

Quasi-periodic solutions associated with the extended nonlinear feedback loop of the five-dimensional nondissipative Lorenz model

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On quasi-periodic solutions associated with the extended nonlinear feedback loop of the five-dimensional nondissipative Lorenz model

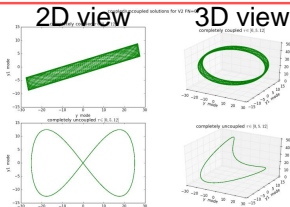
Sara Faghiih-Naini, San Diego State University

Quasi-periodic solutions associated with the extended nonlinear feedback loop of the five-dimensional nondissipative Lorenz model

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Objective

- Analysis and comparison of 5D NLM
- Fundamental role of nonlinear terms in producing quasi-periodic solutions
- Collective role of the nonlinear feedback loop and heating term in producing quasi-periodic solutions



Y-Y1-plots of periodic/quasiperiodic solutions for $t_{end}=0.512$ (top) and $t_{end}=5.12$ (bottom) of coupled (left) and uncoupled (right) LL 5D NLM

Approach

- Performing analytical, symbolic (python scipy) and numerical (Python ODE solver (odeint)) analysis of 5D NLM
- Analyzing quasiperiodic solutions of 5D NLM
- Comparing to solutions of 5D NLM, locally linearized 5D NLM, 3D NLM, ...
- Performing frequency analysis using Python fft package
- Visualizing quasiperiodic solutions using Python visualization packages

Key Milestones

- 09/30/2016: solve for eigenvalues of 5D NLM (analytical and symbolic computation)
- 11/01/2016: verify results, plots for comparison of 3D NLM, LL 3D NLM FN=0 and FN=1, 3D-eigenvaluesolution (analytical and symbolic); plots for comparison of 5D NLM, LL 5D NLM FN=0 and FN=1, 5D-eigenvaluesolution (analytical and symbolic)
- 11/20/2016: finish frequency analysis
- 11/22/2016: create slides for presentation
- 11/30/2016: finish oral presentation
- 12/06/2016: draft of paper

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Figure: QuadChart

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- ▶ Frequency ratio $\frac{\omega_2}{\omega_1}$ of two modes determines behavior of solution:
- ▶ If $\frac{\omega_2}{\omega_1}$ is rational, the motion is periodic and has a closed orbit.
- ▶ If $\frac{\omega_2}{\omega_1}$ is irrational, these two frequencies are called incommensurate;
- ▶ The composite motion is quasiperiodic and its period is infinite;
- ▶ The trajectory is dense, that means, it comes arbitrarily close to each point of torus.

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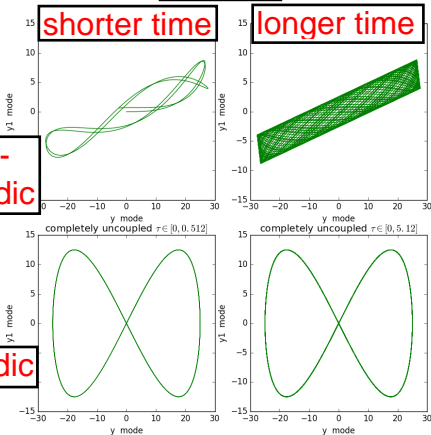
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2D view coupled solutions for V2 FN=0,

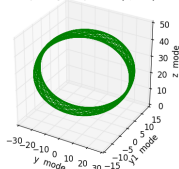
3D view

**quasi-
periodic**

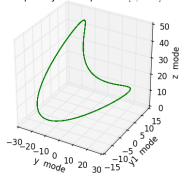


periodic

completely coupled $\tau \in [0, 5.12]$



completely uncoupled $\tau \in [0, 5.12]$



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Figure: Quasiperiodic solution of coupled LL 5D NLM with a frequency ratio of $\frac{1}{2} (3 - \sqrt{5})$ (top) and periodic solution of uncoupled LL 5D NLM with a frequency ratio of 2 (bottom) up to time 0.512 (left) and 5.12 (middle)

