

1 Negentropy anomaly analysis of the borehole strain associated with the Ms 8.0 2 Wenchuan earthquake

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6 **Abstract:** A large earthquake of 8.0 magnitude occurred on 12 May 2012, 14:28 UTC, with the
7 epicenter in Wenchuan. To investigate the pre-earthquake anomalous strain changes, negentropy is
8 introduced to borehole strain data at Guza station, approximated by skewness and kurtosis revealing
9 the non-Gaussianity of recorded fluctuations. We separate the negentropy anomalies from the
10 background by Otsu's method and accumulate the anomaly frequency in different scales. The results
11 show the long-scale cumulative frequency of negentropy anomalies follows a sigmoid behaviour,
12 while the inflection point of the fitting curve is close to the occurrence of the earthquake. For the
13 short-scale analysis before the earthquake, there are two cumulative acceleration phases
14 corresponding to the two crustal stress releases, indicating the preparation process of the Wenchuan
15 earthquake. We consider that negentropy exhibits potential for the analysis of earthquake precursor
16 anomalies.

17 **Keywords:** Negentropy anomaly, Otsu's thresholding, Cumulative frequency, Wenchuan
18 [earthquakeearthquake](#)

19 1. Introduction

20 Changes in crustal deformation fields over time have preceded at least some large earthquakes
21 (Thatcher, W. et al., 1981), such as the 2011 Tohoku earthquake (Hitoshi Hirose, 2011) and the
22 Ruisui earthquake in Taiwan in 2013 (Canitano A. et al., 2015). Borehole strain data, which record
23 the direct crustal changes, provide an opportunity to investigate preparation process prior to
24 earthquakes (Linde et al., 1996, Hsu, Y.-J. et al. 2015).

25 Various methods are used in identifying borehole strain anomalies based on large amount of
26 monitoring data. Experienced scholars extract borehole strain anomalies by discriminating patterns
27 of waveform behaviors compared to those during the normal stage ([Johnston M.J.S.](#) ~~Johnston~~ et al.,
28 2006, Chi S. L. et al., 2014). In the time domain, Qiu Z. H. et al. (2010) identified abnormal strain

29 changes by overrun rate and wavelet decomposition for the Wenchuan earthquake. While in the
30 frequency domain, Qi L. et al. (2011) thought the signal with a period of 10 to 60 minutes might be
31 anomalies through S-transform compared with the background signal. In addition, statistical methods
32 are proved effective in distinguishing borehole strain anomalies with regard to large earthquakes,
33 such as principal component analysis (Zhu K. G. et al., 2018) and correlation coefficients along with
34 the consistency relation (Kong X. et al., 2018).

35 The probability distribution function (PDF) of observation data is also an informative way of
36 extracting potential anomalies contained in earthquake generation processes. P. Manshour et al.
37 (2009) extracted variance anomalies of the probability density of the Earth's vertical velocity
38 increments, and successfully found a pronounced transition from Gaussian to non-Gaussian prior to
39 12 moderate and large earthquakes. Before the Wenchuan earthquake, the high-frequency fluid
40 observational data deviated from Gaussian distributions at 16 water level and 14 water temperature
41 stations (Sun X. L. et al., 2016).

42 Rather than the whole PDF, often its moments are utilized, moments may be estimated quite
43 reliably from relatively small amounts of data (Sattin, F. et al., 2009). In 2016, Chen H. J. et al.
44 applied skewness and kurtosis (the third- and fourth-order moments) of the geoelectric data to pick
45 up non-Gaussian anomalies to predict impending large earthquakes in Taiwan. On the other hand, for
46 turbulent or disordered systems, the non-Gaussian distribution of time series in skewness-kurtosis
47 domain attracts attention. Observation data series from various fields of geophysics indicate that a
48 parabolic relation between skewness and kurtosis holds in fields such as seismology (M. Cristelli, et
49 al., 2012), oceanography (Sura, P. et al., 2008) and atmospheric science (A. Maurizi, 2006).

50 Hence, it is implied that possible precursor anomalies lead to an increase in disordered
51 components of observation data deviate from Gaussian distribution during earthquake preparation
52 processes. K. Eftaxias et al. (2008) proved that the pre-catastrophic stage could break the persistency
53 and high organization of the electromagnetic field through studying fractional-Brownian-motion-type
54 model using laboratory and field experimental electromagnetic data. In view of Lévy flight and
55 Gaussian processes, Lévy flight mechanism prevents the organization of the critical state to be
56 completed before earthquakes, since the long scales are cut-off due to the Gaussian mechanism (S.M.
57 Potirakis et al., 2019).

