Influence of Atmospheric Stratification on the Integral Scale and Fractal Dimension of Turbulent Flows.

M. Tijera[1] , G. Maqueda[2], C. Yagüe[3]

[1] Applied Mathematics Dpt. (Biomathematics). Complutense University of Madrid, Madrid, Spain.

[2] Astronomy, Astrophysics and Atmospheric Science Dpt. Complutense University of Madrid, Madrid, Spain.

[3] Geophysics and Meteorology Dpt. Complutense University of Madrid, Madrid, Spain.

Correspondence to: M.Tijera (mtijera@fis.ucm.es)

Abstract

In this work the relation between integral scale and fractal dimension and the type of stratification in fully developed turbulence is analyzed. Integral scale corresponds to that in which energy from larger scales is incoming into turbulent regime. One of the aims of this study is the understanding of the relation between the integral scale and the Bulk Richardson number, which is one the most widely used indicators of stability close to the ground in atmospheric studies. This parameter will allow us to verify the influence of the degree of stratification over the integral scale of the turbulent flows in the Atmospheric Boundary Layer (ABL).The influence of the diurnal and night cycle in the relationship between the fractal dimension and integral scale is also analyzed. Fractal dimension of wind components is a turbulent flow characteristic as it has been shown in previous works, where its relation to stability was highlighted. Fractal dimension and integral scale of the horizontal (u`) and vertical (w`) velocity fluctuations have been calculated using the mean wind direction as framework. The scales are obtained using sonic anemometer data from three elevations 5.8 m, 13 m and 32 m above the ground measured during the SABLES- 98 field campaign. In order to estimate the integral scales a method that combines the normalized autocorrelation function and the best gaussian fit ($R^2 \geq 0.70$) has been developed. Finally, by comparing,

at the same height, the scales of u` and w` velocity components it is found that almost always the turbulent flows are anisotropic.

1 Introduction

The aim of this paper is to investigate the possible correlations between the integral scale of the turbulent stratified flows in the Atmospheric Boundary Layer and parameters charactering topological features of the wind velocity field, such as fractal dimension and its stability properties, studied through the Bulk Richardson number. We are aware that there is a lack of investigations between the integral scale and fractal dimension. The size of the integral scale of the horizontal and vertical components and fractal dimension of wind velocity near the earth's surface in boundary layer are determined. Also, these magnitudes are compared between them and versus other parameters such as the Bulk Richardson number. It is assumed that the turbulence is the primary agent that causes changes in the boundary layer. In turbulent flows it is observed that time series of meteorological variables as wind velocity, temperature, pressure and other atmospheric mechanical magnitudes fluctuate in a disordered way with peaks extremely sharp and irregular space and time variations. The complicated nature of these series indicates that the motion of the air is turbulent. If we take a good look at the variety of fluctuations of different periods and amplitudes observed in them we could explain the complicated structure of turbulence. The irregularity of the time series obey to the existence of different size and time scales and also to the nonlinear transfer of energy that exists between them in the turbulent flows (Monin and Yaglom, 1971).

The irregular behavior of these flows is also due to waves and turbulence that are often superimposed on a mean wind (Stull, 1998). If we filter the mean wind and waves in the appropriate range we will only have turbulence. Some previous works present results about this procedure (Tijera et al., 2008). In this work, the series of wind velocities in the three directions x, y and z recorded by the anemometer are divided in series of non-overlapping five minutes length. Each of these series applies the necessary rotations to get the x-axis in the mean wind direction (mean v is zero) and zero mean vertical velocity (w vertical component) (Kaimal and Finningan, 1994). We filter horizontal and vertical mean wind velocity obtaining the time series of fluctuations of the velocity in both directions ($u' = u - \bar{u}$, $w' = w - \bar{w}$).

When we observe these time series such as wind velocity, they vary in an irregular shape and in spite of their complexity presents a self-similarity structure (Frisch,1995) . This is a common property of the fractals, so that wind velocity could be considered as a fractal magnitude. The modern physical notion of fractals is largely known due to Mandelbrot (1977, 1985), but the mathematical notion of curves lines or sets having noninteger dimensions is much older (Hausdorf, 1919, Besicovitch, 1929). An analysis that compare the Haussdorff dimension and Kolmogorov capacities of self-similar structure with non integer fractal dimensions (Kolmogorov capacity or box counting dimension) was presented by Vassilicos (Vassilicos and Hunt, 1991). The wind velocity versus time are irregular curves of this type, with noninteger dimensions. These values correspond to the fractal dimension. A way of measuring the complexity of these series is by means of fractal dimension. The Fractal Dimension of wind components is a characteristic of turbulent flow as it has been shown in previous works where its relation to stability was highlighted (Tijera et al., 2012)

In this paper the integral scale of u` and w` component are compared. The scales are calculated using sonic anemometer data from three elevations 5.8 (~ 6), 13 and 32 m above the ground at the main tower site of the Sables 98 field campaign. Turbulent motion of the atmospheric flows occurs through a broad range of scales from the smallest ones that are usually defined as the scales at which the motion dissipates into heat due to the viscosity of the fluid until the larger scales corresponding to the integral scale. The integral scale can be defined in several ways: the larger scale of the flow, the scale above which the Fourier transform has a slope inferior a -5/3 slope, as which the turbulent kinetic energy (TKE) is maximum. Micrometeorological studies have found integral scale varying in a huge range, from around a hundred to a thousand meters (Teunissen, 1980, Kaimal and Finningan, 1994).

We study the anisotropy of the turbulent atmospheric flows in these scales comparing integral scale of fluctuations of the velocity component along of the mean wind direction and the vertical component at three different levels above the ground (5.8 m, 13.5 m, 32 m).

2 Theoretical background

The irregular behavior of the atmospheric turbulent fluxes in the boundary layer at large Reynolds number leads us to be interested in calculating their fractal dimension. Fractal dimension could help us to classify the irregularity of these flows. The more irregular the flow the greater its fractal dimension. Turbulent flows are characterized by the formation of many eddies of different length scales. Theses irregularities are due to the superimposition of eddies of different sizes and it is related with a broad range of scales which exist in turbulence. These scales vary from the smallest scales as dissipative scale to larger scale as integral scale. This paper is concerned with the analysis of the relationship between the integral scale and fractal dimension. As well as the relationship between the integral scale with the Bulk Richardson number, which provides a measure of the degree of stability in the flow, and how this turbulent flow is prone to develop instabilities. It is also used as a criterion for the existence or non-existence of turbulence in a stably stratified environment (a large positive value over a critical threshold, is indicative of a decaying turbulence or a completely non-turbulent) (Arya, 2001).