5D NLM V1

Removing the dissipative terms from the five dimensional Lorenz model (5D LM), the the nondissipative five dimensional Lorenz model (5D NLM) results:

$$\frac{dX}{d\tau} = \sigma X + \sigma Y, \quad (1)$$

$$\frac{dY}{d\tau} = -XZ + rX - Y, \quad (2)$$

$$\frac{dZ}{d\tau} = XY - XY_1 - bZ, \quad (3)$$

$$\frac{dY_1}{d\tau} = XZ - 2XZ_1 - d_0 Y_1, \quad (4)$$

$$\frac{dZ_1}{d\tau} = 2XY_1 - 4bZ_1. \quad (5)$$

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5D NLM V2

Perturbation method leads to the locally linearized 5D NLM (LL 5D NLM):

$$\frac{dX'}{d\tau} = \sigma Y', \quad (6)$$

$$\frac{dY'}{d\tau} = (r - Z_c)X' - X_c Z' - FN(X'Z'), \quad (7)$$

$$\frac{dZ'}{d\tau} = (Y_c - Y_{1c})X' + X_c Y' - X_c Y'_1 + FN(X'Y' - X'Y'_1), \quad (8)$$

$$\frac{dY'_1}{d\tau} = (Z_c - 2Z_{1c})X' + X_c Z' - 2X_c Z'_1 + FN(X'Z' - 2X'Z'_1), \quad (9)$$

$$\frac{dZ'_1}{d\tau} = 2Y_{1c}X' + 2X_c Y'_1 + 2FN(X'Y'_1). \quad (10)$$

Choosing $FN = 0$ makes the system linear with respect to the critical point and for $FN = 1$ the system is fully nonlinear.

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Model Assumptions and List of Simulations

Parameters: $\sigma = 10, r = 25, \Delta\tau = 0.001$

Method	Model	Equations	IC	Python Packages
analytical	V2 FN=0	6-10	$X_0 = \sqrt{\frac{5}{2}\sigma r},$ $X_c = \sqrt{\frac{5}{2}\sigma r + X_0^2}$	-
symbolic	V2 FN=0	6-10	$X_0 = \sqrt{\frac{5}{2}\sigma r},$ $X_c = \sqrt{\frac{5}{2}\sigma r + X_0^2}$	linalg
numerical	V2 FN=0 V2 FN=1 V1	6-10 6-10 1-5	$X_c = \sqrt{\frac{5}{2}\sigma r + X_0^2}$ for $X_0 = 0.25\sqrt{\frac{5}{2}\sigma r},$ $X_0 = 0.5\sqrt{\frac{5}{2}\sigma r},$ $X_0 = \sqrt{\frac{5}{2}\sigma r},$ $X_0 = 4\sqrt{\frac{5}{2}\sigma r},$	odeint

Impact of ICs

Table: Methods, corresponding model assumptions and computing packages used for verifying results and analyzing 5D NLM

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Analytical Solution for Eigenvalues

Setting $FN = 0$, plugging in the basic state values, the Jacobian matrix becomes:

$$A^{5DNLM} = \begin{pmatrix} 0 & \sigma & 0 & 0 & 0 \\ 0 & 0 & -X_c & 0 & 0 \\ 0 & X_c & 0 & -X_c & 0 \\ 0 & 0 & X_c & 0 & -2X_c \\ 0 & 0 & 0 & 2X_c & 0 \end{pmatrix}$$

To obtain the eigenvalues $\lambda_1, \dots, \lambda_5$, with $\lambda = i\beta$, the equation $\det(A^{5DNLM} - i\beta\mathbb{I}) = 0$ was solved applying Laplace's formula:

$$\beta_1 = \sqrt{(3 + \sqrt{5})} X_c$$

$$\beta_2 = -\sqrt{(3 + \sqrt{5})} X_c$$

$$\beta_3 = \sqrt{(3 - \sqrt{5})} X_c$$

$$\beta_4 = -\sqrt{(3 - \sqrt{5})} X_c$$

$$\beta_5 = 0$$

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For $k = 1, 2, 3, 4$ the form of the corresponding eigenvectors is

$$v_k = \begin{pmatrix} \sigma \left(\frac{-\beta_k^2}{2X_c^3} + \frac{5}{2X_c} \right) \\ i\beta_k \left(\frac{-\beta_k^2}{2X_c^3} + \frac{5}{2X_c} \right) \\ \frac{-\beta_k^2}{2X_c^2} + 2 \\ \frac{i\beta_k}{2X_c} \\ 1 \end{pmatrix} \text{ and } v_5 = \begin{pmatrix} \sigma \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}.$$

This leads to the general solution

$S(t) = C_1 e^{i\beta_1 t} v_1 + C_2 e^{i\beta_2 t} v_2 + C_3 e^{i\beta_3 t} v_3 + C_4 e^{i\beta_4 t} v_4 + C_5 e^{i\beta_5 t} v_5$, which can be expressed in the sine-cosine representation.

