1798 **Magnetopause:** The boundary surface between the solar wind and magnetosphere, where the
1799 pressure of the magnetic field of the object effectively equals the ram pressure of the solar wind
1800 plasma.

1801 **Magnetosheath:** The region between the bow shock and the magnetopause, characterized by very
1802 turbulent plasma. This plasma has been heated (shocked) and slowed as it passed through the
1803 bow shock. For the Earth, along the Sun-Earth axis, the magnetosheath is about 3 Earth radii
1804 thick.

1805 **Magnetosonic Speed** (acoustic speed): The speed of the fastest low frequency magnetic waves in
1806 a magnetized plasma. It is the equivalent of the sound speed in a neutral gas or non-magnetized
1807 plasma.

1808 **Magnetosphere:** The magnetic cavity surrounding a magnetized planet, carved out of the passing
1809 solar wind by virtue of the planetary magnetic field, which prevents, or at least impedes, the
1810 direct entry of the solar wind plasma into the cavity.

1811 **Magnetospheric Multiscale Mission (MMS):** A NASA mission designed to spend extensive
1812 periods in locations where magnetic reconnection at the magnetopause/magnetotail is expected
1813 to occur. The critical electron diffusion region will be studied. The mission consists of 4
1814 spacecraft flown in a tetrahedron configuration.

1815 **Magnetotail:** The extension of the magnetosphere in the anti-sunward direction as a result of
1816 interaction with the solar wind. In the inner magnetotail, the field lines maintain a roughly
1817 dipolar configuration. But at greater distances in the anti-sunward direction, the field lines are
1818 stretched into northern and southern lobes, separated by a plasmashet. There is observational
1819 evidence for traces of the Earth's magnetotail as far as 1000 Earth radii downstream, in the anti-
1820 solar direction.

1821 **Maxwellian Distribution:** The minimum energy particle distribution for a given temperature. It
1822 is also the equilibrium distribution in the absence of losses due to radiation, collisions, etc.

1823 **Mean Free Path:** The statistically most probably distance a particle travels before undergoing a
1824 collision with another particle or interacting with a wave.

1825 **Mesosphere:** The region of the Earth's atmosphere between the upper limit of the stratosphere
1826 (approximately 30 km altitude) and the lower limit of the thermosphere (approximately 80 km
1827 altitude).

1828 **MeV:** One million electron Volts (Megaelectron Volt). See also Electron Volt.

1829 **Mirror Point:** The point where the charged particles reverse direction (mirrors). At this point, all
1830 of the particle motion is perpendicular to the local ambient magnetic field. See also Magnetic
1831 Mirror.

1832 **Parker Solar Probe:** a NASA robotic spacecraft to probe the outer corona of the Sun. It will
1833 approach to within 9.9 solar radii (6.9 million kilometers or 4.3 million miles from the center
1834 of the Sun and will travel, at closest approach, as fast as 690,000 km/h (430,000 mph).

1835 **Photosphere:** The lowest visible layer of the solar atmosphere; corresponds to the solar surface
1836 viewed in white light. Sunspots and faculae are observed in the photosphere.

1837 **Pickup Ion:** An ion which has entered the solar system as a neutral particle and then becomes
1838 ionized either through charge exchange or photoionization. It is called a pickup ion because as
1839 soon as the neutral atom is ionized, it becomes attached to the magnetic field carried by the
1840 solar wind and so is "picked up" by the solar wind.

1841 **Pitch Angle:** In a plasma, the angle between the instantaneous velocity vector of a charged particle
1842 and the direction of the ambient magnetic field.

1843 **Plasma** (ions, electrons): A gas that is sufficiently ionized so as to affect its dynamical behavior.
1844 A plasma is a good electrical conductor and is strongly affected by magnetic fields. See also
1845 Anisotropic Plasma; Isotropic Plasma.

1846 **Plasma Instability** (ion, electron): When a plasma is sufficiently anisotropic, plasma waves grow,
1847 which in turn alter the distribution via wave-particle interactions. The plasma is “unstable”.

1848 **Plasma Sheet:** A region in the center of the magnetotail between the north and south lobes. The
1849 plasma sheet is characterized by hot, dense plasma and is a high beta plasma region, in contrast
1850 to the low beta lobes. The plasma sheet bounds the neutral sheet where the magnetic field
1851 direction reverses from Earthward (north lobe direction) to anti-Earthward (south lobe
1852 direction).

