



**Multifractal analysis
of mercury inclusions
in quartz**

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Multifractal analysis of mercury inclusions in quartz by X-ray computed tomography

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Abstract

In order to refine our understanding how fluid inclusions were trapped in the host minerals, we non-destructively observed mercury inclusions (liquid Hg^0) in quartz samples using X-ray computed tomography (CT) technique. The X-ray CT apparatus can observe internal structures of the samples and give cross-sectional images from the transmission of the X-rays through the samples. From the cross-sectional images, we obtained three-dimensional spatial distributions of mercury inclusions, and quantitatively analyzed them using fractal and multifractal methods. Although the samples were from different mines, the resultant fractal dimensions were about 1.7 for the samples. The fractal dimensions were also close to those predicted by diffusion-limited aggregation models and percolation theory, which are controlled by the irreversible kinetics. Then, the mercury-bearing fluids were not primary fluid inclusions, but migrated into the pre-existing cracks of quartz crystals by diffusion processes.

1 Introduction

Fluid inclusions in minerals record important clues to past geologic processes that the host minerals were subjected to. We can obtain information on physical and chemical factors, such as the pressure, temperature, density and composition of the fluids, from the fluid inclusions in minerals (e.g., Takeuchi, 1975; Roedder, 1984). However, studying fluid inclusions poses many difficulties for precise analysis because of the high mobility and evaporation involved. Therefore, non-destructive analytical methods, such as micro-Raman spectroscopy (Burke, 2001; Frezzotti et al., 2012), synchrotron-radiation X-ray fluorescence (SR-XRF) (Schmidt and Rickers, 2003; Tsuchiyama et al., 2009) and proton-induced X-ray emission (PIXE) (Kurosawa et al., 2003), have been developed for the chemical analysis of individual fluid inclusions in minerals. It is also necessary to observe the spatial distribution of different fluid inclusion populations within a mineral grain to establish their paragenetic succession relative to the mineral

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idiomorphic and polycrystalline quartz, and contain visible clusters and/or films of 1–2 mm mercury inclusions.

We used a microfocuss X-ray CT system (SMX-225CT; Shimadzu Corp., Kyoto, Japan), which can distinguish fluid inclusion of mercury (13.59 g cm^{-3}) from the quartz matrix (2.65 g cm^{-3}) based on the difference between their densities. This apparatus forms cross-sectional images by the transmission of the X-ray through the samples. The nominal resolution of the cross-section thickness and the interval between the cross-sections were 0.120 and 0.073 mm, respectively. The cross-sectional images were composed of a 512 pixel \times 512 pixel array that corresponded to an area of 9.88 cm \times 9.88 cm (therefore, one pixel corresponds to an area of 0.193 mm \times 0.193 mm). One of the cross-sectional images is shown in Fig. 1. The grey circle in Fig. 1a represents the field of view for the X-ray CT system, and the irregular region marked in the center of the field of view represents the quartz. The region outside of the quartz is air, containing no solid matter. The white dots represent mercury inclusions in the quartz.

The sequenced cross-sectional images were reconstructed and incorporated onto the images of each quartz crystal using the Image Processing and Analysis software in Java-ImageJ (Rasband, 1997–2011; Abramoff et al., 2004). Areas of mercury inclusions in the quartz were extracted from the images using threshold filtering. The spatial distributions of mercury inclusions analyzed by fractal and multifractal theory.

3 Fractal and multifractal analysis

Fractal and multifractal behavior is common in nature, and that the spatial distributions of mercury inclusions have fractal and multifractal shapes. Because a fractal shape typically has a self-similar structure and scale-free properties, the degree of distribution of the shape follows a power law in the form

$$N(r) \propto r^{-D}, \quad (1)$$

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where $N(r)$ is the number of objects, r is the scale and D is the capacity (fractal) dimension. The capacity dimension, D , which is generally estimated using a box-counting technique, is defined by the relationship between the scaling properties of the distribution by covering it with boxes of size, r , and counting the number of boxes containing $N(r)$, as follows:

$$D = - \lim_{r \rightarrow 0} \frac{\log N(r)}{\log r}. \quad (2)$$

Then, the capacity dimension can be approximately determined as the negative slope of $\log N(r)$ vs. $\log r$. Although the capacity dimension is a fundamental and quantitative parameter of the fractal, the dimension cannot completely describe complex and heterogeneous structures. Therefore, we applied the multifractal theory to analyze the spatial distribution of the mercury inclusions, as described below.