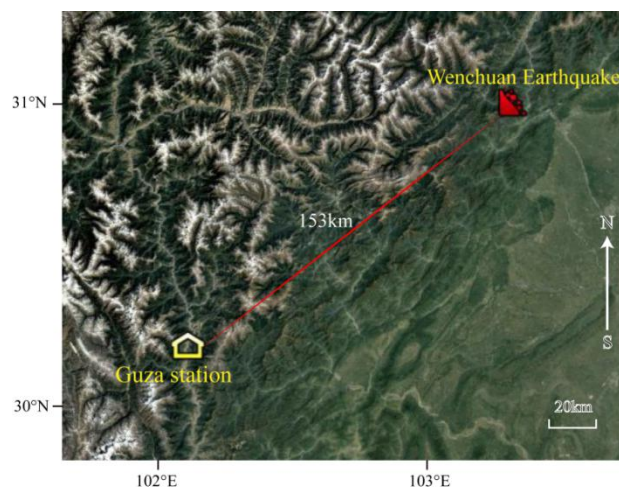
58 Entropy can serve as a measure of the unknown external energy flow into the seismic system
59 (Akopian, S. T., 2014). K. Karamanos et al. (2006, 2005) quantified and visualized temporal changes
60 of the complexity by approximate entropy, they claimed significant complexity decrease and
61 accession at the tail of the preseismic electromagnetic emission could be diagnostic tools for the
62 impending earthquake. Yukio Ohsawa (2018) detected earthquake activation precursors by studying
63 the regional seismic information entropy on earthquake catalog. Angelo De Santis et al. (2011)
64 recalled the Gutenberg - Richter law and considered the negative logarithm of b-value is the entropy
65 of the magnitude frequency of earthquake occurrence associated with two earthquakes in Italy.

66 Negentropy definition is based on the entropy and it is also widely used to detect non-Gaussian
67 features. Yue Li et al. (2018) proposed an arrival-time picking method based on negentropy for
68 microseismic data.—In this study, the negentropy is applied to borehole strain at Guza station
69 associated with the Wenchuan earthquake, approximated by skewness and kurtosis revealing the
70 non-Gaussianity of borehole fluctuations. Subsequently we study the extracted negentropy anomalies
71 in different scales to investigate correlations with crustal deformation.

72 2. Observation

73 YRY-4 borehole strainmeters, which are designed to record continuous deformation occurring
74 over periods of minutes to years, have been deployed at depths of dozens of metres at more than 40
75 terrain-sensitive locations within China. These strainmeters are capable of resolving strain changes
76 of less than one-billionth. The data sampling rate is once per minute.

77 The study period is from January 1, 2007, to June 30, 2009. The object of the study is the
78 Wenchuan Ms 8.0 earthquake in the region, which is shown in Fig. 1.



80 Fig. 1. Location map showing the epicentre of the Wenchuan earthquake epicentre and the Guza station. The
 81 Wenchuan earthquake occurred at 14:28:04 on May 12, 2008 (UTC+8). The magnitude of the earthquake was Ms
 82 8.0, and the focal depth was approximately 14 km. The epicentre was located in Wenchuan County, Sichuan
 83 Province, at 31.01°N, 103.42°E.

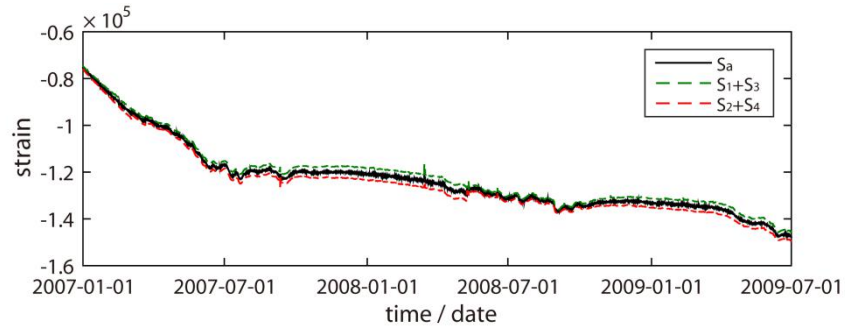
84 Because the four gauges of the YRY-4 borehole strainmeter are arranged at 45° intervals, this
 85 design has improved its self-consistency. This arrangement produces four observation values: S_i ,
 86 ($i=1, 2, 3, 4$) (Qiu et al., 2013a). The self-consistency as shown in equation (1), which can be used to
 87 test the reliability of the data among the four gauges.