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- ▶ $C = C_1, \dots, C_5$ can be determined by applying the initial condition $I = (X_0 - X_c, 0, -r, 0, -\frac{r}{2})$.
- ▶ Using the analytical solutions of the eigenvectors, which include oscillatory modes with two incommensurate frequencies, as a basis for a solution, the coefficients of the eigenvectors can be determined as $C_1 + C_2 = \frac{-\beta_3^2 r}{2(\beta_3^2 - \beta_1^2)}$,
 $C_3 + C_4 = \frac{\beta_1^2 r}{2(\beta_3^2 - \beta_1^2)}$ and $C_5 = \frac{X_0 - X_c}{\sigma} + \frac{5r}{4X_c}$ (constant mode) (for more detailed discussions see Appendix)
- ▶ The **frequency ratio** $\frac{\beta_3}{\beta_1} = \frac{\sqrt{(3-\sqrt{5})X_c}}{\sqrt{(3+\sqrt{5})X_c}} = \frac{1}{2} (3 - \sqrt{5})$ is **irrational**, so solution is **quasiperiodic**.

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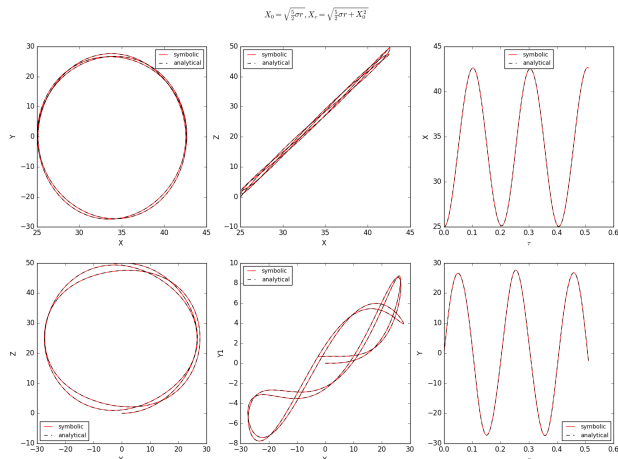
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Analytical vs. Symbolic Solution-shorter time



analytical and symbolic solutions are identical

Figure: Analytical vs. Symbolic Solutions for $X_0 = \sqrt{\frac{5}{2}\sigma r}, X_c = \sqrt{\frac{5}{2}\sigma r + X_0^2}$ and $\tau \in [0, 0.512]$

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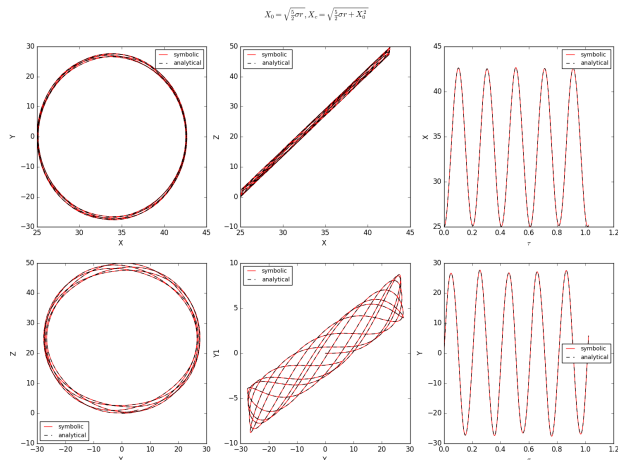
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Analytical vs. Symbolic Solution-longer time



analytical and symbolic solutions are identical

Figure: Analytical vs. Symbolic Solutions for $X_0 = \sqrt{\frac{5}{2}\sigma r}, X_c = \sqrt{\frac{5}{2}\sigma r + X_0^2}$ and $\tau \in [0, 1.024]$