In this section we describe the methodology applied to calculate the fractal dimension and the integral scale. The estimation of the fractal dimension of time series has been the most commonly used criteria to measure their chaotic structure, there exist different works in that direction (Grassberger and Procaccia, 1982, Shirer et al, 1997). One of the methods most commonly used to estimate to fractal dimension of atmospheric flows has been the mean slope method through box-counting dimension using mean slopes of the graph of ln N (L) versus ln (L) for small ranges of L, where N(L) is the number of the boxes of side L necessary to cover the different points that have been registered in the physical space (velocity-time) (Falconer, 2000, Peitgen et al., 2004). As $L \to 0$ then N(L) increases, N meets the following relation:

$$N(L) \cong kL^{-d} \quad (1)$$

The value d is the box-counting dimension that is an approximation of the Hausdorff dimension and is calculated approximately by means of least-square-fitting of the representation of log N(L) versus log L obtaining the straight line regression given by the following equation:

$$\log N(L) = \log k - d \log L \quad (2)$$

The fractal dimension d will be given by the slope of this equation as is shown in the Fig.1.

In this paper we focus on calculating the integral scales for horizontal and vertical component fluctuations u' and w', and we studied their variations with respect to the fractal dimension and with the Bulk Richardson number, a turbulent parameter of stability.

These integral scales have been estimated using the normalized autocorrelation function and a Gaussian fit. The velocity autocorrelation function as a function of τ (lags number) for u' component is:

$$R(\tau)=\frac{\overline{u'(t)u'(t+\tau)}}{\overline{u'^{2}(t)}} \tag{3}$$

Integral time scale is:

$$T_L=\int_0^{\infty}R(\tau)d\tau\approx\int_0^{\tau'}R(\tau)d\tau \tag{4}$$

The integral time scale provides a measure of the scales of eddies in the x direction of a flow field. In the Eq. (2) we observed that τ' denotes the last lag in the data series. In boundary layer observations this time scale can be related to a length by multiplying the mean wind velocity by time scale. This requires the assumption of frozen turbulence known as Taylor's hypothesis (Panofsky and Dutton, 1984). The integral length scale can be defined as:

$$\lambda=\bar{v}T_L \tag{5}$$

The used method is based on Gaussian fit of the normalized autocorrelation function R(τ) and we calculated the value of τ that verifies the following equation:

$$\tau-\int_0^{\tau}R(\tau)d\tau=\int_{\tau}^{\tau'}R(\tau)d\tau \tag{6}$$

The Fig. 2 shows the Gaussian fit for an example of a data series of wind velocities with τ that verifies Eq. 6. This value allows us to calculate the integral time scale multiplying it by the time interval between each lag.

3 Description of Data

The data set was recorded in the Research Centre for the lower Atmosphere (CIBA in the Spanish acronymus), located in Valladolid province (Spain) and were measured in the experimental campaign Sables-98. This research centre was set up primarily to study the atmospheric boundary layer. The campaign took place from 10th to 27th September 1998 (Cuxart et al., 2000). This experimental site is a quite flat and homogeneous which forms a high plain of nearly 200 Km^2, surrounded by crop fields and some small bushes strewn over ground. Duero river flows along the SE border of the high plain. The synoptic conditions during the period of study of eight consecutive days (from 14 to 21 September) were controlled by a high pressure terrain system which produces thermal convection during the diurnal hours and from moderate to strong stable stratification during the nights.

In this work, data from sonic anemometers measured at a sampling rate of 20 Hz installed at 5.8 (~ 6), 13 and 32 m are analyzed. 5 - minute non - overlapping series are used to evaluate the different parameters. At a rate of 20 data points per second, sonic anemometers can resolve integral scales between about 10 m to 2000 m of u' horizontal component and 1 m to 1000 m of the w' vertical component, depending on the height in which the sonic anemometer is positioned and at the wind speeds typically measured in the Sables-98 experiment. We detect vertical scales over a broad range of scales from 1 m to 1000 m. The integral scales here are calculated based on the autocorrelation function, the mean wind velocity and integral time scale, and each of them can be expected to vary significantly. As the integral scale are the larger scales of turbulent flows it is possible to detect vertical scales larger than heights at which the sonic anemometer are located.

4 Results

4.1 Fractal Dimension, integral scale and stability of stratification.

In this paper we analyze the influence of stability of stratification on fractal dimension and integral scale. Different atmospheric surface-layers data are separated into thermal and dynamics stability classes based on a dimensionless parameter such as the Bulk Richardson number $R_{iB.}$ This parameter represents the ratio of the production or destruction of turbulence by buoyancy and by wind shear strain that is caused by mechanical forces in the atmosphere:

$$Ri_B = \frac{g}{\overline{\theta}} \frac{\Delta\overline{\theta}\Delta z}{(\Delta\overline{u})^2} \quad (7)$$

where g is the gravity acceleration and $\bar{\theta}$ the average potential temperature at the reference level, the term $\frac{g}{\bar{\theta}}$ is known as the buoyancy parameter. Ri_B is positive for stable stratification, negative for unstable stratification and approximate zero for neutral stratification (Arya, 2001). The way to calculate this number is described next:

1. Calculation of the mean potential temperatures at height z = 32 m, and close to the surface z = 5.8 m, namely $\overline{\theta_{32}}$ and $\overline{\theta_{5.8}}$ respectively. Being $\Delta\bar{\theta} = \overline{\theta_{32}} - \overline{\theta_{5.8}}$. The potential temperature has been estimated as relative to ground level by using the following formula: $\Delta\theta=\Delta T+\Gamma\Delta z$, $\Gamma=0.0098$ K m^{-1} (Arya 2001)
2. Obtaining of $\overline{u_z}$ the mean wind velocity module at the height z = 32 m and z = 5.8 m, denoted by $\overline{u_{32}}$ and $\overline{u_{5.8}}$ respectively, where $\Delta\bar{u} = \overline{u_{32}} - \overline{u_{5.8}}$.

Once the values of $\Delta\bar{\theta}$, $\Delta\bar{u}$ and Δz have been obtained by means of Eq. (7) we calculate the Bulk Richardson number in the layer between 32m and 5.8m.

In Fig. 3 we present the variation of the fractal dimension of the u' horizontal component of the velocity fluctuations along time at the three considered heights: 5.8 m, 13 m and 32 m. The behaviour of these variations is similar at the three heights. The w' component fluctuation presents an analogous behaviour. The fractal dimension values are in a range between 1.30 to nearly 1.00. We have found that during the diurnal hours the fractal dimension is bigger than at night (Tijera, 2012). We have no theoretical reason to explain this result, but a possible explanation of why this happens could be that fractal dimension is related with atmospheric stability and with the intensity of turbulence. It is well known that the intensity of turbulence grows up as solar radiation increases, producing instability close to the ground, mainly in the hours of noon. Therefore, one of the possible reasons of the increase of FD is the instability of the turbulent flow. In the other hand, during the nights a strong atmospheric stability usually exists, so the fractal dimension is usually smaller than during the diurnal hours.

