
Large-scale linear instability of the generalized Kolmogorov flow - viscous version with a general parallel flow

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U of
any shape,
periodic within the interval $[0, 2\pi]$
with the condition that $\int_0^{2\pi} U[y] dy = 0$
Streamfunction defined as $u = \partial_y \psi$, $v = -\partial_x \psi$

Settings

- Last use

```
DateList[]

{2009, 1, 14, 16, 13, 15.805210}
```

■ General

```
Off[General::spell1]
```

■ Laplacian

```

$$\Delta = - (\epsilon^2 \partial_{x,x} \# + \partial_{y,y} \#) \&$$


$$- (\epsilon^2 \partial_{x,x} \# 1 + \partial_{y,y} \# 1) \&$$

```

Beware of the sign in the definition

■ Definition of the dissipation

Standard viscosity

```
DS = - ν Δ[Δ[#]] &
- ν Δ[Δ[#1]] &
```

Definition of the modified Kolmogorov Flow

■ Velocity and streamfunction

```
Φ[y_] := Integrate[U[s], {s, π/2, y}]
```

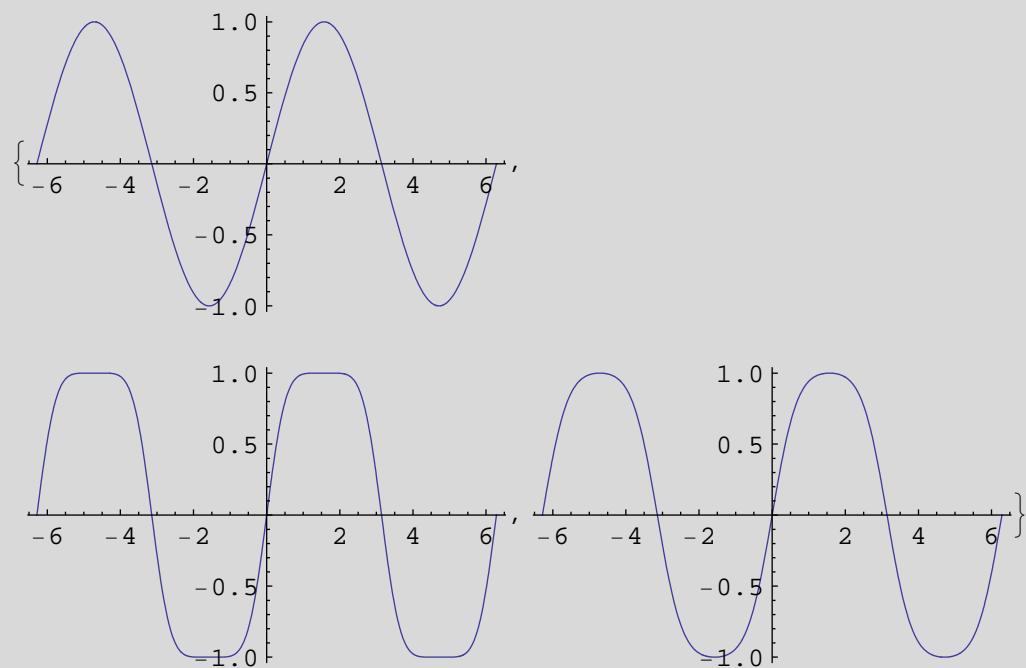
■ Examples of flows

Standard Kolmogorov flow in Flow1

```
Flow2 = U → Evaluate[Sin[# + 1/2 Sin[2 #]] &];
Flow1 = U → Evaluate[Sin[#] &];
Flow3 = U → Evaluate[Sin[# + 1/4 Sin[2 #]] &];
```

■ Plots of the flow

```
{Plot[U[y] /. {Flow1, A→1}, {y, -2 π, 2 π}],
 Plot[U[y] /. {Flow2, A→1}, {y, -2 π, 2 π}],
 Plot[U[y] /. {Flow3, A→1}, {y, -2 π, 2 π}]}  
}
```