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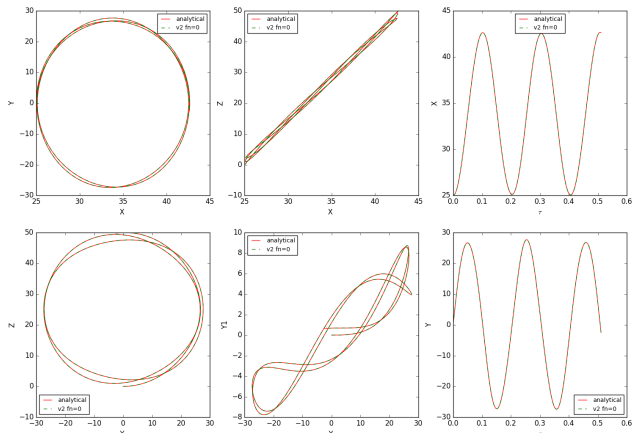
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$$X_0 = \sqrt{\frac{5}{2}\sigma r}, X_c = \sqrt{\frac{5}{2}\sigma r + X_0^2}$$



analytical and numerical solutions are identical

Figure: Analytical vs. Numerical solution of LL 5D NLM (V2, FN=0) for

$X_0 = \sqrt{\frac{5}{2}\sigma r}, X_c = \sqrt{\frac{5}{2}\sigma r + X_0^2}$ and $\tau \in [0, 0.512]$

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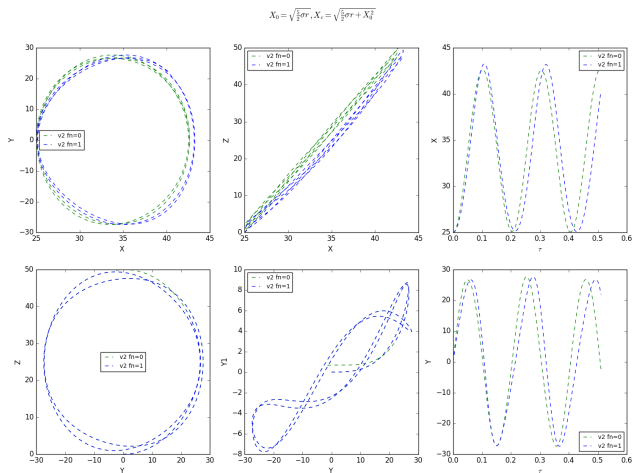


Figure: Numerical solutions of LL 5D NLM FN=0 vs. FN=1 for

$$X_0 = \sqrt{\frac{5}{2}\sigma r}, X_c = \sqrt{\frac{5}{2}\sigma r + X_0^2} \text{ and } \tau \in [0, 0.512]$$

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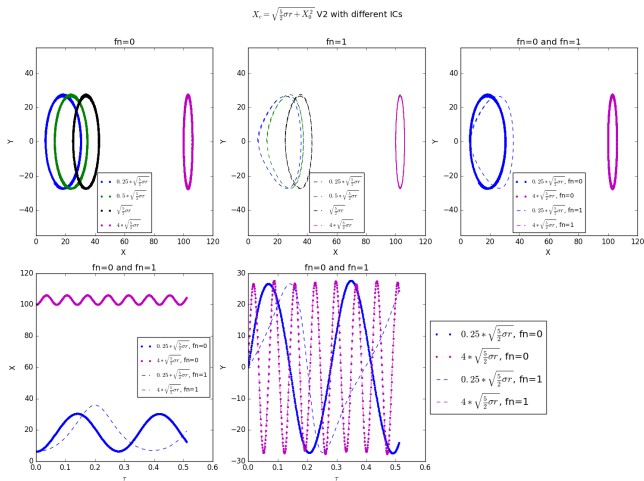


Figure: Comparison of different initial conditions in LL 5D NLM FN=0 and FN=1 for $t \in [0, 0.512]$