In Fig 4 it is observed how the integral scale varies versus time at the three heights. There are some questions that have not been clarified yet in the literature. For example: How does the diurnal and night cycle influence on the integral scale? Which is the mechanism responsible for the growth of this integral scale? It has been observed in previous works that under certain conditions the turbulent flows self – organize and develop large-scale structures that take place through an inverse cascade that occurs in stably stratified anisotropic flows (with or without rotation) (Smith and Waleffe, 2002, Marino et al., 2014). The inverse cascade mechanism might also be responsible for the growth of the integral scale in the stratified atmosphere. It is a fundamental issue that should be clarified in a future research. As it is indicated in Fig 4 the integral scale for u` component varies between around 100 m on their smaller scales, until above 1500 m for its larger scales. The integral scales for w` component are slightly lower than for u` component. It is shown that these vertical scales can reach sizes between a few tens of meters until 1000 m in some occasions. It is observed, for each of them, that the greater is the height at which is located the sonic, the greater is the integral scale in the turbulent flow. Usually, at 32 m these scales are, on average, greater than those of the 13 m and the latter higher than at 5.8 m height.

Although the conceptual model of turbulence as eddies of various sizes is useful, it is difficult to obtain a correlation between the integral scale and fractal dimension in the atmosphere if we consider values throughout the whole day. However, it is much easier to find a relationship between the integral scale and fractal dimension of horizontal and vertical components of the wind velocity if we separate the hours of the day and night, and hence analyze the influence of diurnal and night cycle over these parameters. Daylight hours are from 6-18 UTC and the night from 18-6 UTC. These data set are analyzed in the three studied heights. Fig. 5 shows the variations of the integral scale versus fractal dimension at the level of 5.8 m for horizontal component. As it can be appreciated in Fig. 5 in the diurnal hours the average values of the integral scale versus the fractal dimension can be adjusted to the straight regression line given by the linear equation that appears on the top left of the graph. During those hours these values of the integral scale increase from a few tens of meters until 400 m with increasing values of the fractal dimension until 1.25. During the nights the average values of the integral scale decreases with the increase in the fractal dimension. These values also fit a straight regression line as it is indicated in Fig. 5. One of the possible explanations for this behaviour is that during the diurnal hours the average values of the integral scale increase due to the unstable stratification. During the nights, the existence of the stable stratification decreases the integral scale with an increase in fractal dimension until the approximate value of 1.2. This tendency appears also in the other two heights, at 13 m and at 32 m as it is shown in Fig. 6. Although during the diurnal cycle at 32 m the linear fit it is not so evident, the maximum scales are in the 1.15-1.20 range of the fractal dimension as it is illustrated in Fig. 6 (c).

For the vertical component in the three studied levels, the behaviour is slightly different. The averages values have fitted a quadratic function as it is indicated in Fig. 7. During the diurnal hours the averages values of integral scale reach maximum scales around the value of the fractal dimension of 1.15 at the three heights. From this value the integral scale decreases when fractal dimension increases.These maximum integral scales depend on the height. At the level of 5.8 m their sizes reach 50 m in average and the scattering of the values shows higher values that could reach 100 m. At the level of 13 m the values are about 100 m, the dispersion of these scales can reach sizes of 200 m and at the height of 32 m their larger average scales are approximately around 200 m and due to the variances of the data set could reach sizes of 400 m. From the value of the fractal dimension value of 1.15, the scales decrease until a few meters.

Throughout the night the average values of the integral scales decrease with the increase of the fractal dimension in a parabolic way as it is indicated in Fig. 7. This happens due to the stable stratification that occurs at nights. This behavior during the diurnal and night hours for w` component of the integral scale is similar to the results obtained for u` component, although the fits of averages values are parabolic and not linear. In all these cases our R^2 values and confidence levels are high as it is indicated in Fig. 7.

4.2 Relationship between integral scale and Bulk Richardson number

Among the numerous parameters existing to characterize the degree of stratification in the atmosphere we will use the Bulk Richardson number. The interpretation of this number has already been mentioned in the previous section. Here, we analyze how the integral scale of each one of the u` horizontal and w` vertical components varies with the Bulk Richardson number in diurnal and night cycle in the studied period. These results are shown in Fig. 8 for horizontal component and in Fig. 9 for vertical component. During the daylight hours appear the three kinds of stratification: unstable, neutral and stable as it is shown in the three graphs on the left side of Fig. 8 and the Fig. 9, each one corresponds to the different heights. In the stratification unstable and neutral the integral scales are higher than the integral scales under the influence of the stable stratification. At 5.8 m for the horizontal component these scales vary between 200 m and values slightly higher than 400 m and in the case of neutral stratification could increase until 600 m. This same behaviour occurs in the other two studied heights 13 m and 32 m although their scales are slightly higher as it is illustrated in Fig. 8. During the nights it is observed the biggest stability due to positive values of the Bulk Richardson number.

The same results are obtained for the integral scales of vertical component, although their sizes are smaller. At 5.8 m during the diurnal hours the average values reach about 50 m and during the night hours their values are below 50 m. At 13 m and 32 m in the diurnal hours the average values could reach about 150 m and 200 m and at the night hours are below 100 m and 200 m respectively.

4.3 Analysis of the anisotropy with the integral scale

In the last section we study the relationship between the integral scales of the horizontal and vertical components at different heights: 5.8 m, 13 m and 32 m. In Fig. 10 we represent the integral scale of u\` component versus the integral scale of w\` component at three studied heights and we find linear relations with the averages values of these scales. All integral scales measured during the period of study from 14 to 21 of September appear in this figure. The linear fits obtained are acceptable, with high R^2 values at 13 m and 32 m, as it is indicated in Fig. 10. The linear regression appears on the top left of each graph: at 5.8 m $L_{intu}(5.8\ m)= 1.46\ L_{intw}(5.8\ m) + 178$, at 13 m $L_{intu}(13\ m)=0.957\ L_{intw}(13\ m)+275.6$ and at 32 m $L_{intu}(32\ m)=0.646 L_{intw}(32\ m) +370$, being L_{intu} and L_{intw} the average values of the integral scale for horizontal and vertical component respectively.

The data in Fig. 10 appear quite scattered and the average values could be representative to find relationships between these scales. This scatter is due to the large number of uncontrolled variables, nonlocal disturbance, the presence of waves, horizontal inhomogeneity, low frequency disturbances, etc. These graphs are showing that the scale measured at 32 m is nearly always larger than the integral scale measured at 5.8 m. On the basis of the results obtained, we find slight differences between these components, thus there is anisotropy in atmospheric turbulent flows. In isotropic turbulence the integral scales of both components should be the same at the same height. Only under certain conditions and over limited scales is isotropy a property of turbulence in the stratified atmosphere (Thorpe, 2005).