$$S_1 + S_3 = S_2 + S_4 \quad (1)$$

89 In practical application, the higher the correlation between both sides of the equation (1), the
 90 more reliable the data. Generally, we use S_a for the areal strain in describing the subsurface strain
 91 state of the observation area as shown in equation 2 (Qiu et al., 2013a):

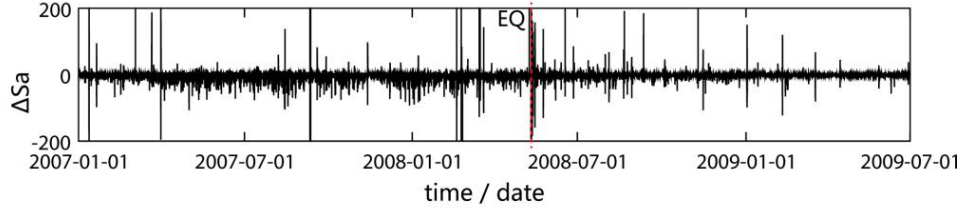
$$S_a = (S_1 + S_2 + S_3 + S_4) / 2 \quad (2)$$

93 The borehole strain is highly consistent among the four gauges at the Guza station (Qiu, et al.,
 94 2009), as shown in Fig. 2.



95
 96 Fig. 2. Self-consistency of the borehole strain at Guza from January 1, 2007, to June 30, 2009.

97 We first calculate the difference in the data because in the borehole strain data, the change in the
 98 strain is a concern (Li Jinwu et al., 2014). We then remove the components associated with the solid
 99 tide frequencies through a daily harmonic analysis. The remaining high-frequency signals are shown
 100 in Fig. 3.



101

102 Fig. 3. High-frequency areal strain after removing harmonic information at Guza from January 1, 2007 to June 30,
 103 2009.

104 3. Methodology

105 3.1 Negentropy and non-Gaussianity

106 The entropy-based negentropy is a statistically justified measure of non-Gaussianity (A.
 107 Hyvarinen, et al., 2000). The entropy of a random variable $X = \{x_1, x_2, \dots, x_i, \dots\}$ is defined as

$$108 \quad H(X) = -\sum_i P(X = x_i) \log P(X = x_i) \quad (3)$$

109 where P is the probability density function. Entropy measures the randomness of a random variable.
 110 The Gaussian random variable has the largest entropy of all other random variables with equal
 111 variance (T.M. Cover et al., 1991). The definition of negentropy is given by

$$112 \quad J(X) = H(X_{gauss}) - H(X) \quad (4)$$

113 in which X_{gauss} is a Gaussian random variable with the same mean and covariance matrix as X . The
 114 entropy of a Gaussian random variable can be estimated by

$$115 \quad H(X_{gauss}) = \frac{1}{2} \log |\det \Sigma| + \frac{n}{2} (1 + \log 2\pi) \quad (5)$$

116 where n is the dimension of the variable, and Σ is its covariance matrix.

117 However, the theoretical calculation of negentropy also depends on the prior probability density
 118 of random variables and other information that is difficult to determine accurately. In practical
 119 applications, higher order statistics (HOS) and density polynomial expansion are usually used to
 120 approximate one-dimensional negentropy (Jones and Sibson, 1987). The approximation results are as
 121 follows:

$$122 \quad J(x) \approx \frac{1}{12} skewness^2(X) + \frac{1}{48} kurtosis^2(X) \quad (6)$$

123 This definition suggests that any deviation from a Gaussian distribution will increase the
 124 negentropy $J(x)$. The skewness and kurtosis are the third- and fourth-order statistics, respectively,
 125 which are defined as

$$126 \quad \text{skewness}(X) = \frac{\mu_3}{\sigma^3} = \frac{E[(X - \mu)^3]}{E[(X - \mu)^2]^{3/2}} \quad (7)$$

127 and

$$128 \quad \text{kurtosis}(X) = \frac{\mu_4}{\sigma^4} = \frac{E[(X - \mu)^4]}{E[(X - \mu)^2]^2} - 3 \quad (8)$$

129 where μ is the mean of X and σ is the standard deviation of X . Skewness is a measure of
 130 asymmetry in a PDF. A symmetric distribution has zero skewness. Kurtosis is a measure of the
 131 heaviness of the tails. Distributions that are more outlier-prone than the normal distribution have
 132 kurtosis values greater than zero.

