

1 **The Transient Variation of the Complexes of the Low Latitude Ionosphere within the**
2 **Equatorial Ionization Anomaly Region of Nigeria.**

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12 **Abstract**

13 The quest to find an index for proper characterization and description of the dynamical response
14 of the ionosphere to external influences and its various internal irregularities has led to the study
15 of the day to day variations of the chaoticity and dynamical complexity of the ionosphere. This
16 study was conducted using Global Positioning System (GPS) Total Electron Content (TEC) time
17 series, measured in the year 2011, from 5 GPS receiver stations in Nigeria which lies within the
18 Equatorial Ionization Anomaly region. The nonlinear aspect of the TEC time series were
19 obtained by detrending the data. The detrended TEC time series were subjected to various
20 analyses to obtain the phase space reconstruction and to compute the chaotic quantifiers which
21 are Lyapunov exponents LE, correlation dimension, and Tsallis entropy for the study of
22 dynamical complexity. Considering all the days of the year the daily/transient variations show no
23 definite pattern for each month but day to day values of Lyapunov exponent for the entire year
24 show a wavelike semiannual variation pattern with lower values around March, April, September
25 and October, a change in pattern which demonstrates the self-organized critical phenomenon of
26 the system. This can be seen from the correlation dimension with values between 2.7 and 3.2
27 with lower values occurring mostly during storm periods demonstrating a phase transition from
28 higher dimension during the quiet periods to lower dimension during storms for most of the
29 stations. The values of Tsallis entropy show similar variation pattern with that of Lyapunov
30 Exponent, with both quantifiers correlating within the range of 0.79 to 0.82. These results show

31 that both quantifiers can be further used together as indices in the study of the variations of the
32 dynamical complexity of the ionosphere. The presence of chaos and high variations in the
33 dynamical complexity, even at quiet periods in the ionosphere may be due to the internal
34 dynamics and inherent irregularities of the ionosphere which exhibit non-linear properties.
35 However, this inherent dynamics may be complicated by external factors like Geomagnetic
36 storms. This may be the main reason for the drop in the values of Lyapunov exponent and Tsallis
37 entropy during storms. The dynamical behavior of the ionosphere throughout the year as
38 described by these quantifiers, were discussed in this work.

39

40 **1.0 Introduction**

41 The behavior of natural systems like the ionosphere is a function of changes that occur in the
42 underlying dynamics that exists in such system. These underlying dynamics however can be
43 sometimes complex and nonlinear due to superposition of different changes in dynamical
44 variables that constitute it. When the dynamical states of a system changes suddenly due to
45 sudden changes in the external factor affecting the system, then such a system is said to be
46 deterministic.

47 However, there is no totally deterministic system in nature, because all natural systems exhibit a
48 mixture of both deterministic properties. Although few natural systems have been found to be
49 low dimensional deterministic in the sense of the theory, the concept of low-dimensional chaos
50 has been proven to be fruitful in the understanding of many complex phenomena (Hegger et al.,
51 1999) The degree of determinism or stochasticity in most natural systems is dependent on how
52 much the system can be influenced by external factors, the nature of these external factors among
53 others .The ionosphere like every other natural system possess its intrinsic dynamics and it can
54 also be influenced by other external factors. The typical characteristics of a dynamical system
55 like the ionosphere is expected to naturally show the interplay between determinism and
56 stochasticity simply because of the fact that the ionosphere which has an inherent internal
57 dynamics is also influenced by the influx of stochastic drivers like the solar wind, since it is
58 influenced by external dynamics like every other natural system. This has made pure
59 determinism impossible in the ionosphere, a situation that is common to all natural system and its
60 surrounding.

61 The intensity of the solar wind coming into the ionosphere varies with the solar activity and an
62 extreme solar activity can lead to geomagnetic storms and substorms drive in high intensity
63 plasma wind at enormous speed and it serves as major stochastic driver leading to storm. The
64 solar wind is driven from the sun into the ionospheric system during the quiet and storm and
65 during relatively quiet periods of each month of the year. However other processes which
66 include various factors like local time variations of the neutral winds, ionization processes,
67 production-recombination rates, photoionization processes, plasma diffusion and various
68 electrodynamics processes. (Unnikrishnan, 2010). The mesosphere and the lower thermospheric
69 dynamics as reported by Kazimirovsky and Vergasova (2009) and also the influence of gravity
70 waves as reported by Sindelarova (2009) can also be of great influence on the internal dynamics
71 of the ionosphere.

72 Therefore, it is of great importance to study the chaoticity and dynamical complexity of the
73 ionosphere and its variations in all geophysical conditions. However a good number of
74 investigations have been carried out on concept of chaos in the upper atmosphere before now
75 which includes the study on magnetospheric dynamics and the ionosphere. The study of chaos in
76 magnetospheric index time series such as AE and AL were initially carried out by (Vasiliadis et
77 al., 1990, Shan et al, 1999; Pavlos et al, 1992). These previous efforts made by the
78 aforementioned researchers has led to the development of the concept of investigating and
79 revealing the chaoticity and the complex dynamics of the ionosphere, and as a result, studies on
80 the chaoticity of the ionosphere have been conducted, by some investigators like Bhattacharyya
81 (1990) who studied chaotic behavior of ionospheric diversity fluctuation using amplitude and
82 phase scintillation data, and found the existence of low dimension chaos. Also, Wernik and Yeh
83 (1994) further revealed the chaotic behavior of the ionospheric turbulence using scintillation data
84 and numerical modeling of scintillation at high latitude. They showed that the ionospheric
85 turbulence attractor (if it exists) cannot be reconstructed from amplitude scintillation data and
86 their measured phase scintillation data adequately reproduce the assumed chaotic structure in the
87 ionosphere. Also Kumar et al., (2004) reported the evidence of chaos in the ionosphere by
88 showing the chaotic nature of the underlying dynamics of the fluctuations of TEC power
89 spectrum indicating exponential decay and the calculated positive value of Lyapunov exponent.
90 This is also supported by the results of the comparison of the chaotic characteristics of the time
91 series of variations of TEC with the pseudochaotic characteristic of the colored noise time series.

92 Xuann et al., (2006) studied chaos properties of ionospheric total electron content (TEC) using
93 TEC data from 1996 to 2004, and analyze possibility to predict it by using chaos. They found
94 the presence of chaos in the TEC measured in the study area, as indicated by the positive
95 Lyapunov exponent computed from their data. The correlation dimension was 3.6092 from their
96 estimation. They were also able to show that the TEC time series can be predicted using chaos.

97

98 Also, Unnikrishnan et al (2006a,b) have analyzed the deterministic chaos in mid latitude and
99 Unnikrishnan (2010), Unnikrishnan and Ravindran (2010), analyzed some TEC data from some
100 Indian low latitude stations for quiet period and major storm period and found in Their results the
101 presence of chaos which was indicated by a positive Lyapunov exponent, and they also inferred
102 that storm periods exhibits lower values compared to quiet periods. The dynamical complexity
103 of magnetospheric processes and the ionosphere have been studied by a number of researchers.
104 Balasis et al., (2008) investigated the dynamical complexity of the magnetosphere by using
105 Tsallis entropy as a dynamical complexity measure in D_{st} time series also Balasis et al., (2009)
106 investigated the dynamical complexity in D_{st} further by considering different entropy measures.
107 Coco et al (2011) using the information theory approach studied the dynamical changes of the
108 polar cap potential which is characteristic of the polar region ionosphere by considering three
109 cases (i) steady IMF $B_z > 0$, (ii) steady IMF $B_z < 0$ and (iii) a double rotation from negative to
110 positive and then positive to negative B_z . They observed a neat dynamical topological transition
111 when the IMF B_z turns from negative to positive and vice versa, pointing toward the possible
112 occurrence of an order/disorder phase transition, which is the counterpart of the large scale
113 convection rearrangement and of the increase of the global coherence. Further studies in chaotic
114 behavior and nonlinear dynamics is however needed to improve our understanding of the
115 dynamical behavior of the ionosphere of low latitude ionosphere especially over Africa during
116 quiet and storm for different season of the year some as to be able to characterize chaoticity for
117 different season of the year for quiet and storm periods. Recently Ogunsua et al (2014) studied
118 comparatively the chaoticity of the equatorial ionosphere over Nigeria using TEC data,
119 considering five quietest day classification and five most disturbed day classification. They were
120 able to show the presence of chaos as indicated the positive Lyapunov exponents and also were
121 able to show that Tsallis entropy can be used as a viable measure of dynamical complexity in the

122 ionosphere with portions showing lower values of Tsallis entropy indicating lower dynamical
123 complexity, with a good relationship with Lyapunov exponents. They found a phase transition
124 from higher dimension during quiet days to Lower dimension during storm.

125

126 The low latitude region where Nigeria is situated is known as the equatorial anomaly region,
127 where the magnetic field B is almost totally parallel to the equator. Off the equator the E region
128 electric field maps map along the magnetic field up to the F-region altitude in the low latitude,
129 this eastward electric field (E) interacts with the magnetic field B at the F region during the day.
130 This results in the electrodynamic lifting of the F-region plasma over the equator, known as
131 EXB drift. The uplifted plasma over the equator moves along the magnetic line in response to
132 gravity, diffusion and pressure gradients and hence, the fountain effect. The fountain effect being
133 controlled by the EXB drift shows the dynamics of the diurnal variation equatorial anomaly
134 (Abdu, 1997; Unnikrishnan 2010). There is a reduction in the F region ionization density at the
135 magnetic equator and also much enhanced ionization density at the two anomaly crests within
136 $\pm 15^\circ$ of the magnetic latitude north and south of the equator (Rama Rao et al., 2006). The
137 equatorial ionization anomaly and other natural processes which includes various ionization
138 processes and recombination; influx of solar wind, photoionization processes and so many other
139 factors that occur due to variations in solar activities, have a great influence on the systems of
140 the ionosphere, due to their effects on internal dynamics of the ionosphere. This portrays the
141 ionosphere as a typical natural system with continuous interaction with its external environment
142 which led to the study of the influence of the sun on the ionosphere (Ogunsua et al., 2014).

143 The ionosphere possesses a significant level of nonlinear variations that requires more
144 investigation which can be studied and characterized using nonlinear approach like the chaoticity
145 and dynamical complexity for the study of its dynamics. The need to study the daily variation in
146 the dynamical complexity of the ionosphere arises from the established knowledge and
147 understanding which shows that the ionosphere is a complex system with so many variations that
148 can arise from various dynamical changes that can be due to various changes in different
149 processes that contribute to the behavior and nature of the ionosphere. Rabiou et al., (2007)
150 affirmed that characterizing the ionosphere is of utmost importance due to the numerous

151 complexities associated with the region. The scale of these numerous complexities interestingly
152 changes at times from one day to another.

153 The concept of chaos as previously applied to ionospheric and magnetospheric studies on quiet
154 and stormy conditions are limited. Most investigations have been based on only quiet and storm
155 conditions for all studies carried out, and none of the previous works involved the use of quiet
156 and disturbed day classification of geophysical conditions until recently by Ogunsua et
157 al.,(2014), where we considered the comparative use of Lyapunov exponent and Tsallis entropy
158 as proxies for the internal dynamics of the ionosphere. This is the main reason for the
159 consideration of day to day variation of these parameters in this work.

160 **2.0 Data and Methodology**

161 The data used for this study is the global positioning system (GPS) total electron content (TEC)
162 data obtained from 5 GPS satellite receiver stations. Table 1 shows the coordinates of the
163 stations. These receivers take the measure of slant TEC within $1m^2$ columnar unit of the cross
164 section along the ray path of the satellite and the receiver which is given by

$$165 \quad STEC = \int_{receiver}^{Satellite} Ndl \quad (1)$$

166 The observation of the total number of free electron along the ray path are derived from the
167 frequency $L_1(1572.42 \text{ MHz})$ and $L_2(1227.60 \text{ MHz})$ of Global Positioning System(GPS), that
168 provide the relative ionosphere delay of electromagnetic waves travelling through the medium
169 (Saito et al.,1998). The Slant TEC is projected to vertical TEC using the thin shell model
170 assuming the height of 350m (Klobuchar,1986).

$$171 \quad VTEC = STEC \cdot \cos[\arcsin(R_e \cos\theta / R_e + h_{max})] \quad (2)$$

172 Where $R_e = 6378km$ (radius of the earth), $h_{max} = 350km$ (the vertical height assumed from
173 the satellite) and $\theta = \text{elevation angle at ground station}$

174

175 In this study, 5 GPS TEC measuring stations lying within the low latitude region were
176 considered, as shown in table 1. The TEC data obtained for January to December 2011 were

177 considered for this study and the data are given at 1min sampling time. The TEC data were
178 subjected to various analyses which will be discussed in the next section. The day to day
179 variations of the chaotic behavior and dynamical complexity were studied for the entire year.
180 The surrogate data tests for non linearity were also conducted for both the dynamical and
181 geometrical aspects.