5 Conclusions

We have calculated the fractal dimension and the integral scale of the horizontal and vertical components using wind velocity data from sonic anemometers at three different heights: 5.8 m, 13 m and 32 m. The numerical results show light significant differences on the diurnal and night cycle when the variation of the integral scale is analyzed versus the fractal dimension. Atmospheric stratification is analyzed for the three heights through the Bulk Richardson number, finding the three classical types of stratification along the diurnal cycle. It would be interesting for future works to study the growth of the integral scale in stratified flows and if it could be due to the inverse cascade on both diurnal and nighttime cycles. The main conclusions of this study are as follows.

Although all data appear quite scattered in this work, the averages values of these magnitudes show interesting results. During the diurnal hours the averages values of the integral scale of the horizontal component increases with the increase in fractal dimension until around 1.25 at 5.8 m and 13 m height. At these heights we have found linear fits between these magnitudes with high coefficients of correlation. While at 32 m the linear fit is not so evident, the maximum scales are in the 1.15-1.20 range of the fractal dimension. One of the possible explanations for this behaviour is that during the diurnal hours the average values of the integral scale increase due to the unstable stratification. During the night hours the average values of the integral scale decreases with the increase in the fractal dimension. These values also fit a straight regression line at the three analysed heights. During the nights the existence of the stable stratification decreases the integral scale with an increase in fractal dimension until the approximate value of 1.2.

For the vertical component of the integral scale the results are similar, even though with slight differences. The averages values have fitted a quadratic function. During the diurnal hours the averages values of integral scale reach maximum around the value of the fractal dimension of 1.15 in the three heights. From this value the integral scale decrease with the increase of the fractal dimension until a few meters. The different degree of stratification along diurnal hours will be reflected in that different behaviour from the value of 1.15. At nights when stability is normally major the integral scale decrease with increasing the fractal dimension of a parabolic way.

In the unstable and neutral stratification the integral scales are higher than the integral scales under the influence of the stable stratification.

To characterize the anisotropy of turbulent flows we have used the comparison of integral scales of horizontal and vertical component showing that the scale of u` component is almost always larger than the scale of the w` component at the same height.

Acknowledgements

This research has been funded by the Spanish Ministry of Science and Innovation (projects CGL2009-12797-C03-03). The GR35/10 program (supported by Banco Santander and UCM) has also partially financed this work through the Research Group "Micrometeorology and Climate Variability" (No 910437). Thank to participant teams in SABLES-98 for the facilities with the data.

References

Arya S. P.: Introduction to Micrometeorology. International Geophysics Series, Academic Press. 420 pp, 2001

Besicovitch A.S.: On linear sets of points of fractional dimension, Mathematische Annalen., 101, 161-193, 1929

Cuxart J., C. Yagüe, G. Morales, E. Terradellas., J. Orbe, J. Calvo, A. Fernández, M.R. Soler, C. Infante, P. Buenestado, A. Espinalt, H.E. Joergensen, J.M. Rees, J. Vilà, J.M. Redondo, I.R.Cantalapiedra, and Conangla, L.: Stable atmospheric boundary layer experiment in Spain (SABLES 98): A report. Bound.-Layer Meteorol., 96, 337-370, 2000

Falconer K. J.: Fractal Geometry Mathematical Foundations and Applications. John Wiley & Sons, Ltd. 288 pp, 1990

Frisch, U.: Turbulence. Cambridge Univesity Press, England, 296 pp, 1995.

Grassberger P. and Procaccia I.: Characterization of Strange Attractors. Phys. Rev Lett. 50, 346-349, 1983

Hausdorff F.: Math. Ann., 79, 157, 1919

Kaimal, J. C. and Finnigan J.: Atmospheric Boundary Layer Flows. Their Structure and Measurement, Oxford University Press. 289 pp, 1994

Mandelbrot,B.B.: Fractal: Form, Chance and Dimension. Freeman, San Fracisco, 1977

Mandelbrot, B.B.: The Fractal Geometry of Nature. Freeman, San Fracisco, 1985

Marino R, Mininni P.D., Rosenberg D. and Pouquet A.: Large-sale anisotropy in stably stratified rotating flows. Phys. Rev. E 90, 023018, 2014

Monin A.S. and Yaglom A.M.: Statistical Fluid Mechanics Mechanics Vol 1, Dover Publications, 769pp, 1971

Panofsky H.A. and Dutton. : Atmospheric Turbulence, Wiley and Sons, 397 pp, 1984

Peitgen H., Jürgens H. and Saupe D.: Chaos and Fractals. Springer-Verlag 971pp, 2004

Shirer H.N., Fosmire C.J., Wells R., Suciu L.: Estimating the Correlation Dimension of Atmospheric Time Series. Journal of the Atmospheric Sciences: 54, pp. 211-229, 1997

Smith L and Waleffe F.: Generation of slow large scales in forced rotating stratified turbulence. J. Fluid Mech. 451, 145-168, 2002

Stull R. B.: An introduction to Boundary Layer Meteorology. Kluwer Academic Publishers, 670pp, 1988

Teunissen H W.: Structure of mean winds and turbulence in the planetary boundary layer over rural terrain. Boundary-Layer Meteorol 19, 187-221, 1980.

Thorpe S.A.: The turbulent ocean. Cambridge Universe Press. 439pp., 2005

Tijera, M., Cano J.L., Cano, D., Bolster, B. and Redondo J.M.: Filtered deterministic waves and analysis of the fractal dimension of the components of the wind velocity. Nuovo Cimento C. Geophysics and Space Physics 31, 653-667, 2008

Tijera M., Maqueda G.,Yaque C., and Cano J.: Analysis of fractal dimension of the wind speed and its relations with turbulent and stability parameters , Intech, Fractal Analisis and Chaos in Geosciences, 29-46, 2012

Vassilicos J. C. and J. C. R. Hunt.: Fractal dimensions and spectra of interfaces with application to turbulence. Proc. R. Soc. London Ser. A 435-505., 1991

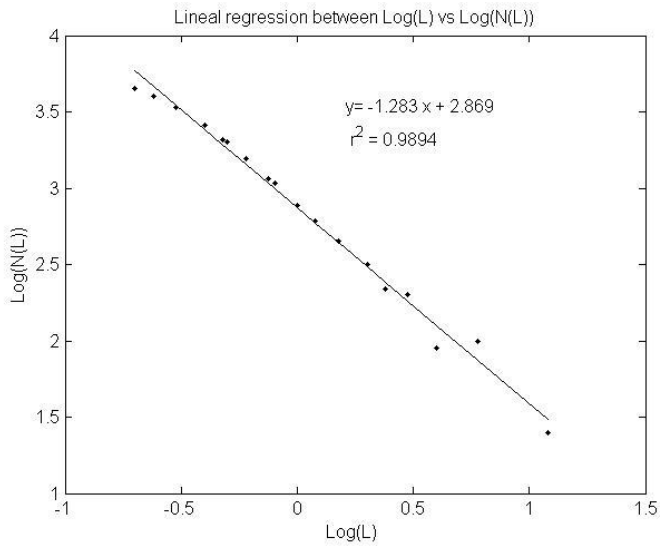

Figure 1.Example of linear regression between number of not empty boxes and length side of the box. The slope (*d*) is the fractal dimension of the w` component, *d* =1.28 ± 0.03 for a example of the w' component of the wind velocity.