133 Moreover, the relation between the skewness and kurtosis is universal, they approximately align
 134 along a quadratic curve (Sattin, F., et al., 2009):

$$135 \quad \text{kurtosis}(X) = A \cdot \text{skewness}^2(X) + B. \quad (9)$$

136 Here we calculate the normalized skewness and kurtosis in the study period, so equation (9) can
 137 be derived into

$$138 \quad \underline{\text{kurtosis}(X) = A \cdot (\text{skewness}^2(X) - 1)} \quad (10)$$

139 indicating the test day is super-Gaussian when the skewness is outside the range (-1,1).

140 This relation is trivial in a Gaussian fluctuating system; it reduces to a fixed mass around zero
 141 (skewness= 0 and kurtosis= 0). In a turbulent environment where fluctuating quantities obey
 142 non-Gaussian statistics, the moments obey the above relation.

143 3.2 Otsu's thresholding method

144 To solve the negentropy anomaly detection problem, we designed a simple thresholding
 145 hypothesis test using the Otsu method (Otsu, 1979) that provides an optimal separation between
 146 background and seismic-related activities. For any given value k , we can separate the previously
 147 calculated $J(x)$, as shown in equation (6), into the following two classes:

148

$$\begin{aligned}
C_0(k) &= \{J(x) \leq k\}, \\
C_1(k) &= \{J(x) > k\}.
\end{aligned}
\tag{10}$$

149 Using these classes, the weighted average value $\mu_T(x)$ of $J(x)$ can be expressed as follows:

150

$$\begin{aligned}
\mu_T(x) &= \lambda_0(k)\mu_0(x; k) + \lambda_1(k)\mu_1(x; k), \\
\lambda_0(k) + \lambda_1(k) &= 1.
\end{aligned}
\tag{11}$$

151 where $\mu_0(x; k), \mu_1(x; k)$ are the mean values of the class $C_i(k)$, $i=1, 2$, and $\lambda_i(k)$ is the
152 percentage of points belonging into each class. Following the thresholding scheme of Otsu (1979),
153 we define the following cost function:

154

$$\begin{aligned}
\sigma_B^2(k) &= \lambda_0(k)(\mu_0(x; k) - \mu_T(x; k))^2 + \lambda_1(k)(\mu_1(x; k) - \mu_T(x; k))^2 \\
&= \lambda_0(k)\lambda_1(k)(\mu_1(x; k) - \mu_0(x; k))^2
\end{aligned}
\tag{12}$$

155 where σ_B^2 is the within-class variance of negentropy. Then, by finding the k^* value searching for k
156 when σ_B^2 becomes the maximum

157

$$k^* = \arg \max_k \sigma_B^2(k) \tag{13}$$

158 the optimal value k^* here separates the background set and anomaly set.

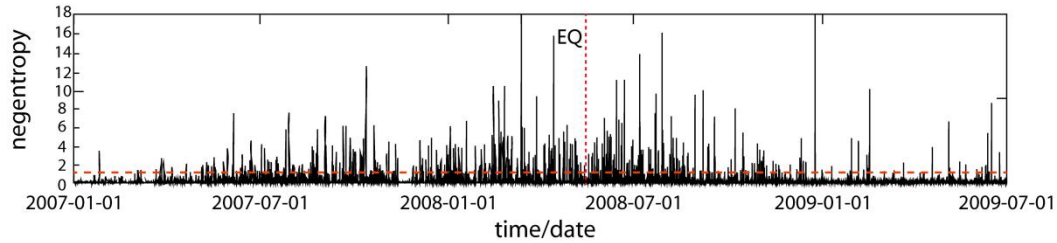
159 In this test, our initial assumption is that the sliding window is composed of a Gaussian signal of
160 non-seismic-related activities. When our test negentropy exceeds the critical value k^* , this initial
161 hypothesis is not valid, and the alternative is true, indicating the presence of a negentropy anomaly
162 within the window.

