

1 **Explanation of the values of Hack's drainage basin, river length scaling exponent**

2

3 A. G. Hunt

4 Department of Physics and Department of Earth & Environmental Sciences

5 Wright State University

6 3640 Colonel Glenn Highway

7 Dayton, OH 45435

8 e-mail allen.hunt@wright.edu

9 Tel. (937) 775-3116

10 Fax (937) 775-2222

11

12 **Abstract**

13

14 Percolation theory can be used to find water flow paths of least resistance. The application of

15 percolation theory to drainage networks allows identification of the range of exponent values that

16 describe the tortuosity of rivers in real river networks, which is then used to generate the

17 observed scaling between drainage basin area and channel length, a relationship known as

18 Hack's law. Such a theoretical basis for Hack's law allows interpretation of the range of

19 exponent values based on an assessment of the heterogeneity of the substrate.

20

21

22 **1. Introduction**

23 River networks display complex organization as documented in numerous studies. This
24 work addresses but one of them. In particular, the relationship between drainage basin area and
25 river length is non-trivial. In Euclidean geometry the basin area, A , would be proportional to the
26 square of the river length, l , i.e., l should be proportional to $A^{1/2}$. In actuality, as determined by
27 Hack (1957), this relationship is

$$l = CA^\beta$$

28 with the value of β approximately 0.6. Later investigations did not always return the identical
29 value of β . Nevertheless, Maritan et al., (1996) consider the value of β to be well constrained and
30 refer to a study of Gray (1961) as having established that “the accepted values for the exponent
31 $[\beta]$ are in the range 0.57 to 0.6.”

32 Hack (1957) (page 65) asserts that the relationship was a consequence of the lengthening
33 of drainage basins with increasing size. But Montgomery and Dietrich (1992) compare straight-
34 line basin length, L , to A over seven orders of magnitude of length scale and find precisely $L =$
35 $A^{0.5}$. That result allows stream length to be expressed in terms of the straight line basin
36 dimension, $l = A^\beta = L^\gamma = L^{2\beta}$, so that $\gamma = 2\beta$. The exponent $\gamma > 1$ then defines the tortuosity
37 (sometimes known as sinuosity) of the stream path through the drainage basin.

38 Hack’s law explanations have been sought in fractal (Tarboton et al., 1988; Maritan et al.,
39 1996), constructal (Reis, 2006) and “feasible optimality” (Rigon et al., 1998) theories. Fractal
40 theories produce the required self-similar drainage basins (Peckham, 1995) as well as increasing
41 stream sinuosity downstream. I suggest that Hack’s law can be understood using percolation
42 theory (Stauffer and Aharony, 1994) because 1) the fractal structure of the percolation cluster

43 generates values for γ that constrain the data for β appropriately, and 2) it exploits the concept
44 that water flows along paths of minimal resistance, as in the subsurface also (Hunt et al., 2014).

45 **2. Theory**

46 There are two distinct applications of percolation theory to flow or conduction problems,
47 and these two applications are those that provide the bounds to Hack's exponent values. The
48 more familiar application is to a binary system, where, e.g., bonds either connect neighboring
49 sites (which in the simplest case are located on a lattice, or grid), or they don't. If enough such
50 neighboring sites are connected, a continuous path of interconnected bonds spans the system.
51 This is denoted the percolation threshold. The shortest distance across the system within this
52 connected cluster is called the chemical path length (Porto et al., 1997). Since all bonds have
53 equal resistance, the shortest flow path has also the lowest resistance and optimal dissipation.

54 The second possibility is a system, in which bonds of varying resistance connect each
55 pair of neighboring sites. When the system is strongly heterogeneous, i.e., when the distribution
56 of the natural logarithm of the resistances has variance, $\sigma^2 \gg 1$, the proper application of
57 percolation theory is to find the path of least cumulative resistance. Quantification of this process
58 equates an integral over the local conductance distribution, from a "critical" value, to the largest
59 value, with the percolation threshold (Pollak, 1972). This particular method became known as
60 "critical path analysis," or CPA. The subnetwork so defined is precisely at the percolation
61 threshold (Stauffer and Aharony, 1994). However, the most interesting path across this system is
62 not the shortest, but the optimal path, which provides the least energetic cost in a strongly
63 heterogeneous network (Lopez et al., 2005). Since the optimization is not for length, but for
64 energy costs, the path is longer than in a homogeneous system, meaning that its tortuosity
65 exponent is larger. Nevertheless, the chemical path length in a homogeneous system is the

66 analogue of the optimal path length in a heterogeneous system, because the shortest path in a
67 system with identical links also represents the path of minimum energy dissipation.

68 River paths above the Earth's surface are obviously impossible, and considering the
69 orders of magnitude slower subsurface flow rates, we can also neglect interchanges between
70 surface and subsurface as part of the river flow network, meaning that all hydrologic connections
71 are restricted to the Earth's surface. This lack of alternate paths above or below the surface
72 makes two dimensional (2D) connectivity and universality relevant, regardless of the best
73 particular description of the roughness of the Earth's surface. In two dimensions, the chemical
74 path length scales with the system size (Sheppard et al., 1999), L as $L^{1.13}$, but the optimal path
75 length scales with system size as $L^{1.21}$, and does not depend on the particular percolation model.
76 Calling the scaling exponent, γ , as above, we therefore find that percolation theory constrains its
77 values to be, $1.13 < \gamma < 1.21$. It should be emphasized here that the precision of the numerical
78 calculations of these exponent values by Sheppard et al., (1999) exceeded all other attempts by at
79 least an order of magnitude, and these values are not best viewed as "empirical estimates." These
80 are thus two possible endpoints for the application of percolation theory to the formation of river
81 networks, and are generated from homogeneous and heterogeneous systems, respectively. How
82 could they be realized in nature, or in landscape models?