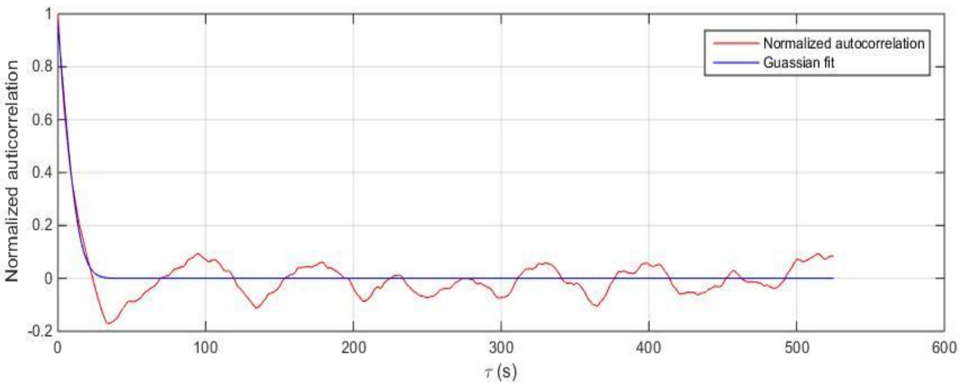

Figure 2. Gaussian fit for a data series of wind velocities u' component that allows us to calculate the integral scale.

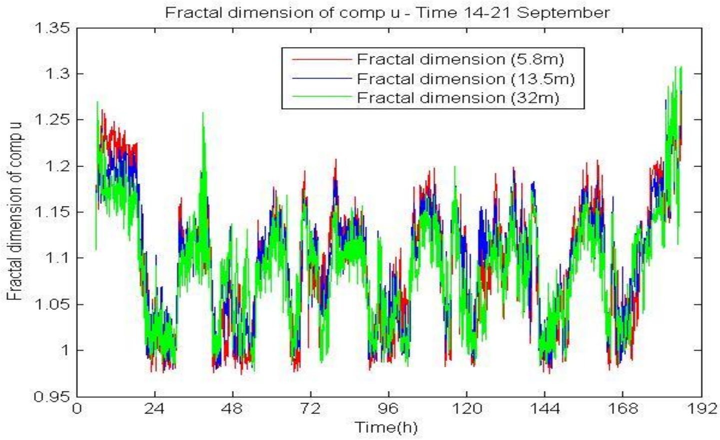

Figure 3. Variation of the fractal dimension versus time for the u' component fluctuation at the three heights, showing the influence of the diurnal cycle.

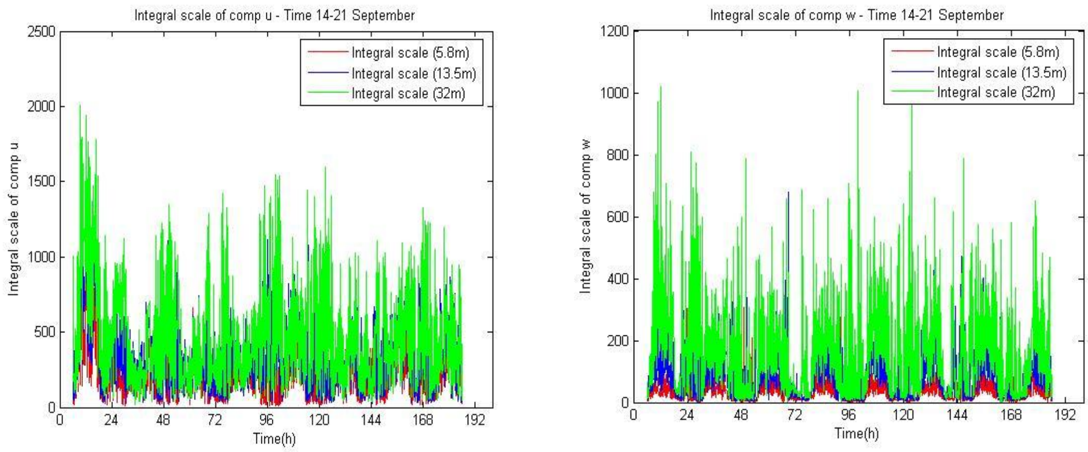

Figure 4. Variation of the integral length scale of horizontal and vertical components versus time at the three heights.

(a) (b)

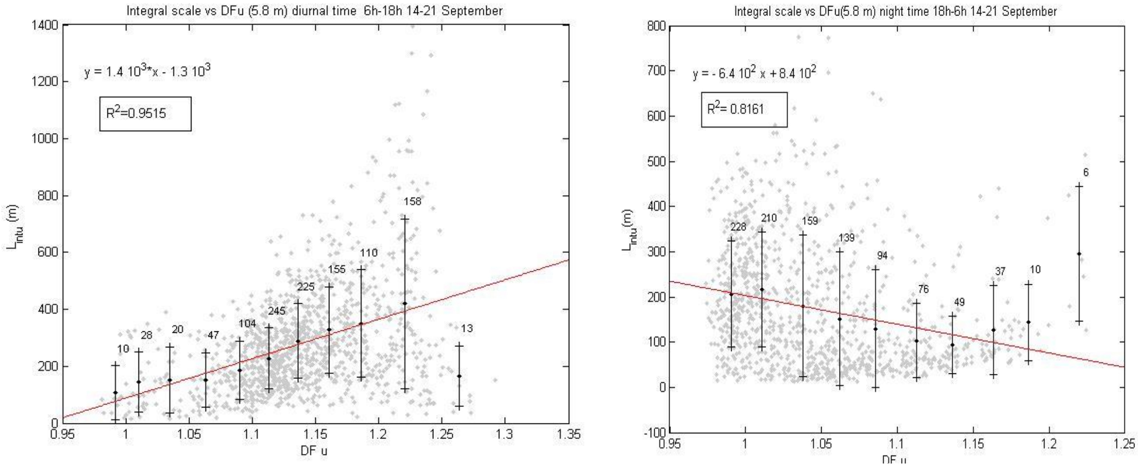

Figure 5. Variations of the integral scale versus the fractal dimension of u\` component of the wind velocity at 5.8 m. (a) diurnal hours 6h-18 h, (b) night hours 18 h - 6h. On the top left of the each graph it is indicated the linear regression of the averages values (a) L_{intu}(5.8 m)= 1.4 10^3 DF_u(5.8 m) - 1.3 10^3 (b) L_{intu}(5.8 m)= -6.4 10^2 DF_u(5.8 m) + 8.4 10^2, being L_{intu} and DF_u the integral scale and fractal dimension for u component respectively.