163 4. Results

164 According to the empirical hypothesis that geophysical signals deviate from the Gaussian
165 distribution when they record seismic-related activities, and based on the results of previous studies,
166 we conduct the following investigation.

167 4.1 Extracting negentropy anomalies

168 As the negentropy is calculated using a 2-hour long sliding window, we assume that it reaches
169 the maximum values when the time window contains anomalies from seismic-related activities. The
170 negentropy during the study period is shown in Fig. 4.



171

172 Fig. 4. Negentropy at Guza from January 1, 2007, to June 30, 2009. The red dotted, horizontal line is the optimal
 173 threshold k^* calculated by Otsu's method.

174

The within-class variance σ_B^2 and negentropy value distribution are compared in Fig. 5.

175

According to equations (10) to (13), when $k^*=1.1130$, σ_B^2 reaches its maximum ~~maximum~~.

176

Therefore, the negentropies were separated by k^* into the quasi-Gaussian background and

177

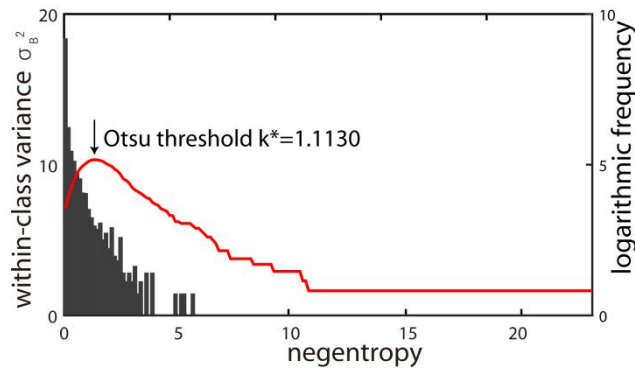
non-Gaussian anomalies from 2007 to 2009. Otsu threshold k^* here is consistent with the accuracy of

178

the negentropy and the strain data. The YRY-4 borehole strainmeter has a measurement accuracy of

179

10^{-9} , however, we usually cutoff four digits after the decimal point in practical calculations.



180

181 Fig. 5. Within-class variance σ_B^2 of the negentropy (red line) and negentropy histogram.

182

In the skewness-kurtosis domain, the statistical relationship of the borehole areal strain is

183

consistent with parabolic behaviour as described in equation (9)(Fig. 6(a)), verifying that the

184

turbulent system of borehole strain is significantly non-Gaussian during the study period. However,

185

the extracted negentropy anomalies are clustered strongly on the left side of the parabola, which

186

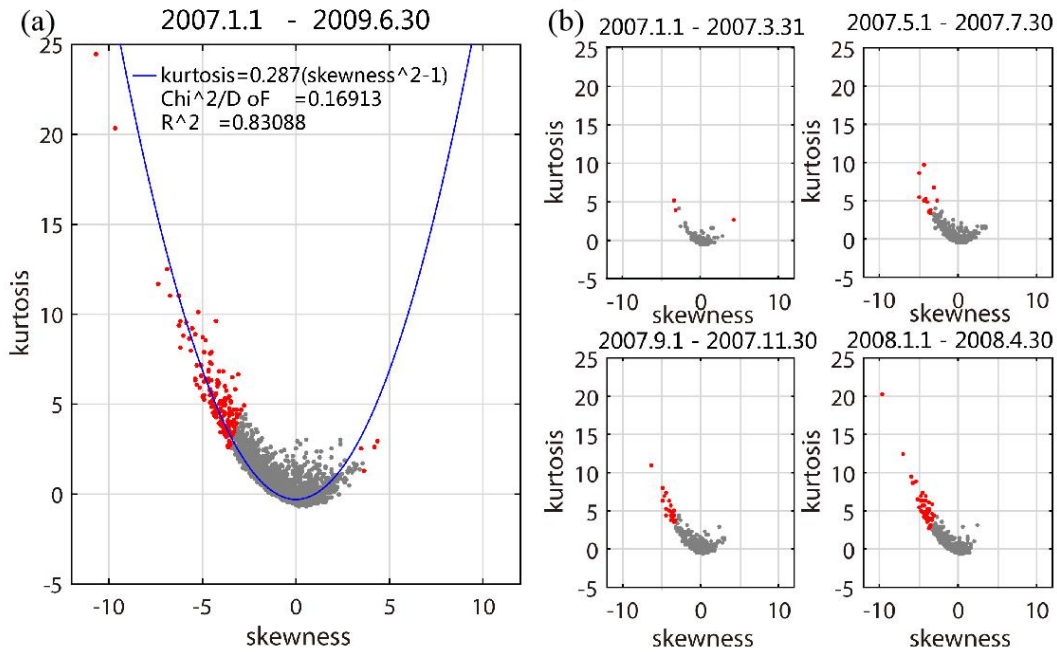
could be a signature of crustal deformation related to the earthquake. Here, there are four points on