1853 **Plasma Wave** (electrostatic/electromagnetic): A wave generated by plasma instabilities or other
1854 unstable modes of oscillation allowable in a plasma. “Chorus” and “Plasmasheric Hiss” are
1855 whistler wave modes. These are electromagnetic waves with frequencies below the electron
1856 cyclotron frequency. Electromagnetic ion cyclotron (EMIC) waves are ion cyclotron waves
1857 with frequencies below the proton cyclotron frequency.

1858 **Polar Cap Absorption Event (PCA):** An anomalous condition of the polar ionosphere whereby
1859 HF and VHF (3-300 MHz) radio waves are absorbed, and LF and VLF (3-300 kHz) radio waves
1860 are reflected at lower altitudes than normal. The cause is energetic particle precipitation into
1861 the ionosphere/atmosphere. The enhanced ionization caused by this precipitation leads to
1862 cosmic radio noise absorption and attenuation of that noise at the surface of the Earth. PCAs
1863 generally originate with major solar flares, beginning within a few hours of the event (after the
1864 flare particles have propagated to the Earth) and maximizing within a day or two after onset.
1865 As measured by a riometer (relative ionospheric opacity meter), the PCA event threshold is 2
1866 dB of absorption at 30 MHz for daytime and 0.5 dB at night. In practice, the absorption is
1867 inferred from the proton flux at energies greater than 10 MeV, so that PCAs and proton events
1868 are simultaneous. However, the transpolar radio paths may still be disturbed for days, up to
1869 weeks, following the end of a proton event, and there is some ambiguity about the operational
1870 use of the term PCA.

1871 **Prominence:** A term identifying cloud-like features in the solar atmosphere. The features appear
1872 as bright structures in the corona above the solar limb and as dark filaments when seen projected
1873 against the solar disk. Prominences are further classified by their shape (for example, mound
1874 prominence, coronal rain) and activity. They are most clearly and most often observed in H
1875 alpha. See also Loop.

1876 **Radiation Belt:** Regions of the magnetosphere roughly 1.2 to 6 Earth radii above the equator in
1877 which charged particles are stably trapped by closed geomagnetic field lines. There are two
1878 belts. The inner belt's maximum proton density lies near 5000 km above the Earth's surface.
1879 Inner belt protons (10s of MeV) and electrons (100s of keV) and originate from the decay of
1880 secondary neutrons created during collisions between cosmic rays and upper atmospheric
1881 particles. The outer belt extends on to the magnetopause on the sunward side (10 Earth radii
1882 under normal quiet conditions) and to about 6 Earth radii on the nightside. The altitude of
1883 maximum proton density is near 16,000-20,000 km. Outer belt protons and electrons are lower
1884 energy (about 200 eV to 1 MeV). The origin of the particles (before they are energized to these
1885 high energies) is a mixture of the solar wind and the ionosphere. The outer belt is also
1886 characterized by highly variable fluxes of energetic electrons. The radiation belts are often
1887 called the "Van Allen radiation belts" because they were discovered in 1958 by a research group
1888 at the University of Iowa led by Professor J. A. Van Allen. See also Trapped Particle.

1889 **Ram Pressure:** Sometimes called "dynamic pressure". The pressure exerted by a streaming
1890 plasma upon a blunt object.

1891 **Reconnection:** A process by which differently directed field lines link up allowing topological
1892 changes of the magnetic field to occur, determining patterns of plasma flow, and resulting in
1893 conversion of magnetic energy to kinetic and thermal energy of the plasma. Reconnection is
1894 invoked to explain the energization and acceleration of the plasmas/energetic particles that are
1895 observed in solar flares, magnetic substorms and storms, and elsewhere in the solar system.

1896 **Relativistic:** Charged particles (ions or electrons) which have speeds comparable to the speed of
1897 light.

1898 **RHESSI:** Reuven Ramaty High Energy Solar Spectroscopic Imager was a NASA solar flare
1899 observatory. It was launched on 5 February 2002 and was operational until 16 August 2018.
1900 RHESSI's primary mission was to explore the physics of particle acceleration and energy
1901 release in solar flares.