182 **3.0 Methods of Data Analysis and Results**

183 **3.1 Time series analysis**

184 Time series can be seen as a numerical account that describes the state of a system, from which it
185 was measured. A given time series, S_n can be defined as a sequence of scalar measurement of a
186 particular quantity taken as series at different portion in time for a given time interval(Δt). The
187 time series describe the physical appearance of an entire system, as seen in Fig 1. However it
188 may not always describe the internal dynamics of that system. A system like the ionosphere
189 possesses a dominant dynamics that can be seen as diurnal so the data should be treated so as to
190 be able to see its internal dynamics. The measured TEC time series were plotted to see the
191 dynamics of the system. A typical plot of TEC usually has a dominant dynamics (see fig 1)
192 which may be seen as the diurnal behavior, however, it can also be seen that there is also a
193 presence of fluctuations (which appear to be nonlinear) in the system as a result of the internal
194 dynamics of the ionosphere and space plasma system, due to different activities in the
195 ionosphere. Therefore there is need to minimize the influence of the diurnal variations since we
196 are more interested in the nonlinear internal dynamics of the system in this study, to do so the
197 TEC time series was detrended by carrying out the following analysis below:

198 Since for the given daily data of 1minute sampling time there are 1440 data points per day. Then
199 there exists a time series t_i , where $i = 1,2,3 \dots 1440$ represents the observed time series, and
200 there also exists a set of u_i where $i = 1,2,3 \dots 1440$, such that the diurnal variation reduced time
201 is given by

$$202 \quad T_i = t_i - u_j \quad (3)$$

203 Where $i = 1, 2, 3, \dots, j = \text{mod}(i, 1440)$, if $\text{mod}(j, 1440) \neq 0$, and $j = 1440$ if $d(j, 1440) = 0$.
 204 This method will give the detrended time series represented by T_i obtained from the original
 205 TEC data as shown in fig 2. This method is similar to that used by (Unnikrishnan et al., 2006,
 206 Unnikrishnan 2010), the further explanations on the dynamical results can be found in (Kumar et
 207 al., 2004). The detrended time series were subjected to further analyses for the Phase space
 208 reconstruction and also to obtain the values of Lyapunov exponents, correlation dimension,
 209 Tsallis entropy and the implementation of surrogate data test.

210

211 **3.1.1 Phase Space reconstruction and Non Linear Time Series Analysis**

212 The study of chaoticity and dynamical complexity in a dynamical system requires a nonlinear
 213 approach, due to the fact that systems described by these phenomena can be referred to as
 214 nonlinear complex systems. The magnetosphere and the ionosphere are good examples of such
 215 systems. To be able to study such phenomena some nonlinear time series analysis can be carried
 216 out on the time series data describing such a system. The detrended time series of TEC
 217 measurement is subjected to some nonlinear time series data analysis to obtain the mutual
 218 information and false nearest neighbours, embedding dimension and delay coordinates for the
 219 phase space reconstruction, and the evaluation of other chaotic quantifiers namely: Lyapunov
 220 Exponents, Correlation dimension, recurrence analysis and Entropy.

221 The phase space reconstruction helps to reveal the multidirectional aspect of the system. The
 222 phase space reconstruction is based on embedding theorem, such that the phase space is
 223 reconstructed to show the multidimensional nature as follows:

$$224 \mathbf{Y}_n = (\mathbf{s}_{n-(m-1)\tau}, \mathbf{s}_{n-(m-2)\tau}, \dots, \mathbf{s}_{n-\tau}, \mathbf{s}_n) \quad (4)$$

225 where \mathbf{Y}_n are vector in phase space. The proper choice of embedding dimension (m) and delay
 226 Time (τ) are essential for phase space reconstruction (Fraser and Swinney, 1986; Kennel et
 227 al., 1992).

228 If the plot showing the time delayed mutual information shows a marked minimum that value
 229 can be considered as a responsible time delay; Fig 3 shows the mutual information plotted

230 against time delay. Likewise, the minimal embedding dimension, which correspond to the
 231 minimum number of the false nearest neighbours can be treated as the optimum value of
 232 embedding dimension in (Unnikrishnan et al.,2006, Unnikrishnan, 2010). A plot of fraction of
 233 false nearest neighbours against embedding dimension can be seen in Fig 4. It was observed that
 234 for all the daily detrended TEC time series the choice of $\tau \geq 30$ and $m \geq 4$ values of delay and
 235 embedding dimension above these values are suitable for analysis of data for all stations. The
 236 choice of $\tau = 30$ and $m = 5$ were mostly used to analyze the dynamical aspects for all the
 237 stations. The reconstructed Phase space trajectory is shown in Fig 5

238 3.1.2 Lyapunov Exponents

239 The Lyapunov exponent has been a very important quantifier for the determination of chaos in a
 240 dynamical system. This quantifier is also used for the determination of chaos in time series,
 241 representing natural systems like the ionosphere and magnetosphere (Unnikrishnan 2008, 2010).
 242 A positive Lyapunov exponent indicates divergence of trajectory in one dimension, or alternative
 243 an expansion of volume, which can also be said to indicate repulsion, or attraction from a fixed
 244 point. A positive Lyapunov exponent indicates that there is evidence of chaos in a dissipative
 245 deterministic system, where the positive Lyapunov exponent indicates divergence of trajectory in
 246 one direction or expansion of value and a negative value shows convergence at trajectory or
 247 contraction of volume along another direction.

248 The largest Lyapunov exponent (λ_1) can be used to determine the rate of divergence as indicated
 249 by (Wolf et al.,1985)

250 Where

$$251 \lambda_1 = \lim_{r \rightarrow \infty} \frac{1}{t} \ln \frac{\Delta x(t)}{x(0)} = \lim_{r \rightarrow \infty} \frac{1}{t} \sum_{i=1}^t \ln \left(\frac{\Delta x(t_i)}{\Delta x(t_{i-1})} \right) \quad (5)$$

252 The Lyapunov exponent was computed for the TEC values measured from Different stations.
 253 The evolution in state space was scanned with $\tau = 30$, $m = 5$, is shown in fig 6. The day to day
 254 variations of the Lyapunov exponent was computed for the entire year to so as to study the
 255 annual trend of variation. This was implemented using the method introduced by Rosenstein
 256 (1993), and Hegger et al., (1994), both algorithms use very similar methods. Lyapunov

257 exponents were also computed for varying time delay at constant embedding dimension and also
 258 for varying embedding dimension, to check for the stability with changes in trajectory. These can
 259 be seen in fig. 6b and 6c. The day to day values of Lyapunov exponent plotted for the Enugu
 260 station and for Toro station are shown in fig 7a to 7b. The plots of the day to day values show the
 261 transient variation of the ionosphere and a wavelike yearly pattern.

262 **3.1.3 Correlation Dimension**

263 Another relevant method to study the underlying dynamics or internal dynamics of a system is to
 264 evaluate the dimension of the system. The correlation dimension gives a good approximation of
 265 this as suggested by Grassberger and Procaccia (1983a; b). The correlation dimension is
 266 preferred over the box counting dimension because it takes into account the density of points on
 267 the attractor (Strogatz 1994). The correlation dimension D is defined as

$$268 \quad D = \lim_{r \rightarrow 0} \frac{\ln C(r)}{\ln r} \quad (6)$$

269 The term $C(r)$ is the correlation sum for radius (r) where for a small radius (r) the correlation
 270 sum can be seen as $C(r) \sim r^d$ for $r \rightarrow 0$. The correlation sum is dependent of the embedding
 271 dimension (m) of the reconstructed phase space and it is also dependent of the length of the time
 272 series N as follows

$$273 \quad C(r) = \frac{2}{N(N-1)} \sum_{i=1}^N \sum_{j=i+1}^N \Theta(r - \|y_i - y_j\|) \quad (7)$$

274 Where Θ is the Heaviside step function, with $\Theta(H) = 0$ if $H \leq 0$ and $\Theta(H) = 1$ for $H > 0$.

275 The correlation dimension was computed using the Theiler algorithm approach, with Theiler
 276 window (w) at 180. The Theiler window was chosen to be approximately equal to the product of
 277 m and τ . A similar approach to the computation of correlation dimension was used by
 278 Unnikrishnan and Ravindran (2010) to determine the correlation dimension of detrended TEC
 279 data for some stations in India which lies within the equatorial region, like Nigeria. Ogunsua et
 280 al., (2014) also used similar methods for some detrended TEC from Nigerian stations.

281 The correlation dimension for data taken for the quietest day of October 2011 and the most
 282 disturbed day of October 2011 from Birnin Kebbi GPS TEC measuring station were represented

283 by Fig 8a and Fig 8b respectively. The correlation dimension saturates at $m \geq 4$ for the quietest
284 day of the month and at $m \geq 5$ for the most disturbed day. In this illustration the most disturbed
285 day of this month fall within the storm period of October 2011. The classification of days into
286 quiet and disturbed days in the month of October 2011 enables us to compare the quiet and storm
287 periods together while comparing the quiet days with some relatively disturbed days.

288 **3.1.4 Computation of Tsallis Entropy and Principles of Nonextensive Tsallis Entropy**

289 Entropy measures are very important statistical techniques that can be used to describe the
290 dynamical nature of a system. The Tsallis entropy can be used to describe the dynamical
291 complexity of a system and to also understand the nonlinear dynamics like chaos which may
292 exist in a natural system. The use of entropy measure as a method to describe the state of a
293 physical system has been employed into information theory for decades. The computation of
294 entropy allows us to describe the state of disorderliness in a system, one can generalize this same
295 concept to characterize the amount of information stored in more general probability
296 distributions (Kantz & Shrieber 2003, Balasis et al.,2009). The concept of information theory is
297 basically concerned with these principles. The information theory gives us an important
298 approach to time series analysis. If our time series which is a stream of numbers, is given as a
299 source of information such that this numbers are distributed according to some probability
300 distribution, and transitions between numbers occur with well-defined probabilities. One can
301 deduce same average behaviour of the system at a different point and for the future. The term
302 entropy is used in both physics and information theory to describe the amount of uncertainty or
303 information inherent in an object or system (Kantz and schrieber 2003). The state of an open
304 system is usually associated with a degree of uncertainty that can be quantified by the
305 Boltzmann-Gibbs entropy, a very useful uncertainty measure in statistical mechanics. However
306 Boltzmann-Gibbs entropy cannot, describe non-equilibrium physical systems with large
307 variability and multifractal structure such as the solar wind (Burgala et al., 2007, Balasis et al.,
308 2008). One of the crucial properties of the Boltzmann-Gibbs entropy in the context of classical
309 thermodynamics is extensivity, namely proportionality with the number of elements of the
310 system. The Boltzmann-Gibbs entropy satisfies this prescription if the subsystems are
311 statistically (quasi-) independent, or typically if the correlations within the system are essentially
312 local. In such cases the system is called extensive. In general however, the situation is not of this

313 type and correlations may be far from negligible at all scales. In such cases, the Boltzmann-
 314 Gibbs entropy is nonextensive (Balasis et. al., 2008, 2009). These generalizations above were
 315 proposed by Tsallis (1988), who was inspired by the probabilistic description of multifractal
 316 geometries. Tsallis (1988, 1998) introduced an entropy measure by presenting an entropic
 317 expression characterized by an index q which leads to a nonextensive statistics,

$$318 \quad S_q = k \frac{1}{q-1} \left(1 - \sum_{i=1}^W p_i^q \right) \quad (8)$$

319 Where p_i are the probabilities associated with the microscopic configurations, W is their total
 320 number, q is a real number, and k is Boltzmann's constant. The value q is a measure of the
 321 nonextensivity of the system: $q \rightarrow 1$ corresponds to the standard extensive Boltzmann-Gibbs
 322 statistics. This is the basis of the so called nonextensive statistical mechanics, which generalizes
 323 the Boltzmann-Gibbs theory. The entropic index q characterizes the degree of nonadditivity
 324 reflected in the following pseudoadditivity rule:

$$325 \quad \frac{S_q(A+B)}{k} = \left[\frac{S_q(A)}{k} \right] + \left[\frac{S_q(B)}{k} \right] + (1 - q) \left[\frac{S_q(A)}{k} \right] \left[\frac{S_q(B)}{k} \right]. \quad (9)$$

326 The cases $q > 1$ and $q < 1$, correspond to subadditivity (or subextensivity) and superadditivity
 327 (or superextensivity), respectively and $q = 1$ represents additivity (or extensivity). For
 328 subsystems that have special theory probability correlations, extensivity is not a valid for
 329 Boltzmann-Gibbs entropy in such cases, but may occur for S_q with a particular value of the index
 330 q . Such systems are sometimes referred to as nonextensive (Boon and Tsallis, 2005, Balasis et al
 331 2008, 2009). The parameter q itself is not a measure of the complexity of the system, but
 332 measures the degree of nonextensivity of the system. It is the time variations of the Tsallis
 333 entropy for a given $q(S_q)$ that quantify the dynamic changes of the complexity of the system.
 334 Lower S_q values characterize the portions of the signal with lower complexity. In this
 335 presentation we estimate S_q on the basis of the concept of symbolic dynamics and by using the
 336 technique of lumping (Balasis et al. 2008, 2009).