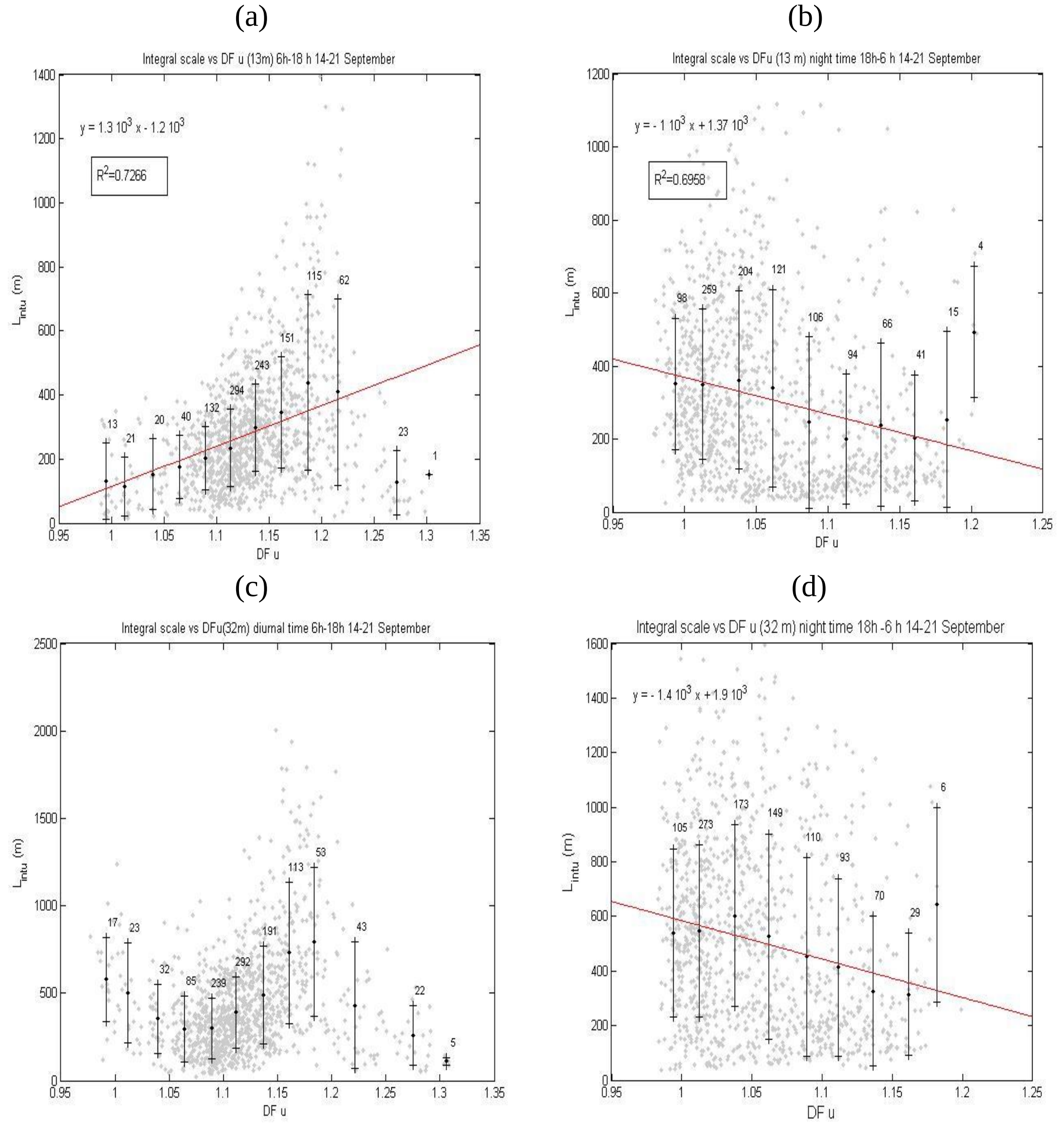

Figure 6. Variations of the integral scale versus the fractal dimension of the u` component at 13 m and 32 m. (a) and (c) diurnal hours, (b) and (d) night hours. In the same manner that in the figure 5 the linear fits are: (a) L_{intu}(13 m)= 1.3 10^3 DF_u(13 m) - 1.2 10^3 , (b) L_{intu}(13 m)= -1 10^3 DF_u(13 m) + 1.37 10^3, (d) L_{intu}(32 m)= -1.4 10^3 DF_u(32 m) + 1.9 10^3