83 One can start from an initially homogeneous landscape, and allow stream incision
84 through random headward erosion, analogous to the processes treated in early landscape
85 evolution models (Willgoose et al., 1991), which generate hierarchical structures from random
86 chance associated with rainfall magnitude variability. A connected path with the lowest
87 dissipation (shortest length) will soon acquire the highest flow, through channel erosion

88 feedbacks. Thus, once a river makes a random choice, the enhanced erosion power from the
89 stream reinforces the initial random choice.

90 The optimal path exponent describes the tortuosity of a channel, when the channel is
91 determined by a global optimization of the flow path in a heterogeneous substrate, and could not
92 be a simple product of headward erosion, which might produce only a local optimization. In such
93 a case geological constraints from varying erodibility can dominate as channels extend either
94 upward, by headward erosion, or downward, (e.g., by overtopping of sills).

95 Using the above result that $\gamma = 2\beta$, we find for Hack's (1957) original result, $l = L^{1.2}$. The
96 range quoted by Maritan et al. (1996), $0.57 < \beta < 0.6$ generates $1.14 < \gamma < 1.20$. The predicted
97 range of tortuosity exponent values, γ , $1.13 < \gamma < 1.21$, generated by percolation theory appears
98 to differ by less than 1% from the observed range of values, and to be slightly larger, consistent
99 with interpreting this range as bounds on observed values.

100 Note that, while e.g., Willemin (2000) found a wider range of β (0.5 to 0.7) than did Gray
101 (1961), these values were for limited statistics (as small as 11 data points). Although this range is
102 wide, compared with our predictions, when all statistics were put together (Willemin's figure 11)
103 the resulting value of β was 0.58. Further, individual values did increase monotonically with
104 increasing geologic heterogeneity. Northwestern Iowa, in the middle of the North American
105 craton produced 0.5, New York, 0.64, and coastal Oregon, in a region of active tectonics, 0.7.
106 Finally, the range of values for β quoted by Gray (1961) arose from his consideration of studies
107 over different regions with distinct terrain; uncertainty in a given region was reported in the
108 variation of the numerical prefactor, rather than the exponent.

109 **4. Conclusions**

- 110 • Percolation predictions generate the range of exponents observed in Hack’s law,
111 including the tendency for the largest exponent values to occur in geologically
112 heterogeneous environments.
- 113 • The statistical nature of percolation theory is in accord with the tendency of the
114 spread in Hack’s exponent values to diminish with increasing sample size,
- 115 • The source of the tortuosity in the “optimal paths” of lowest energy dissipation is
116 in general accord with the “feasible optimality” (Rigon et al., 1998) proposed to
117 explain Hack’s law.

118 **References**

119 Gray, D. M., 1961, Interrelationships of watershed characteristics, *J. Geophys. Res.* 66 1215-
120 1223.

121 Hack, J.T., 1957. Studies of longitudinal profiles in Virginia and Maryland. USGS Professional
122 Papers 294-B, Washington DC, pp. 46–97.

123 Hunt, A. G., R. P. Ewing, and B. Ghanbarian, 2014, *Percolation Theory for Flow in Porous*
124 *Media*, 3rd edition, Springer, Berlin.

125 Lopez, E., S. V. Buldyrev, L. A. Braunstein, S. Havlin, and H. E. Stanley, Possible connection
126 between the optimal path and flow in percolation clusters, *Phys. Rev. E* **72**, 056131
127 _2005_

128 Maritan, A., A. Rinaldo, R. Rigon, A. Giacometti, and I. Rodriguez-Iturbe, 1996, Scaling laws for
129 river networks, *Phys. Rev. E* 53(2): 1510-1515.

130 Montgomery, D. R., and W. E. Dietrich, 1992, Channel initiation and the problem of landscape
131 scale, *Science*, 255 826-830.

132 Peckham, S. D., 1995, New results for self-similar trees with applications to river networks,
133 Water Resour. Res. 31 (4) 1023-1029.

134 Pollak, M., 1972, A percolation treatment of dc hopping conduction, *J. Non Cryst. Solids*, **11**: 1-
135 24, doi:10.1016/0022-3093(72)90304-3.

136 Porto, M, A. Bunde, S. Havlin, and H. E. Roman, 1997, Structural and dynamical properties of
137 the percolation backbone in two and three dimensions, *Phys. Rev. E* 56: 1667-1675

138 Reis, A. H., 2006, Constructal view of scaling laws of river basins, *Geomorphology* 78: 201-
139 206.

140 Rigon, R., I. Rodriguez-Iturbe, and A. Rinaldo, 1998, Feasible optimality implies Hack's law,
141 Water Resour. Res. 34 (11) 3181-3189.

142 Sheppard, A. P., M. A. Knackstedt, W. V. Pinczewski, and M. Sahimi, 1999, Invasion
143 percolation: new algorithms and universality classes, *J. Phys. A: Math. Gen.* **32**: L521-
144 L529.

145 Stauffer, D. and A. Aharony, 1994, *Introduction to Percolation Theory*, 2nd edition, Taylor and
146 Francis, London.

147 Tarboton, D. G., R. L. Bras, I. Rodriguez-Iturbe, I, 1988, The fractal nature of river networks,
148 Water Resour. Res. 24: 1317-1322.

149 Willemin, J. H., 2000, Hack's law: Sinuosity, convexity, elongation, Water Resour. Res. 36 (11)
150 3365-3374.

151 Willgoose, G., R. L. Bras, and I. Rodriguez-Iturbe, 1991, A coupled channel network growth and
152 hillslope evolution model. 2. Nondimensionalization and applications, Water Resour.
153 Res. 27 (7) 1685-1696.