337 A comparison of Tsallis entropy with Lyapunov exponents computed for the same set of data has
 338 been carried out in this work, to see the efficacy of the combined usage of both parameters. This
 339 is based on the established facts that variations in the values of Tsallis entropy can be linked with

340 that of Lyapunov exponents chaotic behavior in systems as seen in (Baranger et al., 2012;
341 Anastasiadis et al., 2005; Kalogeropoulos et al., 2012;2013). Coraddu et al., (2005) showed the
342 Tsallis entropy generalization for Lyapunov exponents. Further details can be found in Ogunsua
343 et al., (2014),

344 they were able to investigate the similarities in their response to the complex dynamics of the
345 ionosphere, and this informs the further use of the two quantities as indices to study the day to
346 day variation of ionospheric behaviour in this work.

347 The values of these entropy measures were also computed in order to study the dynamical
348 complexity of the system under observation (the ionosphere). The day to day values of Tsallis
349 entropy were computed for the entire year for different stations. The day to day values of Tsallis
350 entropy plotted for the Enugu station and for Toro station are shown in fig 9(a and b). The plots
351 of the day to day values show the transient variation of the ionosphere and a wavelike yearly
352 pattern.

353 **3.2 Non linearity Test using surrogate data**

354 The test for non-linearity using the method of surrogate data according to Kantz and Schreiber
355 (2003) has proven to be a good test for non-linearity in time series describing a system. It has
356 been accepted that the method of surrogate data test could be a successful tool for the
357 identification of nonlinear deterministic structure in an experimental data (Pavlos et al., 1999).
358 This method involves creating a test of significance of difference between linearly developed
359 surrogate and original nonlinear time series to be tested. The test is done by carrying out the
360 computation of the same quantity on both surrogates and the original time series and then
361 checking for the significance of difference between the results obtained from the surrogates with
362 the original data. Theiler et al (1992) suggested the creation of surrogate data by using Monte
363 Carlo techniques for accurate results. According to this method, typical characteristic of data
364 under study are compared with those of stochastic signals (surrogates), which have the same
365 auto-correlation function and the power spectrum of the original time series. It can be safely
366 concluded from the test of significance carried out on the surrogate and the original data that, a
367 stationary linear Gaussian Stochastic model cannot describe the process under study provided
368 that the behaviour of the original data and the surrogate data are significantly different.

369

370 In this work 10 surrogate data were generated from the original data set. The geometrical and
371 dynamical characteristics of the original data were then compared to that of the surrogates using
372 the statistical method of significance of difference which can be defined as

373

$$374 \quad S = \frac{\alpha_{Surr} - \alpha_{Original}}{\sigma} \quad (13)$$

375

376 Where α_{Surr} is the mean value of the computed quantity for the surrogate data and $\alpha_{Original}$ is
377 the same quantity computed for the original TEC data, σ is the standard deviation of the same
378 quantity computed for the surrogate data. The significance of difference considered for the null
379 hypothesis to be rejected here is greater than 2, which enables us to be able to reject the null
380 hypothesis that the original TEC data describing the ionospheric system can be modeled using a
381 Gaussian linear stochastic model with confidence greater than 95%.

382 The surrogate data test for all stations used in this study show that the Lyapunov exponent of
383 the surrogate data for the selected days in October are shown in the Table below The results
384 show that the surrogate data test for Lyapunov exponent show a significance of difference
385 greater than 2 for all the selected days for all the stations. Similar results were obtained for
386 Mutual Information, Fraction of False Nearest Neighbours and Correlation Dimension. This
387 result gives us the confidence to reject the null hypothesis that the data used cannot be modeled
388 using a linear Gaussian stochastic model, which shows that the system is a nonlinear system with
389 some level of determinism. Fig. 10 shows the plots comparing the mutual information plotted
390 against time delay for the original detrended data blue with the mutual information for the
391 surrogate data for TEC data measured at Lagos for the quietest day of March 2011, while Fig. 11
392 is comparing fraction of false nearest neighbours for the same set of data. Tables 2a shows the
393 values of Lyapunov exponents for both original detrended and its surrogate data for TEC
394 measured in Lagos during the quietest days and Table 2b shows the values of Lyapunov
395 exponents for both original detrended and its surrogate data for TEC measured in Lagos during
396 the most disturbed days of October 2011.

397 **3.3 Trend filtering using the moving average approach for the daily Values**

398 The trend of a fluctuating time series can be made clearer to reveal the general pattern of that
 399 time series, and to make the fluctuating pattern of the daily variation of the chaoticity and
 400 dynamical complexity measures clearer in the work, the moving average method has been
 401 employed. The method of moving average filtering has found its applications geophysics (e.g.
 402 Bloomfield 1992; Bloomfield and Nychka 1992; Baillie and Chung 2002), and in other areas like
 403 financial time series analysis, microeconomics, biological sciences and medical sciences. The
 404 various fields mentioned require different trend filtering method depending on the structure of
 405 the time series to be analyzed. Different filtering processes that can be used to reveal the trend
 406 includes the moving average filters, exponential filters, band-pass filtering, median filtering etc.

407 Suppose we have a time series $z[t]$ such that $t = 1,2,3 \dots \dots n$, where 'n' could assume any
 408 value. If $z[t]$ consists of a consistent varying trend component that appears over a longer period
 409 of time t given as $u[t]$ and a more rapidly varying component $v[t]$. The goal of trend filtering in
 410 any research is to estimate either of the two components (Kim et al., 2009). The purpose of trend
 411 filtering in this work is to further reveal the general slow varying trend that appears to be obvious
 412 in the daily variation of the values of the chaoticity and dynamical complexity of the ionosphere,
 413 which might appear to be obviously varying with the yearly solar activity (a quantity with slow
 414 varying trend). To make $u[t]$ which represents the general slow varying trend smoother and in
 415 the process reduce $v[t]$ we apply the moving average filter.

416 If we assume $z[t]$ to be our time series representing the daily variation of the values of the
 417 chaoticity and dynamical complexity of the ionosphere, then our smoothing with weighting
 418 vector/filter w_j will create the new sequence u_j as

$$419 \quad u[t] = z[t] * w[n] = \frac{1}{2k+1} \sum_{i=-k}^k x[n - 1]. \quad (14)$$

420 In this work the Savitzky-Golay method of smoothing proposed by Savitzky and Goley (1967),
 421 which is a generalized form of moving average was applied to the trend smoothing of the daily
 422 variation of the chaoticity and dynamical complexity of the ionosphere. In this case it performs a
 423 least square fit to a small set of $L(= 2k + 1)$ consecutive data to a polynomial and then takes
 424 midpoint of the polynomial curve as output. The smoothed time series in this work will now be
 425 given as

426 $u[t] = z[t] * \omega[n] = \frac{\sum_{i=-k}^k A_i * x[n-1]}{\sum_{i=-k}^k A_i}$ (15)

427 where, $\omega[n] = \frac{A_n}{\sum_{i=-k}^k A_i}$, $-k \leq n \leq k$ such that A_i controls the order of polynomial. A similar
 428 method was described in Reddy et al., (2010).

429 The smoothed daily variation and the original data and the plot of the smoothed variation only,
 430 for the Lyapunov exponents of the detrended TEC measured at the Enugu and Toro are shown in
 431 fig 12(a and b). The smoothed day to day variation for Tsallis entropy for the detrended TEC
 432 measured at Enugu and Toro stations respectively are shown in fig 13(a and b).

433 **4.0 DISCUSSION**

434 The results presented in the work reveals the dynamical characteristics of the ionosphere. These
 435 characteristics are being discussed in this section, considering the time series treatment and phase
 436 space reconstruction; the study of chaos using chaotic quantifiers and the use and comparison of
 437 dynamical complexity measures in terms of their response to the variations on ionospheric
 438 dynamics. Also being discussed, is the implication of the nonlinearity test using the surrogate
 439 data and the comparison of the two quantifiers and their viability as indices for the continuous
 440 study and characterization of the ionosphere

441
 442 The time series analysis shows the appearance of some degree of nonlinearity in the internal
 443 dynamics of the ionosphere. The time series plot in Fig. 1 shows the rise in TEC to peak at the
 444 sunlit hours of the day, however it can be seen that the rising to the peak exhibited by the
 445 ionosphere, which is the dominant dynamics during the day, make it impossible to clearly see the
 446 internal dynamics of the system from the TEC time series plot. It can be seen that the TEC time
 447 series curve is not a smooth curve with tiny variations, which probably describes a part of the
 448 internal dynamics. These visible tiny variations around the edges of the time series plot can be
 449 regarded as rate of change of TEC which is a phenomenon that can describe the influence of
 450 scintillations in the ionosphere these variations are however more obvious during the night time
 451 between 1100th and 1440th minutes of the day (that is, between about 1800 and 2400 hours of
 452 the day). It should be noted here that scintillations has been described as a night time phenomena
 453 associated with spread-F, and it occurs around pre-midnight and post-midnight periods (Vyas

454 and Chandra 1994; Vyas and Dayanandan 2011; Mukherjee et al.,2012; Bhattacharyya and
455 Pandit 2014). The detrended data shows the internal dynamics of the system more clearly, with a
456 pattern similar to the values around night period mentioned earlier. The post-sunset values
457 (especially at night time) in Fig.1 show a pattern similar pattern with the detrended TEC plot in
458 Fig 2. It has been established that TEC does not decrease totally throughout the night as expected
459 normally through simple theory that TEC builds up during the day, but it shows some anomalous
460 enhancements and variations and this can occur under a wide range of geophysical conditions
461 (Balan and Rao, 1987; Balan et al., 1991;Unnikrishnan and Ravindran, 2010). The delay
462 representation of the phase space reconstruction shows a trajectory that is clustered around its
463 origin, for all the stations, which can be seen as an indication of the possible presence of chaos.
464 The degree of closeness of these trajectories however varies for different days from one station
465 to another, resulting from varying degrees of variations in stochasticity and determinism. The
466 varying degrees of variations in stochasticity and determinism can be attributed to the daily
467 variations and local time variations of photoionization, recombination, influx of solar wind and
468 other factors that may influence the daily variations of TEC (Unnikrishnan 2010).