- 1902 **Ring Current:** In the magnetosphere, a region of current that flows near the geomagnetic equator
1903 in the outer belt of the two Van Allen radiation belts. The current is produced by the gradient
1904 and curvature drift of the trapped charged particles of energies of 10 to 300 keV. The ring
1905 current is greatly augmented during magnetic storms because of the hot plasma injected from
1906 the magnetotail and upwelling oxygen ions from the ionosphere. Further acceleration processes
1907 bring these ions and electrons up to ring current energies. The ring current (which is a diamagnetic
1908 current) causes a worldwide depression of the horizontal geomagnetic field during a magnetic
1909 storm.
- 1910 **SDO/EVE:** The Solar Dynamics Observatory is a NASA mission designed to understand the Sun's
1911 influence on the Earth and near-Earth space by studying the solar atmosphere on small scales
1912 of space and time in many wavelengths simultaneously. The EVE (Extreme Ultraviolet
1913 Variability Experiment) instrument measures the solar extreme ultraviolet (EUV) spectral
1914 irradiance at high spectral resolution, temporal cadence and precision.
- 1915 **Solar Energetic Particle (SEP):** An energetic particle of solar flare/interplanetary shock origin.
- 1916 **Sheath:** The plasma and magnetic fields in the downstream subsonic region after collisionless
1917 shocks. See Shock Wave.
- 1918 **Shock Wave:** A shock wave is characterized by a discontinuous change in pressure, density,
1919 temperature, and particle streaming velocity, propagating through a compressible fluid or
1920 plasma. Fast collisionless shock waves occur in the solar wind when fast solar wind overtakes
1921 slow solar wind with the difference in speeds being greater than the magnetosonic speed.
1922 Collisionless shock thicknesses are determined by the proton and electron gyroradii rather than
1923 the collision lengths. See also Diffusive Shock Acceleration; Solar Wind Shock.
- 1924 **Solar Corona:** See Corona.
- 1925 **Solar Cycle:** The approximately 11-year quasi-periodic variation in the sunspot number. The
1926 polarity pattern of the magnetic field reverses with each cycle. Other solar phenomena, such as
1927 the 10.7-cm solar radio emission, exhibit similar cyclical behavior. The solar magnetic field
1928 reverses each sunspot cycle so there is a corresponding 22-year solar magnetic cycle.
- 1929 **Solar Energetic Particle (SEP) Event:** A high flux event of solar (low energy) cosmic rays. This
1930 is commonly generated by larger solar flares or CME shocks, and lasts, typically from minutes
1931 to days. See also Cosmic Ray.

- 1932 **Solar Flares:** Transient perturbations of the solar atmosphere as measured by enhanced x-ray
 1933 emission (see x-ray flare class), typically associated with flares. Five standard terms are used
 1934 to describe the activity observed or expected within a 24-h period:
- 1935 Very low - x-ray events less than C-class.
 1936 Low - C-class x-ray events.
 1937 Moderate - isolated (one to 4) M-class x-ray events.
 1938 High - several (5 or more) M-class x-ray events, or isolated (one to 4).
 1939 M5 or greater x-ray events.
 1940 Very high - several (5 or more) M5 or greater x-ray events.
- 1941 **Solar Maximum:** The month(s) during the sunspot cycle when the smoothed sunspot number
 1942 reaches a maximum.
- 1943 **Solar Minimum:** The month(s) during the sunspot cycle when the smoothed sunspot number
 1944 reaches a minimum.
- 1945 **Solar Orbiter:** A European Space Agency-led (ESA) mission intended to perform detailed
 1946 measurements of the inner heliosphere and nascent solar wind to answer the question "How
 1947 does the Sun create and control the heliosphere?" The mission will make observations of the
 1948 Sun from an eccentric orbit moving as close as ~60 solar radii (R_S), or 0.284 astronomical units
 1949 (AU) from the Sun.
- 1950 **Solar Wind:** The outward flow of solar particles and magnetic fields from the Sun. Typically at 1
 1951 AU, solar wind velocities are 300-800 km/s and proton and electron densities of 3-7 per cubic
 1952 centimeter (roughly inversely correlated with velocity). The total intensity of the interplanetary
 1953 magnetic field is nominally 3-8 nT.
- 1954 **SORCE:** Solar Radiation and Climate Experiment. A NASA mission that measures
 1955 electromagnetic radiation produced by the Sun and the power per unit area of that energy on
 1956 the Earth's surface.
- 1957 **Space Weather:** Dynamic variations at the Sun, in interplanetary space, in the Earth's and
 1958 planetary magnetospheres, ionospheres and atmospheres associated with space phenomena.
- 1959 **Stratosphere:** That region of the Earth's atmosphere between the troposphere and the mesosphere.
 1960 It begins at an altitude of temperature minimum at approximately 13 km and defines a layer of
 1961 increasing temperature up to about 30 km.