187

the right side; one occurred in early 2007, and the others occurred after the earthquake. Therefore,

188

we will not discuss them in the following.



189

190 Fig. 6. Negentropy distributions in the skewness-kurtosis domain in (a) January 1 ,2007, to June 30 2009, and (b)

191 four shorter periods before the earthquake. Red denotes that the negentropy is greater than k^* , and grey indicates

192 that the negentropy less than k^* . The blue curve is the quadratic fit with the a 95% confidence.

193 In addition, as shown in Fig. 6(b), at times far from the earthquake, the negentropy distribution
 194 is basically Gaussian in the skewness-kurtosis domain. However, at times closer to the earthquake,
 195 the distribution gradually begins to show non-Gaussianity, with more negentropy anomalies
 196 appearing on the left side of the parabola. While in 2008, almost all of the negentropy present
 197 left-skewed.

198 These phenomena prompt us to study the origin of this left-skewed distribution and its possible
 199 correspondence with the seismogenic process.

200 4.2 Negentropy anomaly frequency accumulation

201 The transition of negentropy anomalies in the skewness-kurtosis domain is quantified as the
 202 change of the anomaly frequency per unit time through a logarithmic-linear model.
 203 Logarithmic-linear models are often used of interest to estimate the expected frequency of the
 204 response variable at the original scale for a new set of covariate values, such as in the famous
 205 Gutenberg-Richer law, in which a linear relationship exists between the logarithm of the cumulative
 206 number of seismic events of magnitude M or greater versus the magnitude M (Gutenberg and Richer,
 207 1954).

208 The logarithmic-linear regression model is proposed as

209
$$\log N = \beta_1 \times k_j + \beta_0 + \varepsilon \quad (14)$$

210 where k_j takes different threshold values according to the $J(x)$ values, N is the number of
 211 occurrences in which J is greater than or equal to the threshold k_j , β_0 and β_1 are the regression
 212 coefficients, where a lower slope β_1 indicates that there are more higher J values, implying there
 213 are more anomalies at that moment, and ε is the random error that represents the model uncertainty.

214 We use the logarithmic-linear model to solve the relationship between the negentropy anomaly
 215 frequency and different thresholds each day using the ordinary least squares method (OLS) method.
 216 Afterwards, an optimal threshold k^* , calculated by the Ostu method, is chosen for all models, where

217
$$N_j(t) = \exp(\beta_1(t) \times k^* + \beta_0(t) + \varepsilon(t)) \quad (15)$$

218 and the $N_j(t)$ under the threshold k^* is shown in Fig. 7. The model theoretically solves the
 219 problem of selecting the length of the time window. In addition, the estimated $N_j(t)$ is considered
 220 as the expected frequency of anomalies.

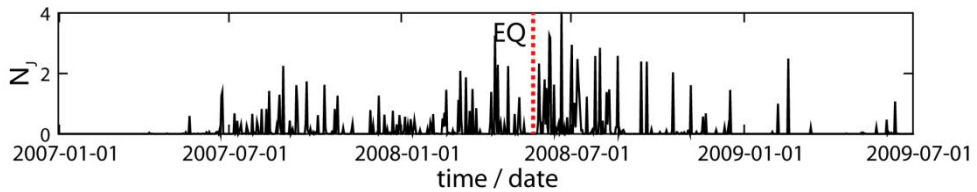
221 The goodness of fit for each logarithmic-linear model was evaluated using analysis of

222
$$R^2 = 1 - \frac{\sum_{i=1}^n (N_i - \hat{N}_i)^2}{\sum_{i=1}^n (N_i - \bar{N}_i)^2} \quad (16)$$

223 and the root-mean-squared error

224
$$RMSE = \sqrt{\sum_{i=1}^n (N_i - \hat{N}_i)^2 / n} \quad (17)$$

225 The R^2 and $RMSE$ values in the study period (912 days) show that the logarithmic-linear
 226 relationship can explain the relationship between the negentropy anomaly frequency and different
 227 thresholds. The mean of R^2 is 0.9695, which is close to 1, and the variance of R^2 is 0.0435. The
 228 mean and variance of the $RMSE$ are also small (0.1098 and 0.1301, respectively).