469
470 The positive values of Lyapunov exponent indicate the presence of chaos (Wolf et al., 1985;
471 Rosenstein et al., 1993; Hegger et al., 1999; Kantz and Schreiber, 2003). The presence of chaos
472 was revealed by the positive Lyapunov exponent computed from all stations and this as a result
473 of the fact that the ionosphere is a system controlled by many parameters influencing its internal
474 dynamics. Because of its extreme sensitivity to solar activity, the ionosphere is a very sensitive
475 monitor of solar events. The ionospheric structure and peak densities in the ionosphere vary
476 greatly with time (sunspot cycle, seasonally and diurnally), with geographical location (polar,
477 auroral zones, mild-latitudes, and equatorial regions), and with certain solar-related ionospheric
478 disturbances. During and following a geomagnetic storm, the ionospheric changes around the
479 globe, as observed from ground site can appear chaotic (Fuller-Rowell et al., 1994; Cosolini and
480 Chang, 2001; Unnikrishnan and Ravindran, 2010). The recorded presence of chaos as indicated
481 by the positive values of Lyapunov exponent was found in all the computations, for all the TEC
482 values obtained for the selected days from all the measuring stations used in this work. This can
483 be expected as it agrees with results from previous works that show that there is a reasonable
484 presence of chaos in the ionosphere, even in the midst of the influence of stochastic drivers like

485 solar wind (Bhattacharyya, 1990; Wernik and Yeh, 1994; Kumar et al., 2004; Unnikrishnan et
486 al., 2006a,b; Unnikrishnan, 2010). However the values of Lyapunov exponents vary from day to
487 day due to variations in ionospheric processes for different days on the same latitude as seen in
488 Fig. 7(a and b) with Fig. 12(a and b) showing the day to day variation (upper panel) and the
489 smoothed curve of the day to day variation (lower panel) for the entire year. There are also
490 latitudinal variations due to spatial variations in the various ionospheric processes taking place
491 simultaneously. The ionosphere is said to have a complex structure due to these varying
492 ionospheric processes.

493 The higher values of Lyapunov exponent during months of low solar activity (the solstices) is an
494 evidence that that the rate of exponential growth in infinitesimal perturbations in the ionosphere
495 leading to chaotic dynamics might be of higher degree during most of the days of those months
496 compared to days of the months with high solar activities showing lower values of Lyapunov
497 exponents (Unnikrishnan 2010, Unnikrishnan and Ravindran, 2010).

498
499 The results of the correlation dimension values computed are within the range of 2.7 to 3.2 with
500 the lower values occurring mostly during the storm periods. The lower dimension during the
501 storm periods compared to the quiet days may be due to the effect of a stochastic drivers like
502 strong solar wind and solar flares, that occurs during geomagnetic storms on the internal
503 dynamics of the ionosphere, this could have been as a result of the fact that the internal dynamics
504 must have been suppressed by the external influence. The restructuring of the internal dynamics
505 of the ionosphere might be responsible for low dimension chaos during storm and also the lower
506 values of other measures like the Lyapunov exponents. The relatively disturbed day however
507 might have a higher dimension so long as it is not a storm period, and sometimes a relatively
508 disturbed day of the month might be a day with storm and in this case there is usually a lower
509 value of chaoticity and sometimes lower values of correlation dimension as well. The lower
510 value of chaoticity and dimension in ionosphere during storms indicates a phase transition from
511 higher values during the quiet periods to lower values during storm periods which may be due to
512 the modification of the ionosphere by the influx of high intensity solar wind during the storm
513 period (Unnikrishnan et al., 2006a, b; Unnikrishnan 2010; Unnikrishnan and Ravindran, 2010).

514

515 The surrogate data test shows significance of difference greater than 2 for all the computed
516 measures which enables rejection of the null hypothesis that the ionospheric system can be
517 represented with a linear model for all the data used from the stations. However it was
518 discovered that the lower significance of difference corresponds to the lower values of Lyapunov
519 exponents during storm and extremely disturbed periods (see tables 2 and 3). This may be due
520 the rise in stochasticity during the storm period as a result of drop in values of computed
521 quantities like Lyapunov exponents. Our ability to reject the Null hypothesis for all stations
522 however shows the presence of determinism and confirms that the underlying dynamics of the
523 ionosphere is mostly non-linear. This further validates the presence of chaos since the surrogate
524 data test for non-linearity show that out detrended TEC is not a Gaussian (linear) stochastic
525 signal (Unnikrishnan 2010).

526

527

528 The Tsallis entropy was able to show the deterministic behavior of the ionosphere considering
529 its response during storm periods compared to other relatively quiet periods as the rapid drop in
530 values of Tsallis entropy during storm show that there is a transition from higher complexity
531 during quiet period to lower complexity during storms, this response in the values of Tsallis
532 entropy is similar to the response of Lyapunov exponent values during storm. This reaction to
533 storm shown by the values of Tsallis entropy computed for TEC was also described by the
534 reaction of Tsallis entropy computed for Dst during storm periods (Balasis et al., 2008, 2009). A
535 closer observation of the day-to-day variability within a month shows that the values were much
536 lower for storm periods compared to the nearest relative quiet period. For example, the storm
537 that occurred on the 25th of October resulted in lower values of Lyapunov exponent and Tsallis
538 entropy compared to relatively quiet days close to it. The reaction to storm may be due to the
539 influence of stochastic driver like strong solar wind flowing into the system as a result of solar
540 flare or CMEs that produces the geomagnetic storms. Although there is always an influence of
541 corpuscular radiation in form of solar wind flowing from the sun into the ionosphere, the
542 influence is usually low for days without storm compared to days with geomagnetic storms as a
543 result of solar flares, CMEs etc (Unnikrishnan et al., 2006a,b; Unnikrishnan, 2010, Ogunsua et al
544 2014).

545

546 The presence of chaos and high variations in the dynamical complexity, even at quiet periods in
547 the ionosphere may be due to the internal dynamics and inherent irregularities of the ionosphere
548 which exhibit non-linear properties. However, this inherent dynamics may be complicated by
549 external factors like Geomagnetic storms. This may be the main reason for the drop in the values
550 of Lyapunov exponent and Tsallis entropy during storms. According to Unnikrishnan et al.,
551 (2006a,b), geomagnetic storms are extreme forms of space weather, during which external
552 driving forces , mainly due to solar wind, subsequent plasmasphere -ionosphere coupling, and
553 related disturbed electric field and wind patterns will develop. This in turn creates many active
554 degrees of freedom with various levels of coupling among them, which alters and modifies the
555 quiet time states of ionosphere, during a storm period. This new situation developed by a storm,
556 may modify the stability/instability conditions of ionosphere, due to the superposition of various
557 active degrees of freedom.

558

559 The observation from the day-to-day variability of Lyapunov exponent and Tsallis entropy also
560 show irregular pattern for all stations. These irregular variations might be due to the same factors
561 mentioned before (i.e internal irregularities due to so many factors described and also due to
562 variation in the influx of the external stochastic drivers). The day-to-day variability for the entire
563 year shows a “wavelike” pattern with the values dropping to lower values during the equinox
564 months especially during March-April equinox. This can be seen as a form of semiannual
565 variation, possibly resulting from the higher energy inputs during equinoxes. This is because
566 solar wind is maximized at the equinoxes which might result in higher energy input that will
567 eventually suppress the internal dynamics to give lower values of chaoticity. The modification of
568 the ionosphere as a result of the higher energy input resulting from the maximized influx of solar
569 wind has been reported to be responsible for the lower values of chaoticity when averagely
570 compared to the days of the year with lower solar wind inputs as reported by Unnikrishnan et al.,
571 2006; 2010; Ogunsua et al., 2014. The semiannual pattern has been found to be similar for
572 different stations as seen in Figs. 7 & 12 and Figs. 9 & 13 for Lyapunov exponents and Tsallis
573 entropy respectively. Figs.9 and 13 show the smoothed curves for Lyapunov exponent and
574 Tsallis entropy respectively, with the drop in values at equinoxes showing more clearly. The
575 phase transition in chaoticity and dynamical complexity is also responsible for the wavelike
576 variations, with values of Lyapunov exponent and Tsallis entropy dropping during the equinoxial

577 months, and this may be due to the influence of the daily influx of the solar wind having higher
578 values during equinoxes due to the proximity of the Earth to the sun during this period compared
579 to the solstice months.

580
581 The wavelike pattern observed has been described to be as a result of self organized critical
582 (SOC) phenomenon, a phenomenon which has been found to exist in both the magnetosphere
583 and the ionosphere or the space plasma system in general, due to coupling between the two
584 systems, since the magnetosphere couples the ionosphere tightly to the solar wind (Lui, 2002).
585 Many literatures has shown the existence of chaos in the SOC in the magnetosphere (chang et
586 al.,1992,1998,1999; Consolini et al., 1996 Chapman et al., 1998; Freeman and Watkins 2002
587 and; Koselov and Koselova, 2001. Uritsky et al., (2003) and; Chang et al., (1992). The existence
588 of SOC in space plasma system involving both the ionosphere the the magnetosphere was
589 described by (Lui, 2002; Chang et al., 2002; Chang et al., 2004)

590
591 The variation along the latitude also shows the inconsistence and complexity of the ionospheric
592 processes. This is the reason why for the same day of the month the values of Lyapunov
593 exponent vary from one station to another. Lyapunov exponent however, appears to respond
594 better to changes in solar activities compared to Tsallis entropy with more distinct results. This
595 may be due to the fact that Tsallis entropy being not only a measure of complexity, but also a
596 measure of disorderliness in a system might not be as perfect in describing chaos as Lyapunov
597 exponent. Kalogeropoulos (2009) and Baranger et.al (2002) observed that Tsallis entropy has a
598 relationship that is not totally linear in all cases at different level of chaos with Lyapunov
599 exponent as a measure of chaos.

600
601 There are also many variations in the internal dynamics of the ionosphere that could lead to
602 changes in chaotic behavior. The variations of Lyapunov exponents during quiet days might be
603 as a result of different variations in the intrinsic dynamics of the ionosphere. Difference in
604 variation pattern at different stations for the same quiet day might also be due to the same reason.
605 It can be affirmed that the ionosphere is a complex system that varies with a short latitudinal or
606 longitudinal interval such that even stations with one or two degrees of latitudinal differences
607 might record different values on the same day for both quiet and disturbed periods and that the

608 same might also occur for storm periods. This is illustrated by the different pattern of variation of
609 TEC recorded from different stations within such a close range as used in this study.

610
611 These Latitudinal variation in the values of Lyapunov exponents and Tsallis entropy can be
612 further described by the behavior of the TEC because there can be a more sporadic rate of
613 change in TEC as seen in the time series plots as a result of irregularities in the internal dynamics
614 of the ionosphere, which might be as a result of plasma bubbles. Irregularities develop in the
615 evening hours at F region altitudes of magnetic equator, in the form of depletions, frequently
616 referred to as bubbles. The edges of these depletions are very sharp resulting in large time rate of
617 TEC in the equatorial ionosphere, even during magnetically quiet conditions. The large gradient
618 of the equatorial ionization persists in the local post-sunset hours till about 2100 h LT.
619 (DasGupta et al., 2007; Unnikrishnan and Ravindran, 2010). The TEC data for one station might
620 experience an extremely sharp rate of change in TEC that may be due to some plasma bubbles in
621 that region while the TEC from the other station stays normal. These variations in the various
622 internal dynamics like plasma bubbles leading to scintillation can cause variations in the
623 dynamical response of the TEC. Hence, the irregular variation in the values of the Lyapunov
624 exponent and Tsallis entropy even in quiet periods for two relatively close stations may be due to
625 these irregularities. This might also be responsible for the quiet days in the same station having
626 lower values of Lyapunov exponent compared to higher values recorded for disturbed days
627 without the external influence of storms.