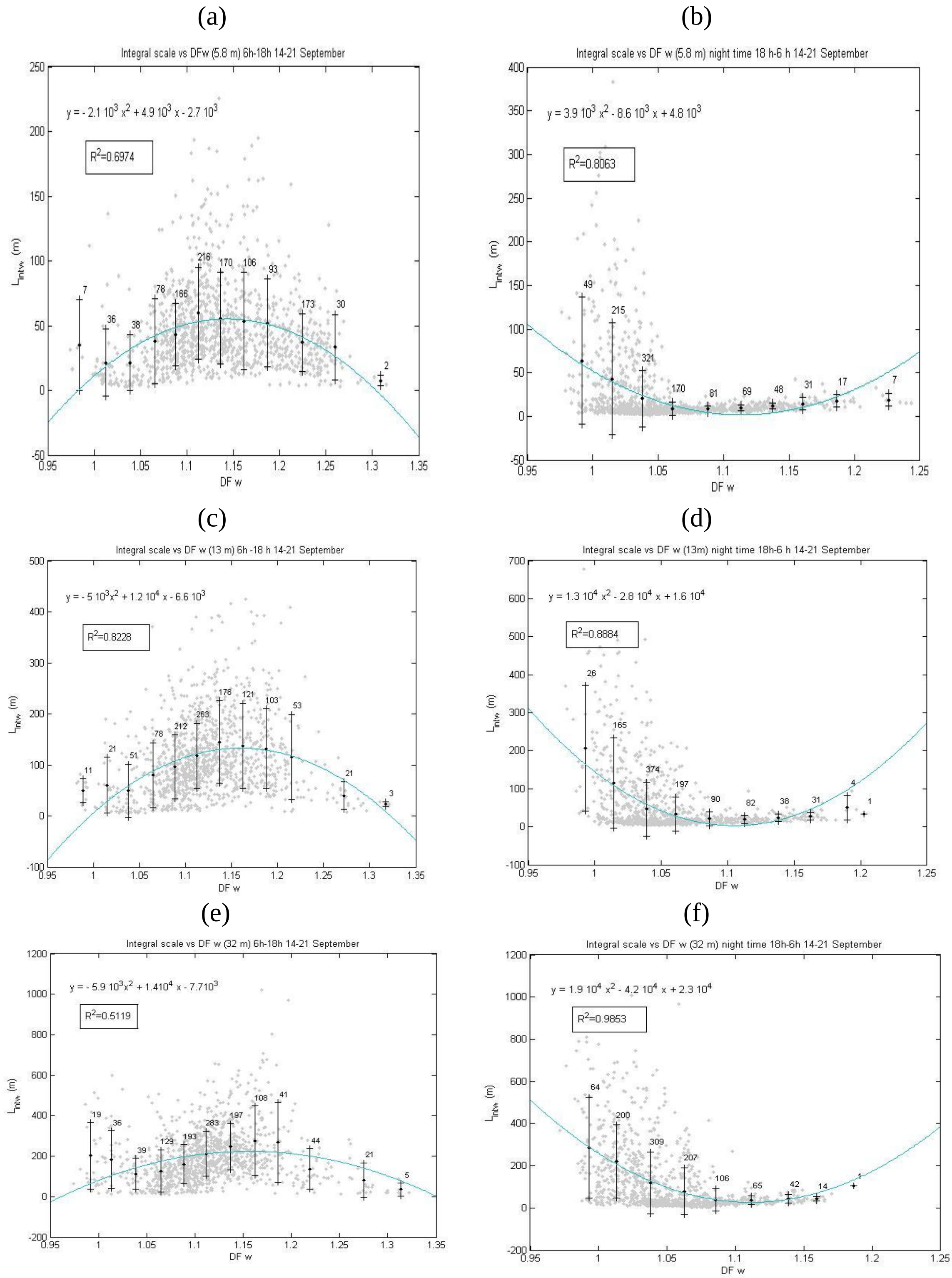

Figure 7. Variations of the integral scale versus the fractal dimension of the w` component at 5.8 m, 13 m and 32 m. (a), (c) y (e) diurnal hours, (b), (d) y (f) night. hours. The fits to a quadratic function of the averages values appear on the top left of the each graph, being y variable L_{intw} and x variable DF_w.

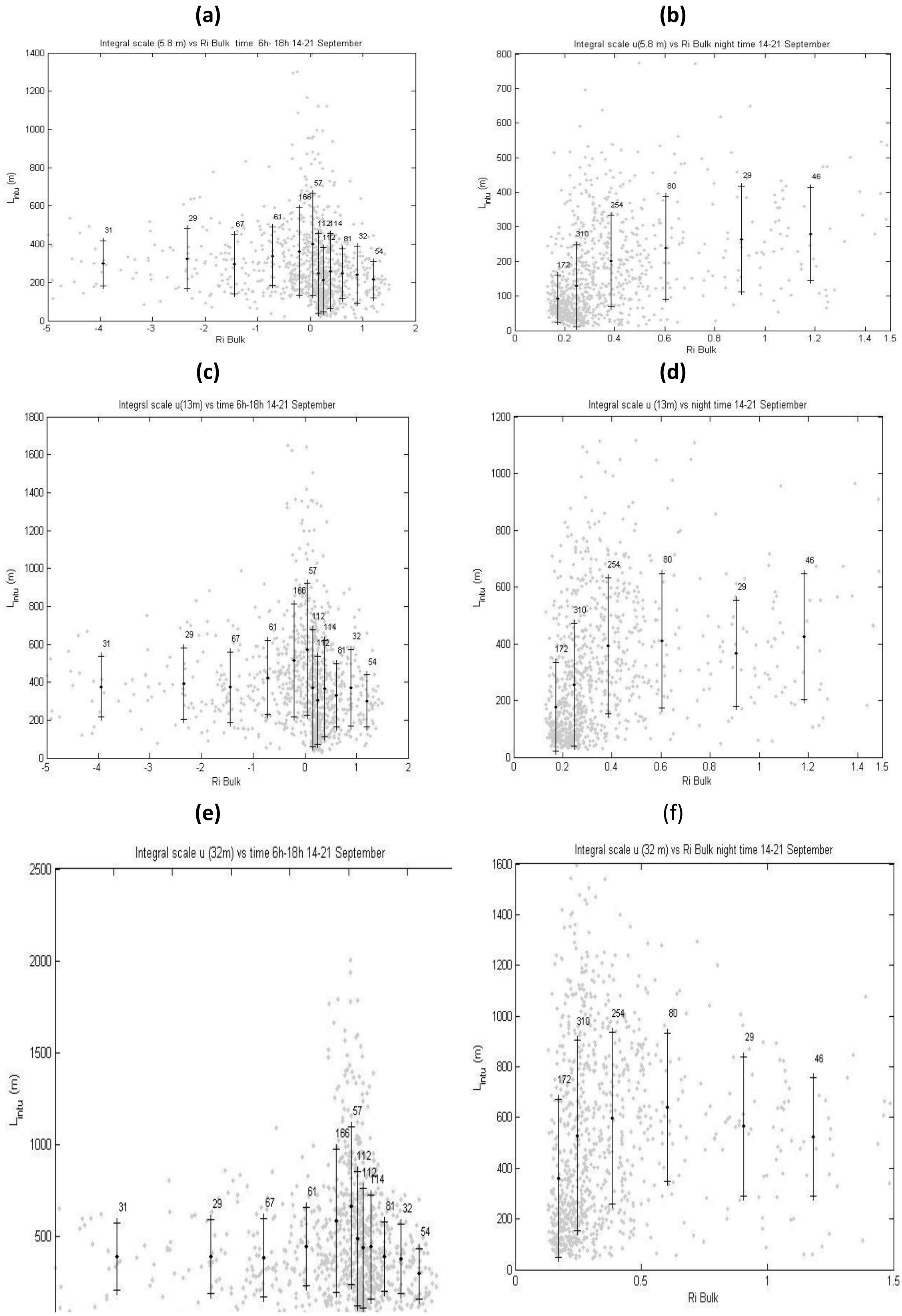

Figure 8. Integral length scales of u` component plotted against the Bulk Richardson number at 5.8 m, 13 m and 32 m. (a), (c) y (e) diurnal hours, (b), (d) y (f) night hours.