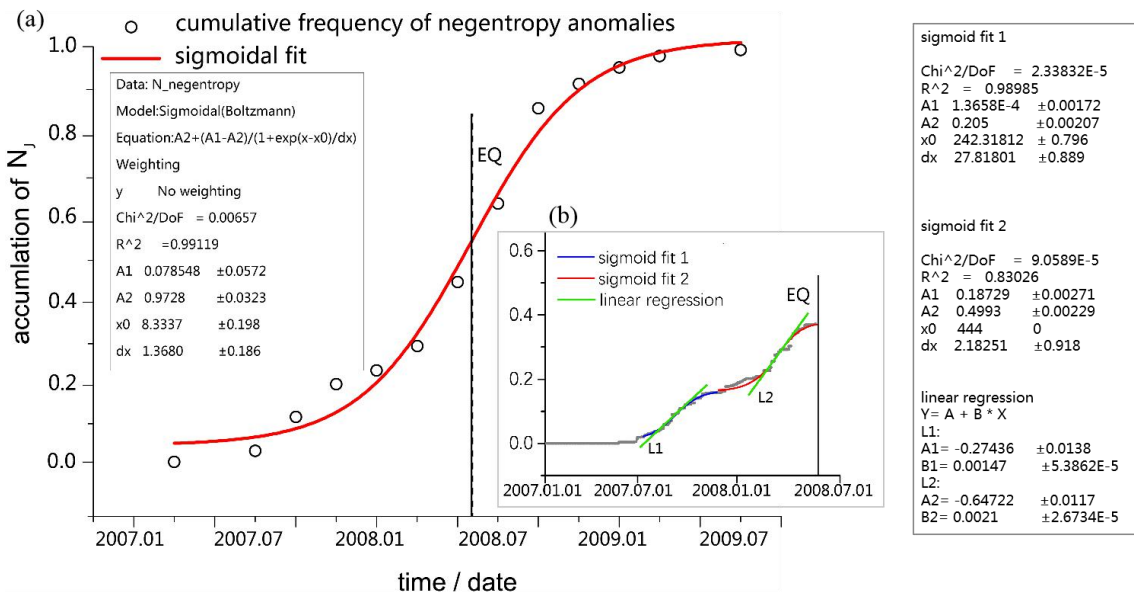


229
 230 Fig. 7. Estimated expected frequency N_j under the optimal threshold k^*

231 In general, accumulated value of a typical random process usually has an linear increase. The
 232 negentropy cumulative frequency of the study period is shown in Fig. 8. We not only do a long-scale

233 analysis of the whole earthquake process, but also carry out a short-scale analysis of the
 234 pre-earthquake process.

235 For the entire earthquake process, a two-month long sliding window is selected for
 236 accumulation as shown in Fig. 8(a). Beginning in July 2007, the negentropy anomalies gradually
 237 accumulated. In particular, we find more frequent negentropy anomalies in 2008 as the earthquake
 238 approaches, and less frequent anomalies after, so a sigmoid function is used to fit the acceleration,
 239 before the earthquake and the deceleration after. According to De Santis, A. et al. (2017), inflection
 240 point in this function is a reasonable estimation of the time of the significant change in the critical
 241 dynamical system. Our calculation shows that the inflection point x_0 of the optimal fitting is 8.3337,
 242 which is surprisingly close to the actual time (8.3871) of the Wenchuan earthquake after conversion.
 243 Thus, the earthquake moment is proved to be a critical time during the whole Wenchuan earthquake
 244 process.



245
 246 Fig. 8. (a) Results of the long-scale negentropy anomaly frequency for the Wenchuan earthquake at Guza station
 247 from January 1, 2007, to June 30, 2009. Each circle represents an anomaly negentropy for 2 months. The
 248 cumulative frequency of negentropy anomaly is represented. The earthquake day is represented as a vertical solid
 249 line. The red line is a sigmoidal fit that underlines a critical point (vertical dashed line) is close to the occurrence of
 250 the earthquake. (b) Results of the short-scale negentropy anomaly frequency prior to the earthquake, every grey
 251 point is an anomaly for one day, the blue and red lines are two segment sigmoid fit results. Two green lines
 252 represent the liner regression for the two phases, the first phase slope is 0.00147, the second one is 0.0021.