628
629 The variations of these chaos and dynamical complexity parameters might also be as a result of
630 the anomalous TEC enhancements that might occur at nights (Balan and Rao (1987); Balan et
631 al., 1991). These effects can also be seen more clearly in the Tsallis entropy values for the five
632 period window for quiet day of January, 2011, because the night time value is higher and it also
633 show a much higher series of fluctuations during this period compared to other periods. As
634 mentioned in Unnikrishnan and Ravindran (2010), the irregular changes in the dynamical
635 characteristics of TEC from the results of Lyapunov exponent and Tsallis entropy also may be
636 due to the collisional Raleigh-Taylor instability which may give rise to a few large irregularities
637 in L band measurements (Rama Rao et al., 2006; Sripathi et al., 2008) all these can be seen as
638 internal factors responsible for variations in the dynamical response of TEC as recorded from the

639 values of the Lyapunov exponents and Tsallis entropy completed for days without storm which
640 might be quiet or disturbed according to classification and also could account for higher values
641 of these qualifiers during disturbed days compared quiet days. During storms however, the
642 values were much lower

643

644 Earlier we, (Ogunsua et al., 2014) showed the appearance and variation of chaoticity quiet and
645 disturbed day classification by international most quiet day (IQD) and internal most disturbed
646 day (IDD) classification, as compared to quiet and storm period used by Unnikrishnan (2006;
647 2010). We were able establish that a relatively quiet day may be less chaotic compared to a
648 relatively disturbed day unlike the result presented by Unnikrishnan (2006; 2010) for quiet and
649 storm period. Also the combined use of both Lyapunov exponent and Tsallis entropy for the first
650 time was found to have a high correlation mostly above 80%, which has stimulated the interest
651 for further research using the two diagnosis for the study of ionospheric dynamics.

652

653 This work on the other hand presents the results for day to day variation and has revealed a
654 seasonal trend for both Lyapunov exponents and Tsallis entropy, which appear in wavelike in
655 form, with troughs during the two equinoxes. This was established for different stations used in
656 this research work. The results show the appearance of seasonal trend in spite of the sporadic
657 daily variation resulting from various changes in the internal dynamics. The seasonal trend has
658 provided another possible evidence of higher energy input during equinoxes, since it reveals the
659 effect of the annual energy input to the ionosphere. The day to day response these parameters has
660 also revealed the variations in the underlying dynamics of the system.

661

662 As a similarity between the present work and Ogunsua et al. (2014) the relationship between
663 Lyapunov exponent and Tsallis entropy can also be seen from this work, as the two quantifiers
664 exhibit similarities in their response to the dynamical behavior of the ionosphere with phase
665 transition at the same periods of time for all stations. A further investigation of this relationship
666 shows that all the daily values of Tsallis entropy correlates positively with the values of
667 Lyapunov exponent at values between 0.78 and 0.83.

668

669 The ability of these quantifiers to clearly reveal the ionospheric dynamical response to solar
670 activities and changes in its internal dynamics due to other factors is a valid proof of the
671 authenticity of the use of these chaotic and dynamical measures, as indices for ionospheric
672 studies.

673 **5.0 Conclusion**

674 The chaotic behaviour and dynamical complexity of low latitude ionosphere over some parts of
675 Nigeria was investigated using TEC time series measured Simultaneously at five different
676 stations namely Birnin Kebbi (geographic coordinates $12^{\circ}32'N$, $4^{\circ}12'E$; dip latitude $0.62^{\circ}N$),
677 Torro (geographic coordinates $10^{\circ}03'N$, $9^{\circ}04'E$; dip latitude $-0.82^{\circ}N$), Enugu (geographic
678 coordinates $6^{\circ}26'N$, $7^{\circ}30'E$; dip latitude $-3.21^{\circ}N$), Lagos (geographic coordinates $6^{\circ}27'N$,
679 $3^{\circ}23'E$; dip latitude $-3.07^{\circ}N$) and Yola (geographic coordinates $9^{\circ}12'N$, $12^{\circ}30'E$; dip
680 latitude $-1.39^{\circ}N$) within the low latitude region. The detrended TEC time series data obtained
681 from the GPS data measurement were analysed using different chaoticity and dynamical
682 complexity parameters.

683 The evidence of the presence of chaos in all the time series data was obtained for all the data
684 used, as indicated by the positive Lyapunov exponent. The results of Tsallis entropy show the
685 variations in the dynamical complexity of the ionosphere, which may be due to geomagnetic
686 storms and other phenomena like changes in the internal irregularities of the ionosphere. The
687 response of the Tsallis entropy to various changes in the ionosphere also shows the deterministic
688 nature of the system. The results of the Tsallis entropy show a lot of similarities with that of the
689 Lyapunov exponents between 0.78 and 0.81, with both results showing a phase transition from
690 higher values in the solstices to lower values during the equinoxial months. The values of
691 Lyapunov exponent were found lower for the days of the months in which storm was recorded
692 relative to the nearest relatively quiet days which agree with previous works by other
693 investigators. A similar pattern of results was obtained for the computed values of Tsallis
694 entropy. The random variations in the values of chaoticity in the detrended TEC describing the
695 internal dynamics of the ionosphere as seen in the result obtained from both Lyapunov exponent
696 and Tsallis entropy depicts the ionosphere as a system with a continuously changing internal
697 dynamics, which shows that the ionosphere is not totally deterministic but also has some
698 elements of stochasticity influencing its dynamical behaviour.

699

700 The phase transition in the systems of the ionosphere resulting in the lower values of the
701 chaoticity and dynamical complexity quantifiers during the geomagnetic storms and the
702 equinoxial months is the evidence that the ionosphere can be greatly modified by stochastic
703 drivers like solar wind and other incoming particle systems. The drop in values during equinoxes
704 can be seen as form of semiannual variation, a phenomenon peculiar to the low latitude regions.

705

706 Although the knowledge of being able to characterize the ionospheric behaviour using the two
707 major quantifiers shows their ability to measure level of determinism when used together, the
708 relationship between these two quantifiers calls for more research, in the use of these qualifiers,
709 to enable proper description and characterization of the state of ionosphere. The response of both
710 Tsallis entropy and Lyapunov exponents to changes in the ionosphere shows that the two
711 quantifiers can be used as indices to describe the processes/dynamics of the ionosphere.

712

713 Even though we cannot conclude totally until further investigations have been carried out on
714 various properties of the ionosphere describing its dynamics. It can be safely established that this
715 study has created roadmap for the use of the chaoticity and dynamical complexity measures as
716 indices to describe the process/dynamics of the ionosphere.

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725 **Acknowledgement**

726 The authors appreciate the editorial team and the referees for their contributions which have led
727 to the final shape of this paper. The GPS data used for this research were obtained from the public
728 archives of the Office of the Surveyor General of the Federation (OSGoF) of the Federal
729 Government of Nigeria, which is the mapping agency of Nigeria.

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920 Table 1: Coordinates of the GPS stations

Station Name	Geographic Coordinates		Dip latitude ($^{\circ}N$)
	Long ($^{\circ}E$)	Lat($^{\circ}N$)	
Birnin Kebbi	$4^{\circ} 12' E$	$12^{\circ} 32' N$	$0.62^{\circ} N$
Torro	$9^{\circ} 04' E$	$10^{\circ} 03' N$	$-0.82^{\circ} N$
Yola	$12^{\circ} 30' E$	$9^{\circ} 12' N$	$-1.39^{\circ} N$
Lagos	$3^{\circ} 23' E$	$6^{\circ} 27' N$	$-3.07^{\circ} N$
Enugu	$7^{\circ} 30' E$	$6^{\circ} 26' N$	$-3.21^{\circ} N$

928 Table 2a : Results of Surrogate data test for Lyapunov exponent for TEC data for the quietest
 929 days of October 2011 at Birnin Kebbi station.

Original Data	Surrogate data
0.1165	0.3921 ± 0.0420
0.0931	0.2029 ± 0.0756
0.1041	0.3860 ± 0.0741
0.0498	0.2891 ± 0.0598
0.1420.	0.3621 ± 0.0504

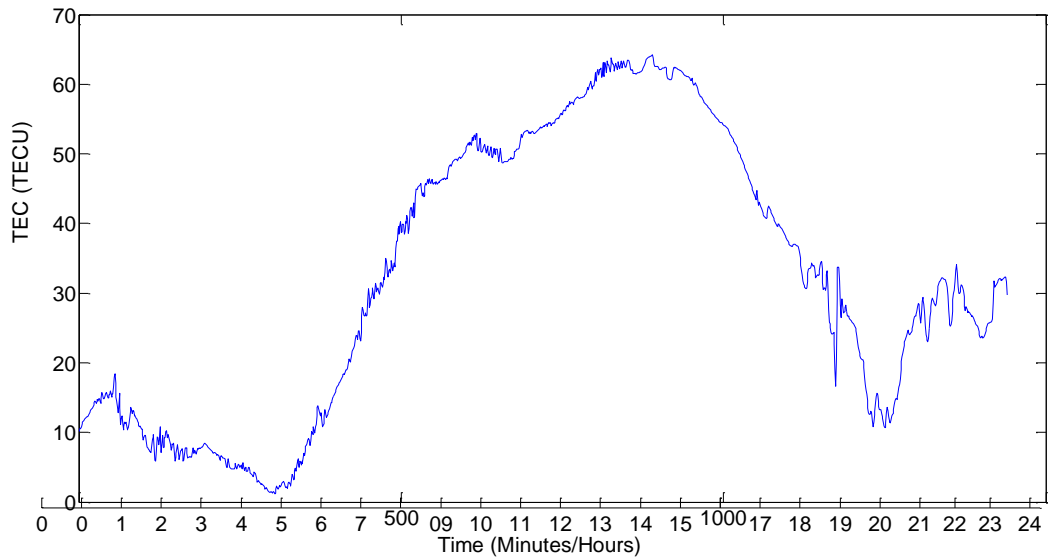
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931 Table 2b: Results of Surrogate data test for Lyapunov exponent for TEC data for the most
 932 disturbed days of October 2011 at Birnin Kebbi station.

Original Data	Surrogate data
0.0579	0.3039 ± 0.0541
0.0502	0.3156 ± 0.0428
0.0786	0.2527 ± 0.0296
0.1795	0.3662 ± 0.0468
0.1038	0.3100 ± 0.0416

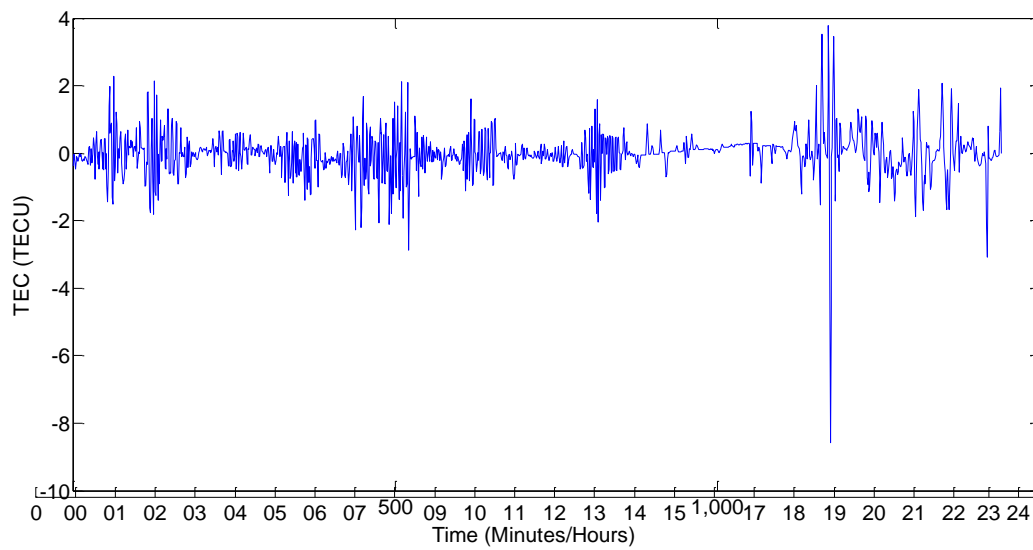
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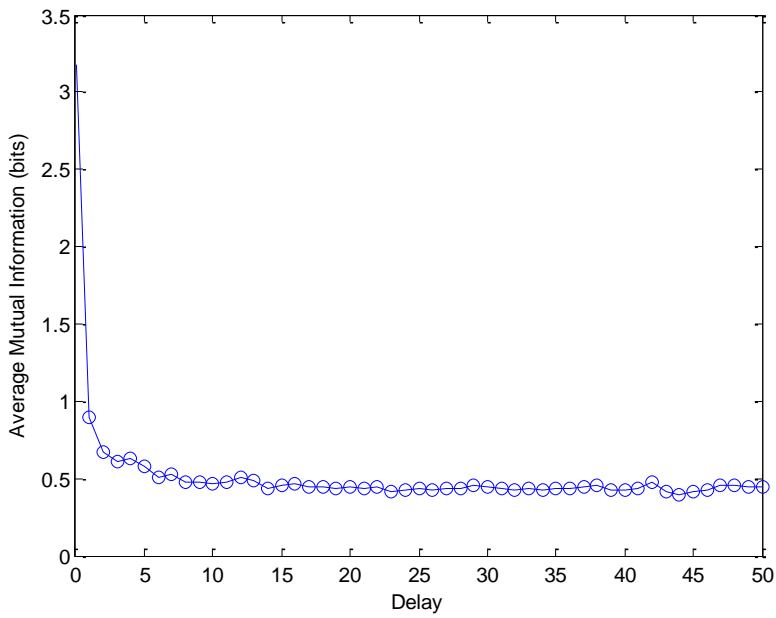
940 Fig 1. A typical time series plot for TEC measured at Lagos for 20 November 2011