The multifractal theory can be characterized on the basis of the generalized dimensions of the q th order moment of a distribution, D_q . The generalized dimensions, D_q , can be defined by the function

$$D_q = \lim_{r \rightarrow 0} \frac{\log \sum_i P_i(r)^q}{(q-1) \log r}, \quad (3)$$

where $P_i(r)^q$ is the probability within the i th region of a measured quantity varying with scale r (Takayasu, 1986). The generalized dimension, D_q , can be rewritten with the singularity exponent, α , and the generalized fractal dimension, $f(\alpha)$, as the equation

$$D_q = \frac{q\alpha(q) - f(\alpha(q))}{(q-1)}. \quad (4)$$

When we can obtain the generalized dimension, D_q , from experiments, the singularity exponent, α , and the generalized fractal dimension, $f(\alpha)$, can be estimated using the following relationships

$$\alpha(q) = \frac{d}{dq} (q-1)D_q \quad (5)$$

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Several studies have been performed using fractal geometry, which is controlled by the irreversible kinetic processes such as diffusion, aggregation and percolation. These processes have been described by the diffusion-limited aggregation (DLA) and percolation mechanism, which involves Brownian particle motion (e.g., Witten and Sander, 1983; Meakin, 1985; Zheng et al., 1998; Stauffer, 1979; Hunt and Ewing, 2009). Based on the DLA models, the fractal dimension, D , is predicted to be equal to $(d^2 + 1)/(d + 1)$, where d is the spatial dimensionality (Muthukumar, 1983). For a two- and three-dimensional system D should be $5/3$ (1.66) and $5/2$ (2.5), respectively. Also, percolation theory shows that the fractal dimensions are 1.89 and 2.54 in two- and three-dimensional systems D , respectively. Estimated fractal dimensions of the mercury inclusions in our quartz samples from San Benito and Itomuka are similar to these obtained by DLA models and percolation theory for two-dimensional system. This suggests that mercury-bearing fluids would imply information on paragenetic succession of mineral and crystallization processes.

Dunning et al. (2005) observed deposit of San Benito sample in geology and mineralogy, and deduced that quartz is formed in earlier than mercury is done. Itomuka sample is also the same formation and age as San Benito sample, although they are from different deposits. From the fractal dimensions and geological deduction, fluids of mercury inclusions would be captured into quartz after its crystallization process. Therefore, the mercury-bearing fluids were not the primary fluid inclusions, but were trapped in cracks that already existed in the quartz samples.

Multifractal analysis allowed us to examine the complex signature of the mercury inclusion clusters more quantitatively. We obtained three sets of multifractal parameters, α , $f(\alpha)$ and D_q . Figure 3 shows multifractal spectra in which the generalized fractal dimensions, $f(\alpha)$, are plotted against the singularity exponent, α . Multifractal spectra generally show parabolic curves that are concave downwards, and the curve maxima occur at $q = 0$, at which point $f(\alpha)$ in the San Benito and Itomuka samples corresponded to the capacity dimensions of 1.70 and 1.71, respectively. The parabola of both samples had similar curves, but the width of the parabola was larger

for the San Benito sample than for the Itomuka sample. The San Benito sample had more fractal structure patterns than the Itomuka sample, because wider curves reflect more heterogeneous structures in the mercury inclusion distributions.

Figure 4 shows the difference between the fractal dimension distributions in the samples, and the relationships between the generalized dimensions, D_q , and the order moments, q . In general, D_q dimensions increase with decreasing q moments. In Fig. 4, the curves for both quartz samples are similar. The q moments are interpreted as a parameter of the inclusion distribution probability densities, i.e., low- q implies a low distribution density and high- q implies a high density. The D_q dimensions at low- q moments in the samples were 2.3–2.7, indicating that the distribution of surface and spatial structures were scarce. However, the D_q dimensions at high- q moments were 1.0–1.3, indicating that linear configurations were common. Consequently, mercury inclusions ramified like a dendritical structure. This implies that mercury-bearing fluids migrated linearly into the quartz sample cracks and then expanded to form surface plane and spatial figurations.

5 Conclusions

We proposed to analyze of mercury inclusions in quartz samples using in situ X-ray computed tomography (CT). The X-ray CT apparatus allows the internal structures of the quartz samples to be analyzed non-destructively, and therefore, mercury inclusions are to be retained in the quartz throughout the experiment. We performed X-ray CT measurements on two quartz samples, one from San Benito, California, USA and one from Itomuka, Hokkaido, Japan, both of which contain visible mercury inclusions. We obtained the spatial distributions of the mercury inclusions based on sequenced X-ray cross-sectional images. Spatial distributions can be explored quantitatively using fractal and multifractal theory. We showed the fractal dimensions, $\alpha-f(\alpha)$, multifractal spectra and the relationship between q and D_q . The mercury inclusions trapped in the quartz samples had similar distribution signatures, even though the quartz samples were from

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different mines. In addition, the fractal dimensions were close to those obtained by diffusion-limited aggregation DLA models and percolation theory for two-dimensional system. The similar mercury inclusion signatures correspond to the samples being formed during a process of diffusion into pre-existing cracks in the quartz. After the formation of crystalline quartz, the mercury-bearing fluids probably migrated into pre-existing cracks in the quartz crystals, and native mercury was lodged in the cavities.