253 When we narrowed the accumulated window to one day, we observed two negentropy anomalies
254 before the earthquake as shown in Fig. 8(b). The first anomaly frequency increase occurred from
255 August to October 2007. In March 2008, there was a second phase of anomaly increase, and the
256 cumulative frequency then slowly increased to a plateau period near the time of the earthquake.
257 These two phases prior to the earthquake are also approximated with sigmoid functions. In order to
258 further compare the anomalies of the two phases, we use linear regression to fit the central part of the
259 two sigmoid curves. We find that the second acceleration is greater than the first acceleration. We
260 think the accumulations of these two negentropy anomalies may be an indication of crustal activity
261 before the Wenchuan earthquake.

262 5. Discussion and Conclusion

263 In the skewness-kurtosis domain, we observed the evolution of the negentropy distribution prior
264 to the earthquake. Negentropy gradually transformed its distribution to a parabolic relation since July
265 2007, indicating a relatively stable state was broken due to the non-Gaussian mechanism.

266 Previous studies for the Wenchuan earthquake are consistent with our findings. Wang (2018)
267 concluded that an apparent stress change occurred after June 2007 based on multiple focal
268 mechanisms. Likewise, we did not find negentropy anomalies in the first six months of 2007. In the
269 large-scale analysis, we show that the cumulative frequency of negentropy anomalies follows the a
270 power-law behaviour approaching a critical time that is close to the earthquake time, and then
271 recovers as a typical recovery phase after the earthquake This process is consistent with the empirical
272 phenomena before and after earthquakes, which is also similar with a potential earthquake precursory
273 pattern in magnetic data from Swarm satellites by A. De Santis (2017) for the 2015 Nepal event.

274 Qiu (2009) and Chi (2014) speculated that the areal strain indicates that the integrity of the
275 medium around the borehole at the Guza station began to change significantly after August 2007,
276 because the continuity of the medium in the source region of the Wenchuan earthquake was
277 gradually deformed during the nucleation process. In our short-scale analysis, negentropy anomalies
278 also present a acceleration in August 2007. The second acceleration in March 2008 is also consistent
279 with the occurrence of a phase measured by GPS (Jiang Z. S., 2009), in which the elastic
280 deformation of the crust reaches its limit, and the deformation is resisted in the seismogenic region
281 before the earthquake.

282 More importantly, we speculate that the two accelerations of the cumulative negentropy
283 anomaly corresponding to two stress releases. Ma Jin (2014) proposed the deformation
284 characteristics in the sub-instability stage of faults before earthquakes based on different experiments
285 and several earthquakes. Fault zones contain relatively weak and relatively strong parts. The former
286 is the area where strain release begins, while the latter is the stress locking part and the beginning of
287 rapid instability (Noda et al., 2013). She thinks accelerated expansion of the strain release zone in
288 fault zone is a sign of entering the inevitable earthquake stage. There are two instabilities before
289 earthquakes, the former is related to the release of weak parts, and the latter is related to the rapid
290 release of strong parts of the fault during strong earthquakes. The accelerated expansion of the
291 former promotes the occurrence of the latter.

292 We speculate that the two negentropy anomaly accelerations may represent the two instabilities
293 associated with the Wenchuan earthquake. The first corresponds to the start of strain release and the
294 second larger one corresponds to the acceleration of instability, indicating that strong earthquakes are
295 likely to occur.

296 In our work, the extracted negentropy anomalies of the short-period signal of borehole areal
297 strain based on Otsu's thresholding associated with the Wenchuan earthquake are analyzed. A
298 logarithmic-linear model is proposed to estimate the expected frequency of the left-skewed
299 negentropy anomaly for everyday, and the evolution processes of the negentropy anomaly frequency
300 are studied in both long- and short-scale during the study period. We consider the negentropy
301 anomalies corresponding to crustal stress changes, which may be a reflection of the subsurface
302 medium and faults activities in the focal area associate with the Wenchuan earthquake. Moreover, we
303 may be able to ensure that the negentropy has great potential in the study of earthquake precursors.

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