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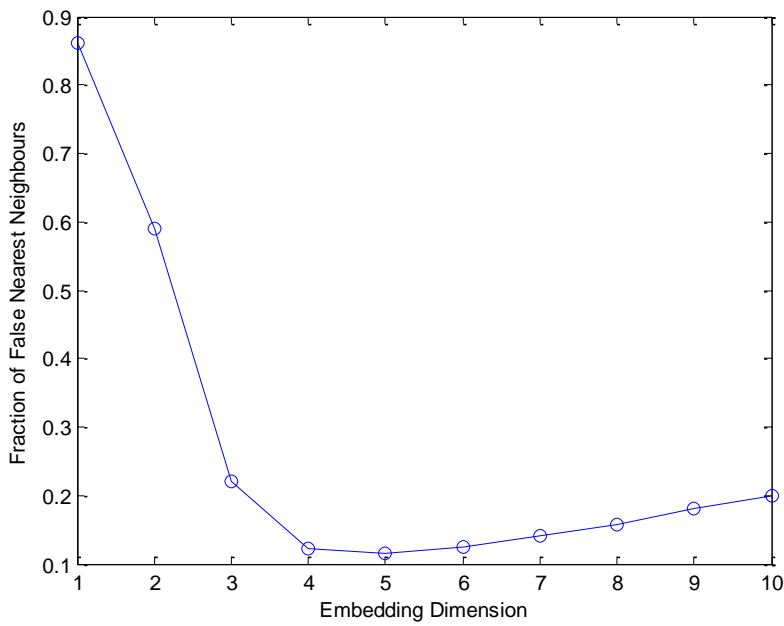
942 Fig 2. The detrended time series plot for TEC measured at Lagos

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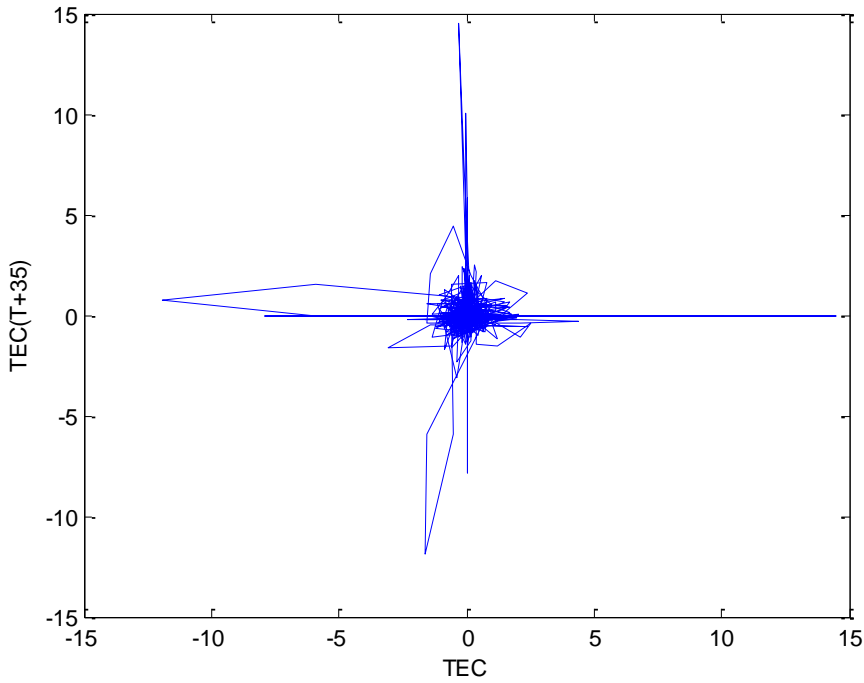
945 Fig. 3 Average mutual information against time Delay for TEC measured at Yola



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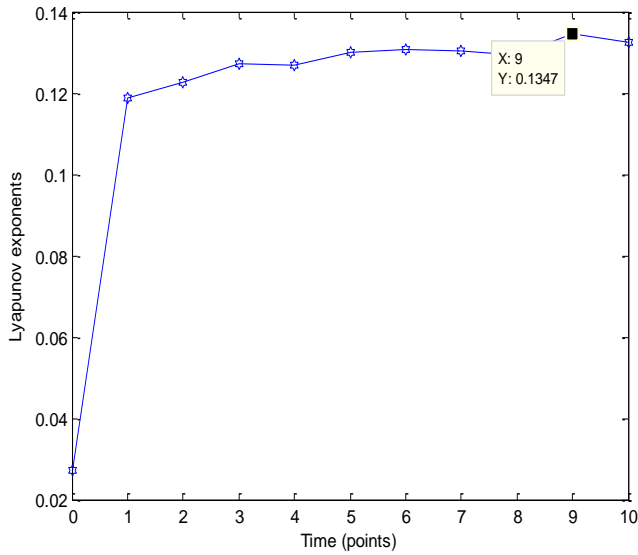
947 Fig. 4 Fraction of false nearest neighbours against embedding dimension for TEC measured at
 948 yola

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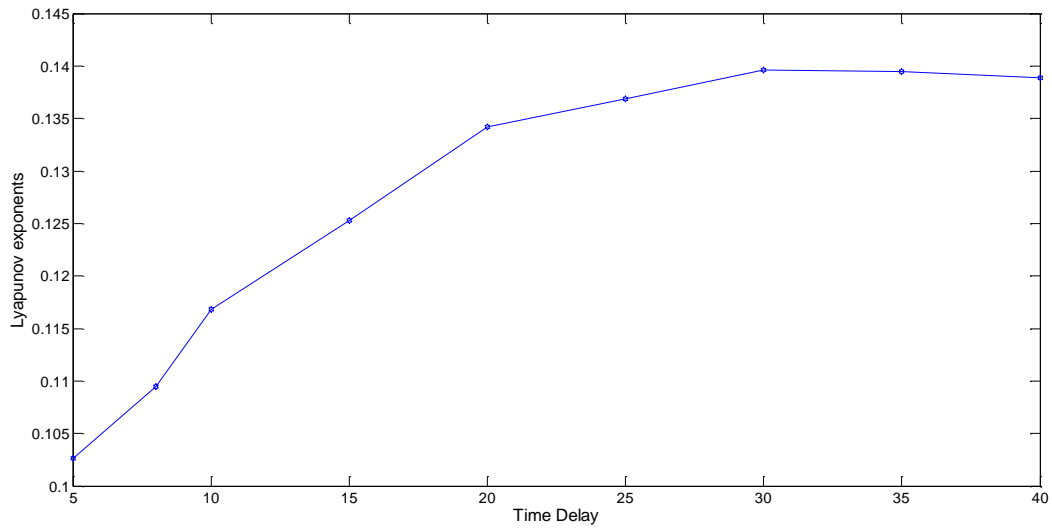
951 Fig.5 The Delay representation of the phase space reconstruction of the detrended TEC



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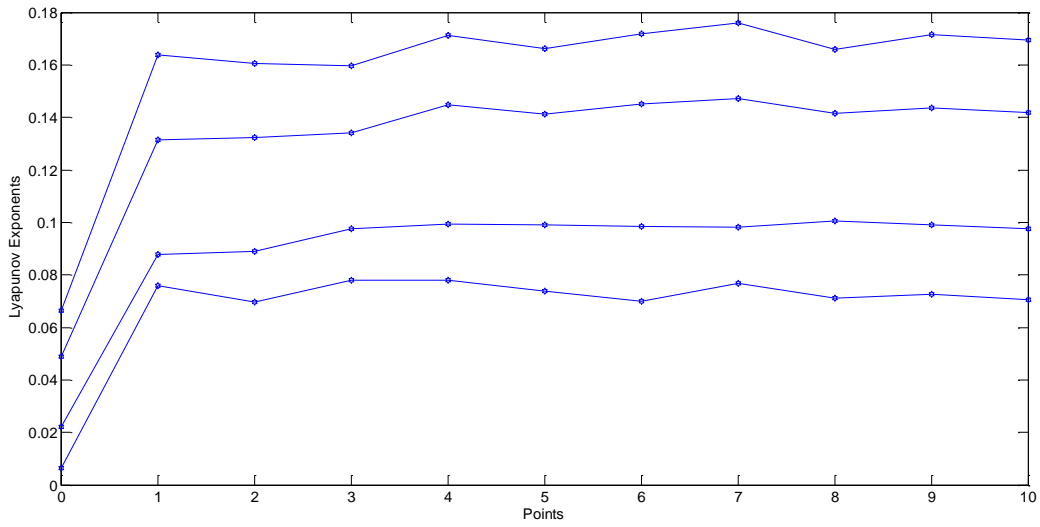
953 Fig. 6 Lyapunov Exponent computed and its evolution, computed as the state space trajectory
 954 scanned with $\tau=30$, $m=5$ for detrended time series measured at Yola with Largest Lyapunov
 955 Exponent equal to 0.1347.

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958 Fig. 6b Lyapunov exponent computed for different time delay with a constant embedding
959 dimension.



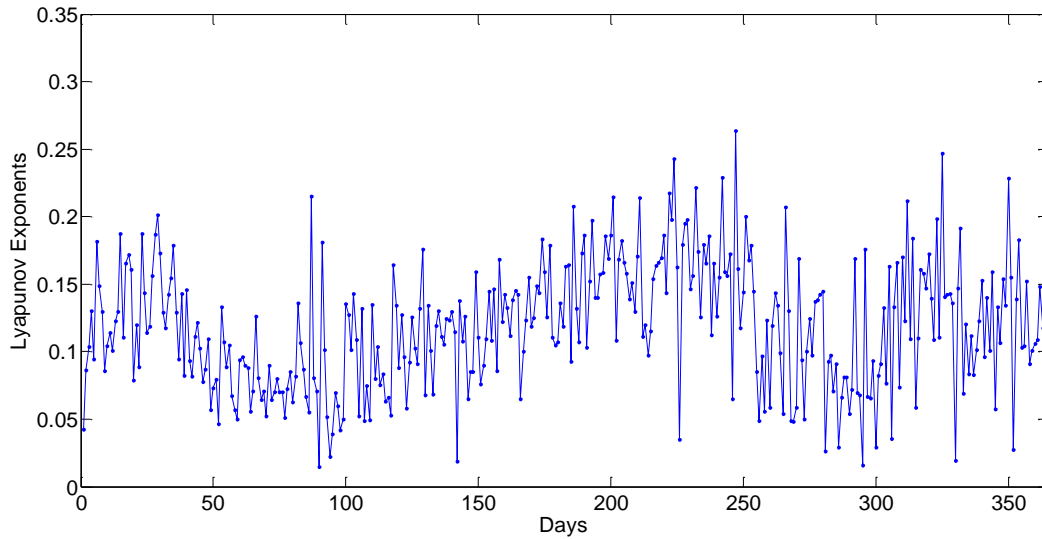
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961 Fig. 6c Lyapunov exponents computed for different embedding dimension at constant time delay