Author contribution. T. Shibata designed the study, T. Maruoka performed the X-ray CT measurements and T. Echigo prepared samples. T. Maruoka. and T. Echigo gave technical supports and conceptual advice, and T. Shibata prepared the manuscript with contributions from all co-authors.

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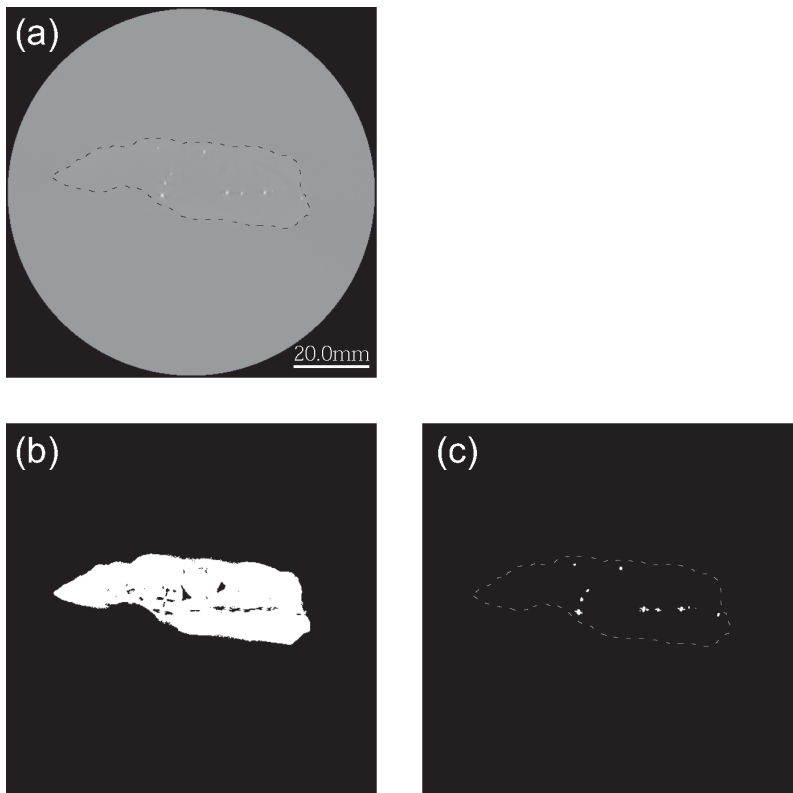


Figure 1. Two-dimensional (2-D) slice image of the San Benito quartz sample obtained with a microfocus X-ray CT system. **(a)** Original grayscale image. The circular dark gray region is the measurable area, the irregular central region is the quartz sample, and the white points are mercury inclusions; **(b)** binarized image of the quartz crystal; and **(c)** binarized image of the mercury inclusions. The quartz area is edged with a dotted line.

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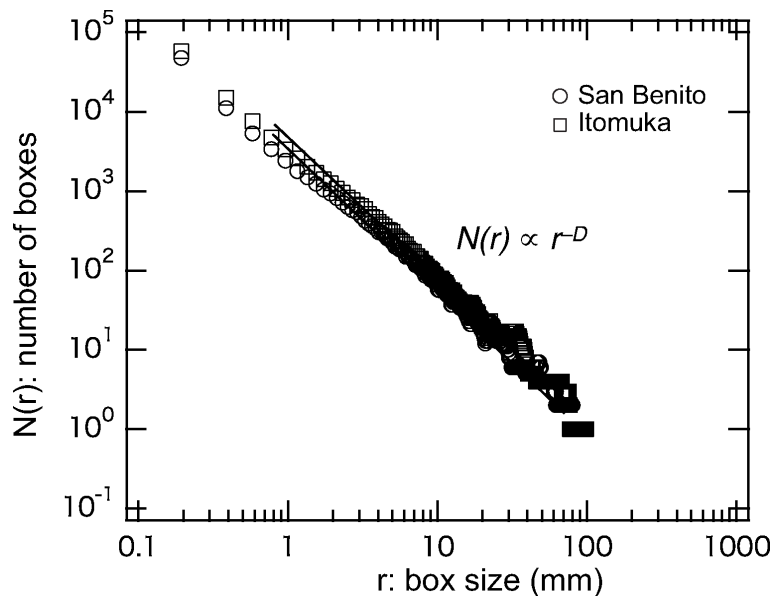


Figure 2. Box counting plot for mercury inclusions in the San Benito and Itomuka quartz samples. To obtain the capacity dimensions, D , only the middle portions of the sequence, from four times the length of the minimal box size to 9/10 of the length of the maximal box size, were analysed to avoid edge effects.

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