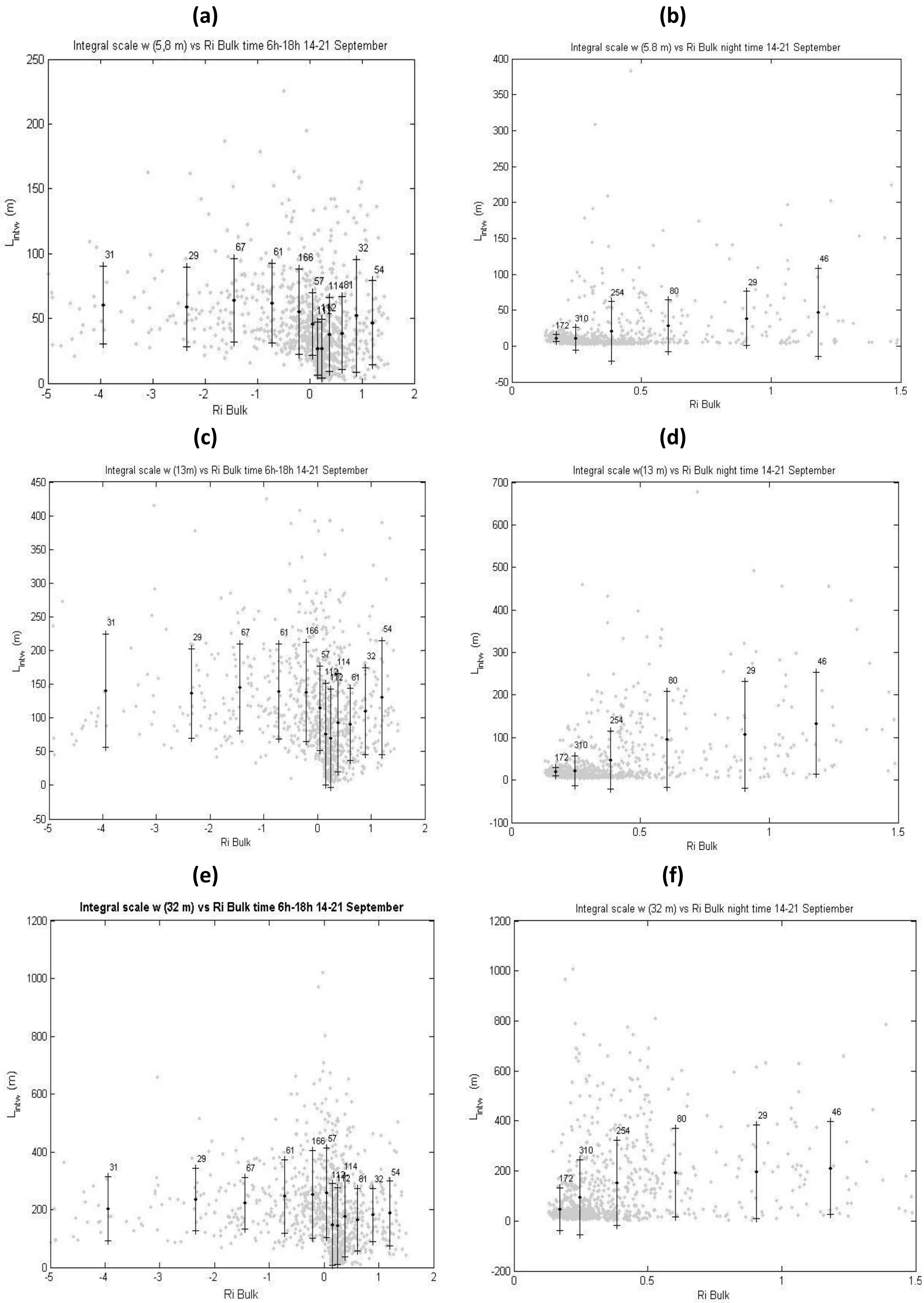

Figure 9. Integral length scales of w` component plotted against the Bulk Richardson number at 5.8 m, 13 m and 32 m. (a), (c) y (e) diurnal hours, (b), (d) y (f) night hours.

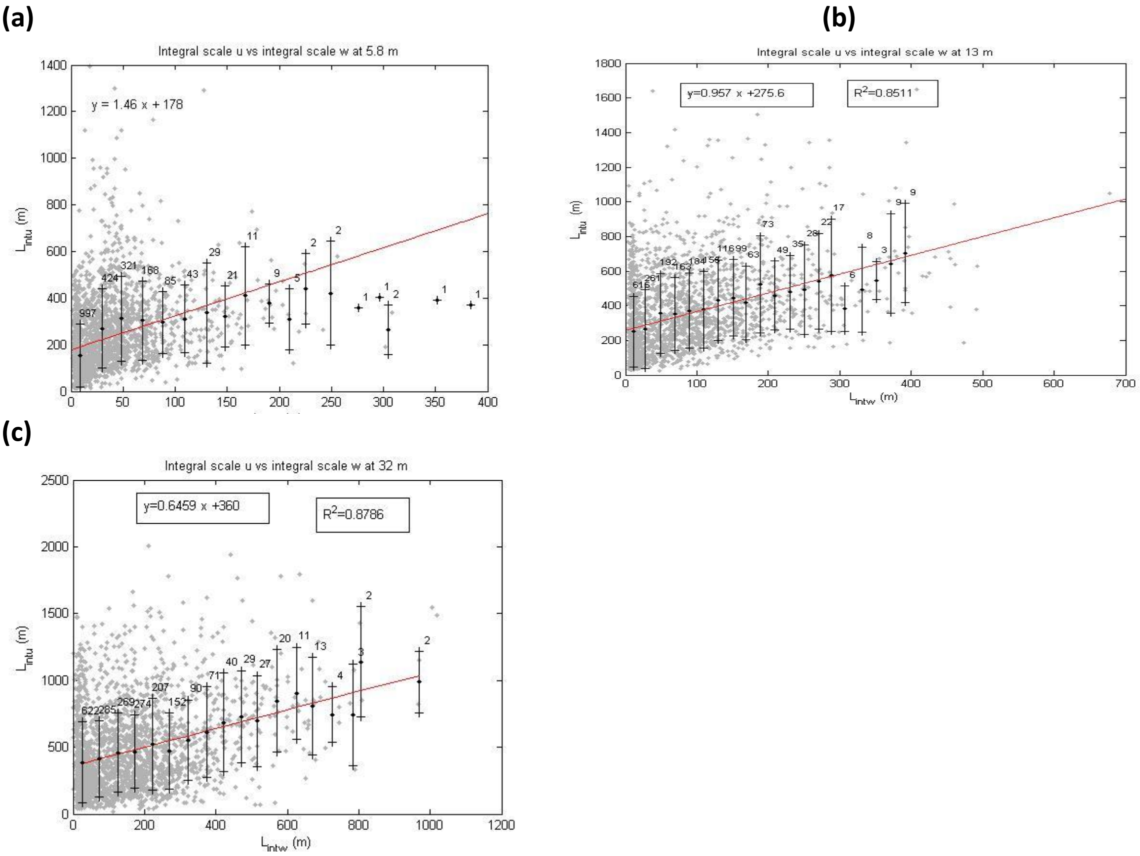

Figure 10. Comparison of integral scales of u` component and w` component (a) at 5.8 m, (c) at 13 m (c) at 32 m, showing that the averages values of theses scales fit to the linear regression indicated on the top left each graph. Data set appear as a cluster around the straight line.