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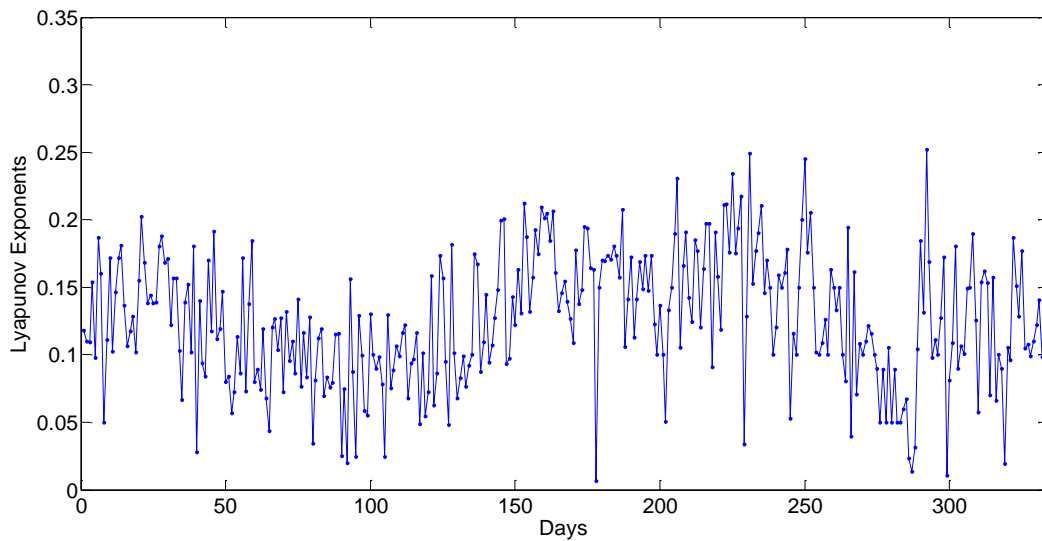
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966 Fig. 7a The transient variations of Lyapunov exponents for 365 days of 2011 for detrended TEC
967 measured at Enugu

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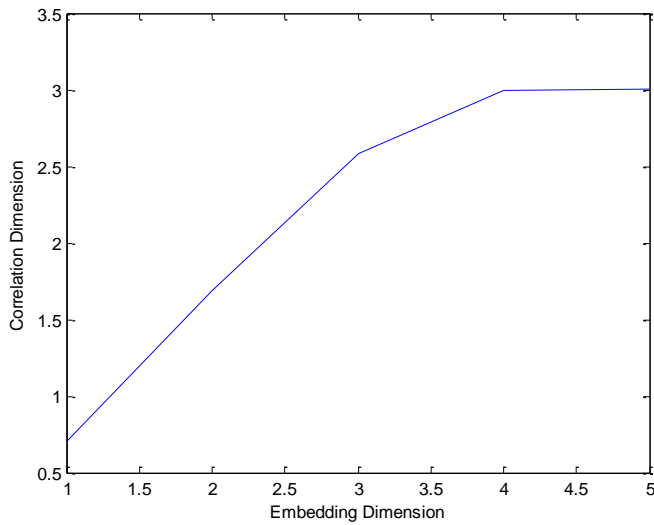


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970 Fig. 7b The transient variations of Lyapunov exponents for 334 days (Jan 1 –Nov30) of 2011 for
971 detrended TEC measured at Toro

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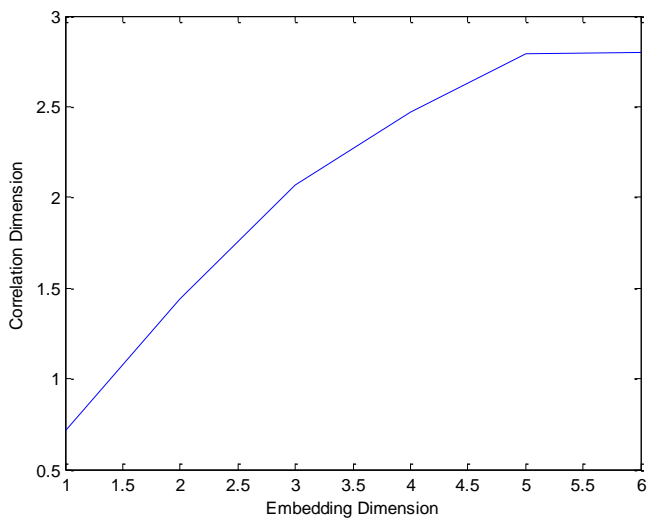
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975 Fig. 8a The correlation dimension of the detrended TEC for the quietest day of October at Birnin
976 Kebbi which saturates at $m \geq 4$ and $\tau = 39$

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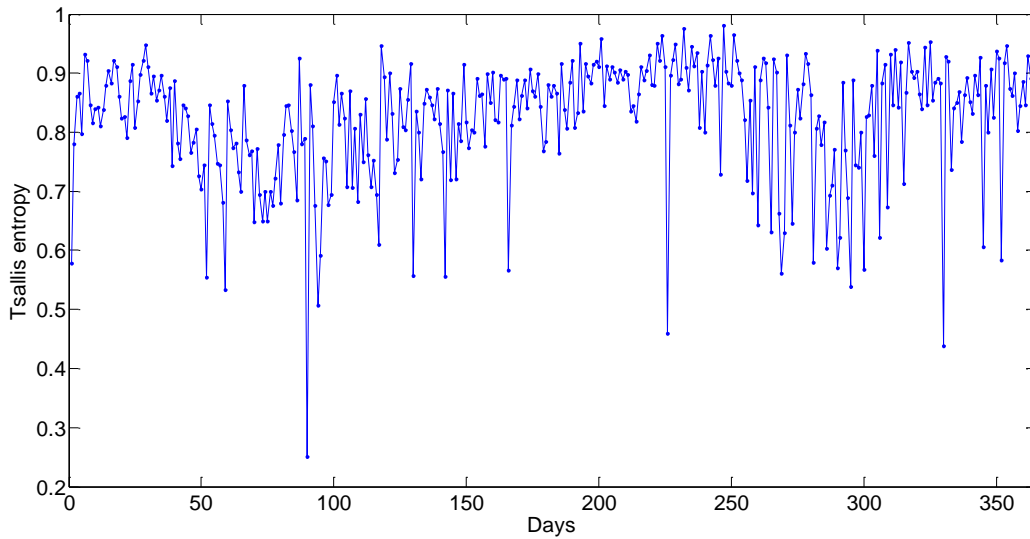


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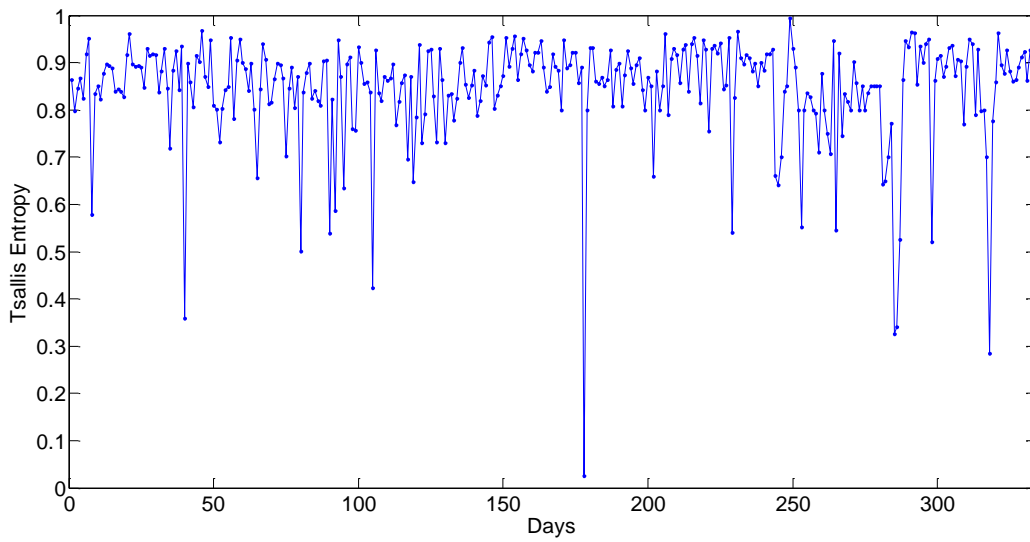
980 Fig. 8b The correlation dimension of the detrended for the most disturbed day of October at
981 Birnin Kebbi which saturates at $m \geq 5$ and $\tau = 34$

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984 Fig. 9a The transient variations of Tsallis Entropy for 365 days (Jan1 –Nov30) of 2011 for
985 detrended TEC measured at Enugu

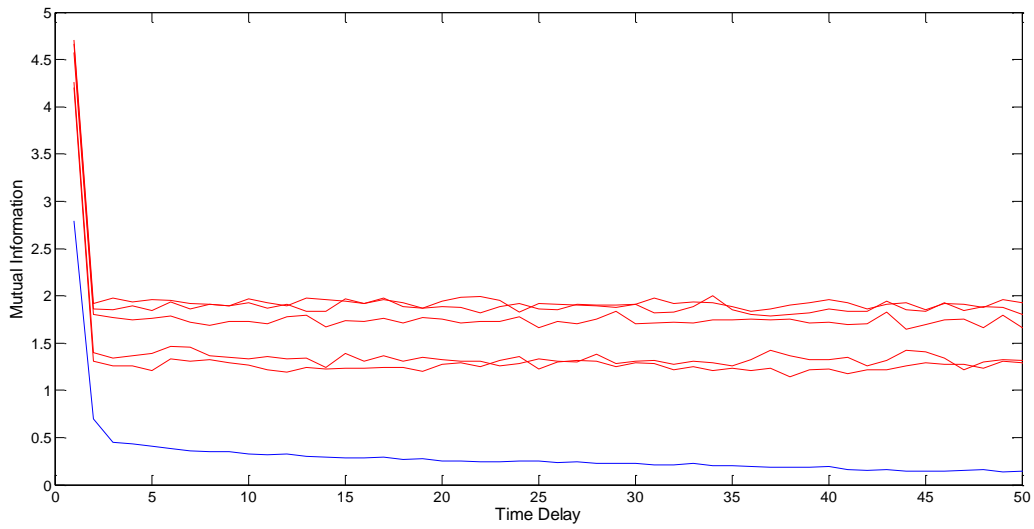


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987 Fig. 9b The transient variations of Tsallis Entropy for 334 days (Jan1 –Nov30) of 2011 for
988 detrended TEC measured at Toro

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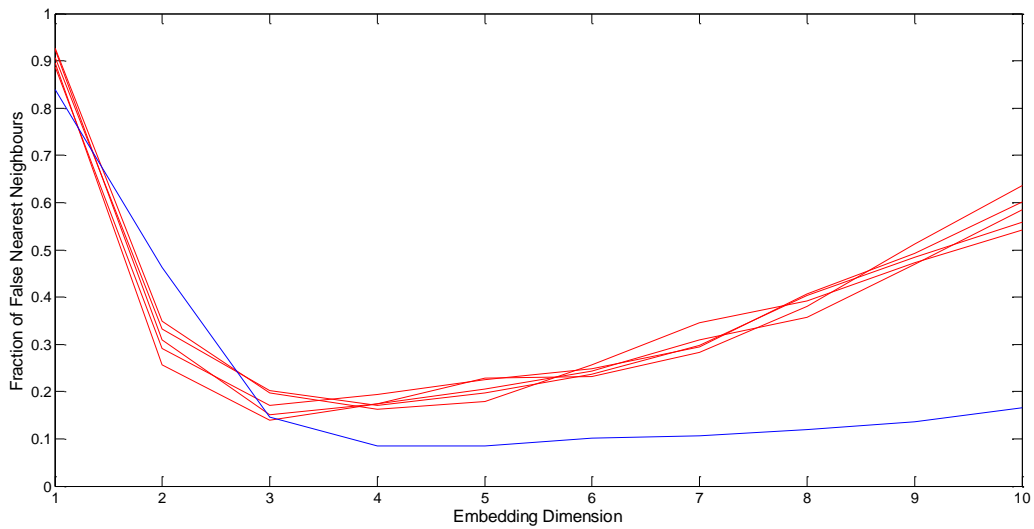
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992 Fig 10 Mutual information plotted against time delay for the original detrended data in (blue
 993 curve) with the mutual information for the surrogate data (red curve) for TEC data measured at
 994 Lagos for the quietest day of march 2011