1962 **Streamer:** A feature of the white light solar corona (seen in eclipse or with a coronagraph) that
1963 looks like a ray extending away from the Sun out to about 1 solar radius, having an arch-like
1964 base containing a cavity usually occupied by a prominence.

1965 **Substorm:** A substorm corresponds to an injection of charged particles from the magnetotail into
1966 the nightside magnetosphere. Plasma instabilities lead to the precipitation of the particles into
1967 the auroral zone ionosphere, producing intense aurorae. Potential drops along magnetic field
1968 lines lead to the acceleration of ~1 to 10 keV electrons with brilliant displays of aurora as the
1969 electrons impact the upper atmosphere. Enhanced ionospheric conductivity and externally
1970 imposed electric fields lead to the intensification of the auroral electrojets.

1971 **Sudden Impulse (SI):** An abrupt (10s of seconds) jump in the Earth's surface magnetic field. The
1972 positive sudden impulses (SI⁺s) are caused by fast forward shock impingement onto the
1973 magnetosphere.

1974 **Sunspot:** An area seen as a dark spot, in contrast with its surroundings, on the photosphere of the
1975 Sun. Sunspots are concentrations of magnetic flux, typically occurring in bipolar clusters or
1976 groups. They appear dark because they are cooler than the surrounding photosphere. Larger and
1977 darker sunspots sometimes are surrounded (completely or partially) by penumbrae. The dark
1978 centers are umbrae. The smallest, immature spots are sometimes called pores.

1979 **Supersubstorm:** Defined as an event with SML < -2500 nT. These auroral zone events appear to
1980 have different evolutionary properties than the standard (Akasofu, 1964) auroral substorms.

1981 **SWARM:** A European Space Agency (ESA) mission originally instrumented to study the Earth's
1982 magnetic field. The current goals are to study magnetospheric-ionospheric coupling and auroral
1983 Space Weather problems.

1984 **Telsa:** A unit of magnetic flux density (Weber/m²). A nano Tesla (nT) is 10⁻⁹ Teslas.

1985 **Termination Shock:** The shock wave in the solar wind which is caused by the abrupt deceleration
1986 of the solar wind as it runs into the local interstellar medium (LISM). It is thought to lie
1987 somewhere between 70 and 150 AU from the Sun.

1988 **Thermal Speed** (ion, electron): The random velocity of a particle associated with its temperature.

1989 **Thermosphere:** That region of the Earth's atmosphere where the neutral temperature increases
1990 with height. It begins above the mesosphere at about 80-85 km and extends upward to the
1991 exosphere.

- 1992 **TIMED:** Thermosphere Ionosphere Mesosphere Energetics and Dynamics (TIMED). A NASA
 1993 mission to investigate and understand the energetics of the Earth's mesosphere and lower
 1994 thermosphere/ionosphere.
- 1995 **Total Electron Content (TEC):** The column density of electrons in the Earth's ionosphere.
- 1996 **Trapped Particle:** Particles gyrating about magnetic field lines (e.g., in the Earth's
 1997 magnetosphere). See also Magnetic Mirror and Pitch Angle.
- 1998 **Troposphere:** The lowest layer of the Earth's atmosphere, extending from the ground to the
 1999 stratosphere, approximately 13 km altitude. In the troposphere, temperature decreases with
 2000 height.
- 2001 **Ultraviolet (UV):** That part of the electromagnetic spectrum between 5 and 400 nm.
- 2002 **Ultra Low Frequency (ULF):** 1 milliHertz to 1 Hertz.
- 2003 **Very High Frequency (VHF):** That portion of the radio frequency spectrum from 3 to 300 MHz.
- 2004 **Very Low Frequency (VLF):** That portion of the radio frequency spectrum from 3 to 300 kHz.
- 2005 **Van Allen Radiation Belt:** See Radiation Belt.

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