Examine coupling terms using 5D NLM V2

$$\frac{dX'}{d\tau} = \sigma Y',$$

$$\frac{dY'}{d\tau} = (r - Z_c)X' - X_c Z' - FN(X'Z'),$$

$$\frac{dZ'}{d\tau} = (Y_c - Y_{1c})X' + X_c Y' - \boxed{X_c Y'_1}^1 + FN(X'Y' - X'Y'_1),$$

$$\frac{dY'_1}{d\tau} = (Z_c - 2Z_{1c})X' + \boxed{X_c Z'}^2 - 2X_c Z'_1 + FN(X'Z' - 2X'Z'_1),$$

$$\frac{dZ'_1}{d\tau} = 2Y_{1c}X' + 2X_c Y'_1 + 2FN(X'Y'_1).$$

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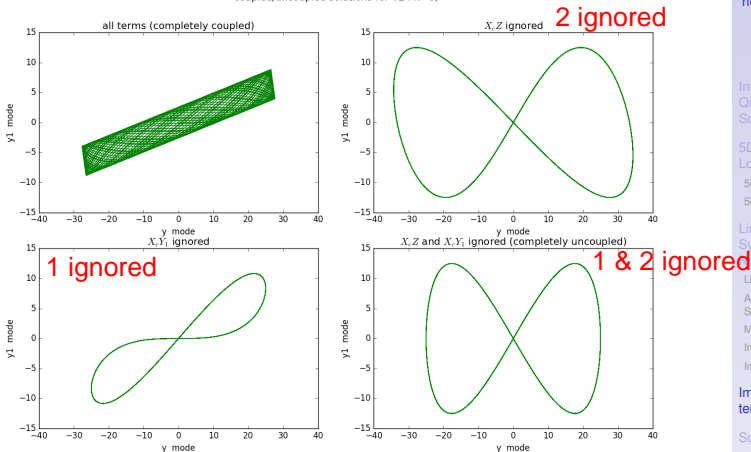


Figure: Solution of LL 5D NLM FN=0 becomes periodic, if coupling terms $X_c Y_1$ in $\frac{dz'}{dt}$ and/or $X_c Z$ in $\frac{dy_1'}{dt}$ are ignored

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$C = C_1, \dots, C_5$ can be determined by setting $t = 0$ and equating S with the initial condition $I = (X_0 - X_c, 0, -r, 0, -\frac{r}{2})$. The the system has the following form:

$$S(t) = \frac{-\beta_3^2 r}{2(\beta_3^2 - \beta_1^2)} \begin{pmatrix} \sigma \left(\frac{-\beta_1^2}{2X_c^3} + \frac{5}{2X_c} \right) \cos(\beta_1 t) \\ -\beta_1 \left(\frac{-\beta_1^2}{2X_c^3} + \frac{5}{2X_c} \right) \sin(\beta_1 t) \\ \left(\frac{-\beta_1^2}{2X_c^2} + 2 \right) \cos(\beta_1 t) \\ \frac{-\beta_1}{2X_c} \sin(\beta_1 t) \\ \cos(\beta_1 t) \end{pmatrix} + \frac{\beta_1^2 r}{2(\beta_3^2 - \beta_1^2)} \begin{pmatrix} \sigma \left(\frac{-\beta_3^2}{2X_c^3} + \frac{5}{2X_c} \right) \cos(\beta_3 t) \\ -\beta_3 \left(\frac{-\beta_3^2}{2X_c^3} + \frac{5}{2X_c} \right) \sin(\beta_3 t) \\ \left(\frac{-\beta_3^2}{2X_c^2} + 2 \right) \cos(\beta_3 t) \\ \frac{-\beta_3}{2X_c} \sin(\beta_3 t) \\ \cos(\beta_3 t) \end{pmatrix} + \left(\frac{X_0 - X_c}{\sigma} + \frac{5r}{4X_c} \right) \begin{pmatrix} \sigma \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

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```
from sympy import Symbol, Matrix

X_c = Symbol('X_c')
sig = Symbol('sig')
la = Symbol('la')
A5D=Matrix([[0, sig, 0, 0, 0],
            »      [0, 0, -X_c, 0, 0],
            »      [0, X_c, 0, -X_c, 0],
            »      [0, 0, X_c, 0, -2*X_c],
            »      [0, 0, 0, 2*X_c, 0]])
eigenvalues=A5D.eigenvals()
eigenvectors=A5D.eigenvects()
```

Figure: Python code for symbolic calculation of eigenvalues and eigenvectors

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Numerical Solution - 5D NLM V1