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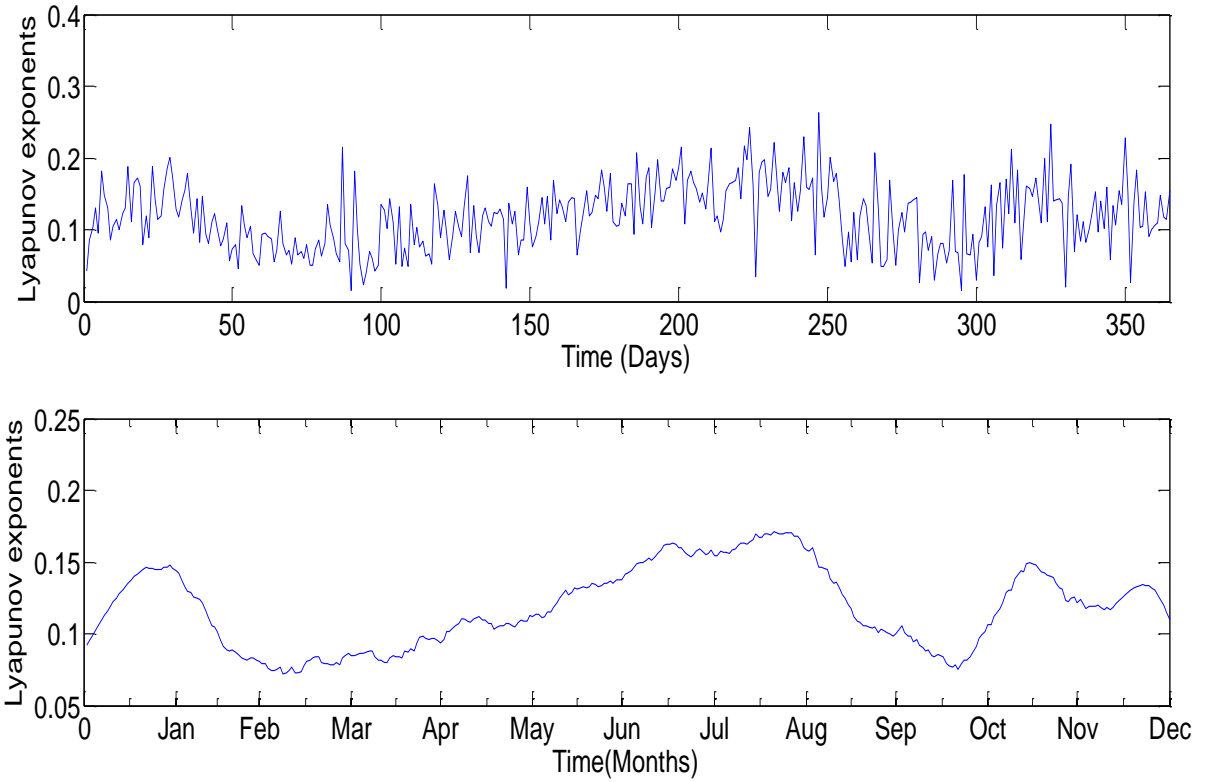


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997 Fig 11 Fraction of false nearest neighbours plotted against time embedding dimension for the
 998 original detrended data in (blue curve) with the mutual information for the surrogate data (red
 999 curve) for TEC data measured at Lagos for the quietest day of march 20

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1003 Fig. 12a Daily variation of Lyapunov exponents for TEC measured at the Enugu station for the
1004 year 2011 showing the Original data (Upper Panel) and the smoothed Plot of daily variation of
1005 Lyapunov exponents for TEC measured at the Enugu station for the year 2011 (Lower panel)

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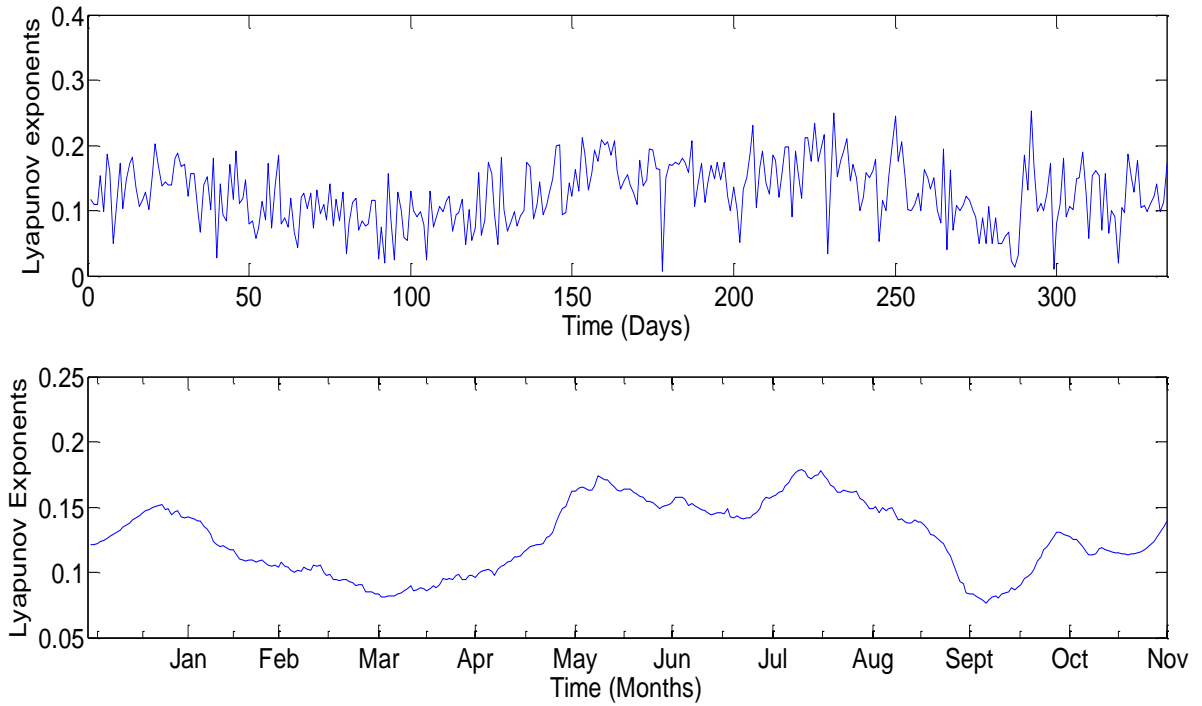
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1015 Fig 12b Daily variation of Lyapunov exponents for TEC measured at the Toro station for the
1016 year 2011 showing the Original data (Upper Panel) and the smoothed Plot of daily variation of
1017 Lyapunov exponents for TEC measured at the Toro station for the year 2011 (Lower panel)

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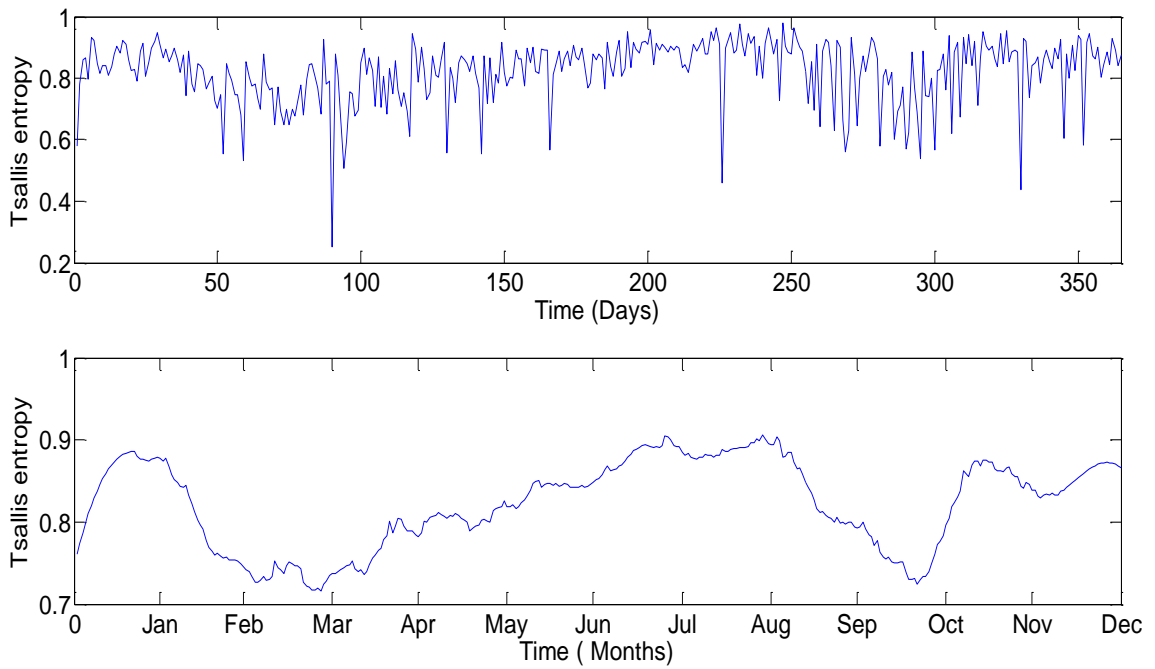
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1027 Fig. 13a Daily variation of Tsallis entropy for TEC measured at the Enugu station for the year
1028 2011 showing the Original data (Upper Panel) and the smoothed Plot of daily variation of
1029 Lyapunov exponents for TEC measured at the Enugu station for the year 2011 (Lower panel)

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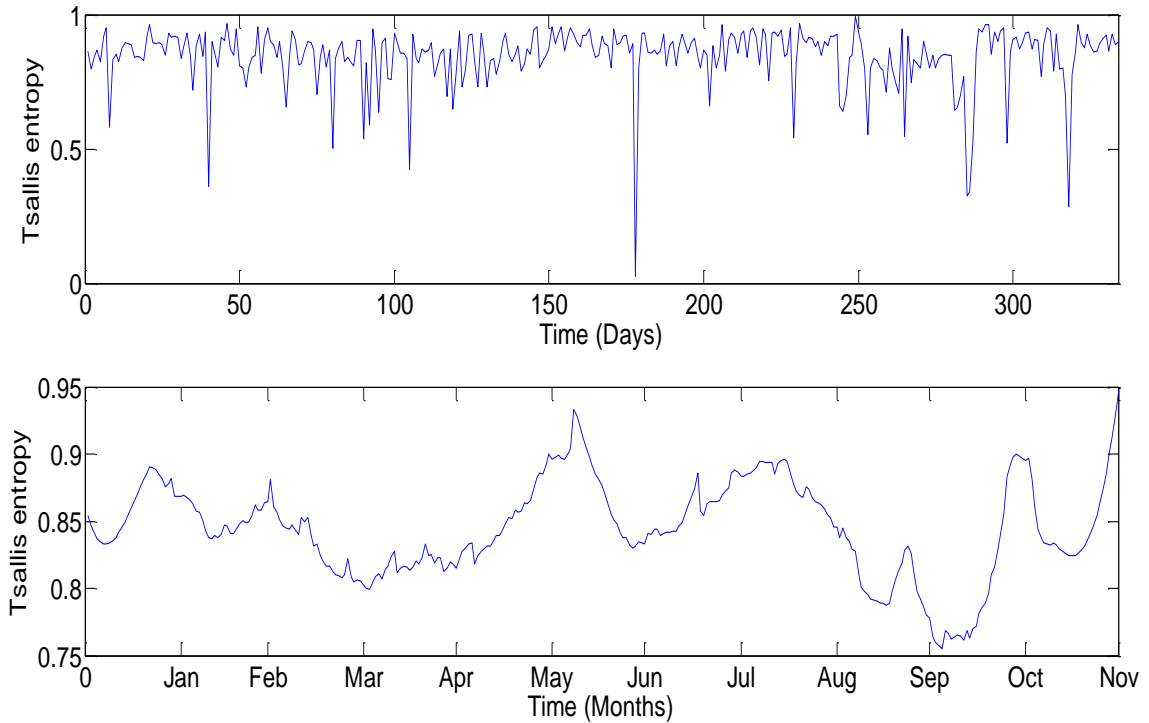
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1041 Fig. 13b Daily variation of Tsallis entropy for TEC measured at the Toro station for the year
1042 2011 showing the Original data (Upper Panel) and the smoothed Plot of daily variation of
1043 Lyapunov exponents for TEC measured at the Enugu station for the year 2011 (Lower panel)

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