```
from scipy.integrate import odeint
import numpy as np

def Lorenz5D_nondissipative(state,t, sigma, r, beta, d0):

    x,y,z,y1,z1 = state

    dx = sigma * (y)
    dy = -x*z + r*x
    dz = x*y - x*y1
    dy1 = x*z - 2*x*z1
    dz1 = 2*x*y1

    return [dx, dy, dz, dy1, dz1]

sigma = 10.0
r = 25.0
beta = 8.0/3.0
d0 = 19.0/3.0

X0=np.sqrt(5./2*sigma*r)
xc=np.sqrt(5./2*sigma*r+X0**2)
zc=r
z1c = r/2
yc= 0.
y1c= 0.

dt=0.001
L=512*dt
t = np.arange(0.0,L, dt)
state = [X0_1, 0, 0, 0, 0]

out = odeint(Lorenz5D_nondissipative, state, t, args=(sigma, r, beta, d0))
```

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Figure: Python code for numerical solution of the 5D NLM using odeint

Numerical Solution - 5D NLM V2

```
from scipy.integrate import odeint
import numpy as np

def LL_Lorenz5D(state,t, parameters):

    x,y,z,y1,z1 = state
    sigma, r, beta, d0, FN, X0, xc, zc, yc, z1c, y1c = parameters

    dx = sigma * (y)
    dy = (r-zc)*x - xc*z - FN*(x*z)
    dz = (yc-y1c)*x + xc*y - xc*y1 + FN*(x*y-x*y1)
    dy1 = (zc-2*z1c)*x + xc * z - 2*xc*z1 + FN*(x*z-2*x*z1)
    dz1 = 2*y1c*x + 2*xc*y1 + 2*FN*(x*y1)

    return [dx, dy, dz, dy1, dz1]

sigma = 10.0
r = 25.0
beta = 8.0/3.0
d0 = 19.0/3.0

X0=np.sqrt(5./2*sigma*r)
xc=np.sqrt(5./2*sigma*r+X0**2)
zc=r
z1c = r/2
yc = 0.
y1c = 0.

dt=0.001
L=512*dt
t = np.arange(0.0,L, dt)
state = [X0,1-xc, 0, -r, 0, -r/2] #=[X', Y', Z', Y1', Y2']

FN=0
parameters = (sigma, r, beta, d0, FN, X0, xc, zc, yc, z1c, y1c)
out0 = odeint(LL_Lorenz5D, state, t, args=(parameters,))
FN=1
parameters = (sigma, r, beta, d0, FN, X0, xc, zc, yc, z1c, y1c)
out1 = odeint(LL_Lorenz5D, state, t, args=(parameters,))
```

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Figure: Python code for numerical solution of the LL 5D NLM using odeint.

Frequency analysis for LL 5D NLM

V2 frequency analysis $X_0 = \sqrt{\frac{2}{3}\sigma r}$, $X_1 = \sqrt{\frac{2}{3}\sigma r + X_0^2}$ magnitudes

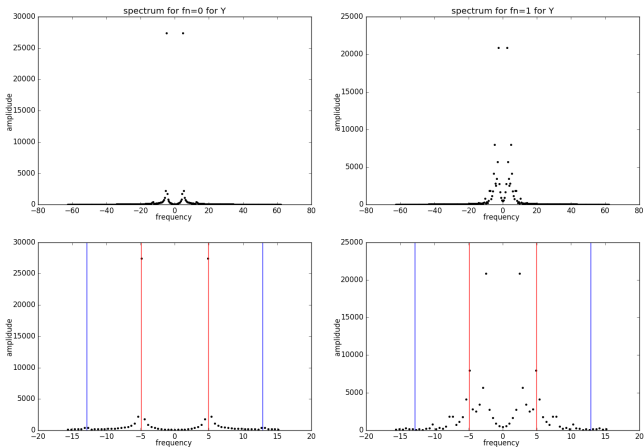


Figure: Magnitudes of fourier modes of LL 5D NLM FN=0 and FN=1 and zoomed-in presentation for comparison with eigenvalue solution for $\tau \in [0, 0.512]$

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