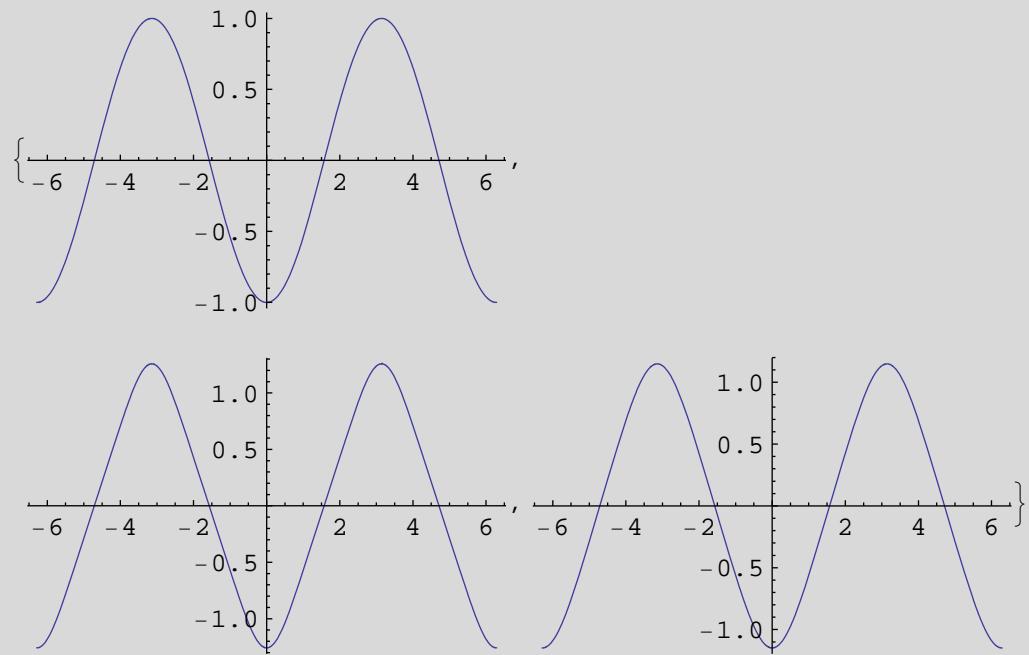


■ Plots of the streamfunction

```

 $\Phi_1 = \text{NIntegrate}\left[\text{Evaluate}[U[y] /. \{\text{Flow1}, A \rightarrow 1\}], \left\{y, \frac{\pi}{2}, \#\right\}\right] \&;$ 
 $\Phi_2 = \text{NIntegrate}\left[\text{Evaluate}[U[y] /. \{\text{Flow2}, A \rightarrow 1\}], \left\{y, \frac{\pi}{2}, \#\right\}\right] \&;$ 
 $\Phi_3 = \text{NIntegrate}\left[\text{Evaluate}[U[y] /. \{\text{Flow3}, A \rightarrow 1\}], \left\{y, \frac{\pi}{2}, \#\right\}\right] \&;$ 
{Plot[\Phi1[u], {u, -2\pi, 2\pi}],
 Plot[\Phi2[u], {u, -2\pi, 2\pi}],
 Plot[\Phi3[u], {u, -2\pi, 2\pi}]}

```



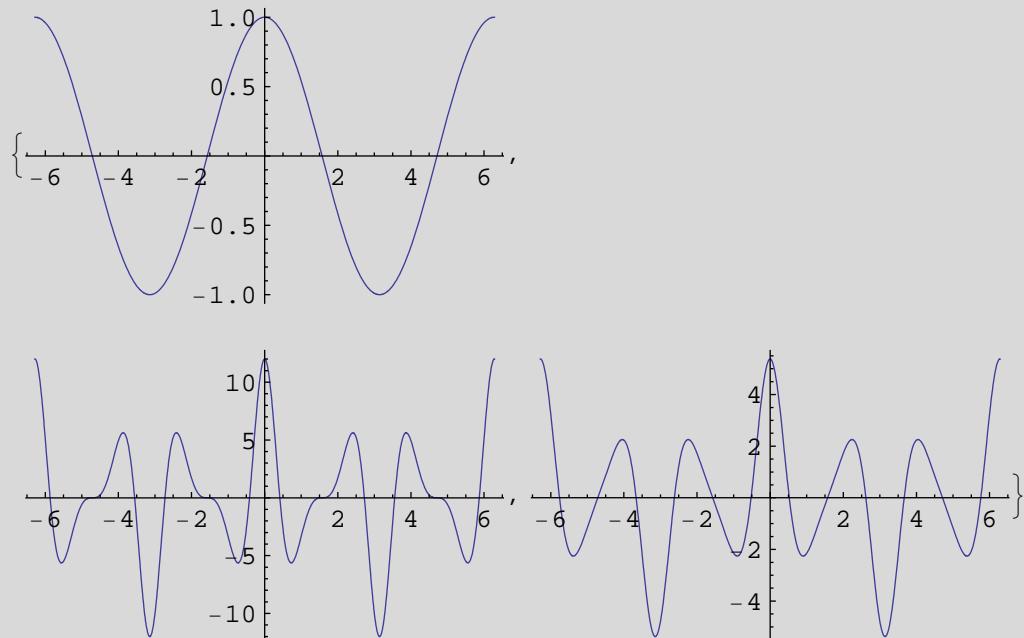
■ Plots of the forcing

```

Forcing = Simplify[DS[\Phi[y]]]
-ν U^(3) [Y]

```

```
{Plot[Forcing //.{Flow1, A→1, ν→1},
{y, -2π, 2π}, PlotRange→All],
Plot[Forcing //.{Flow2, A→1, ν→1},
{y, -2π, 2π}, PlotRange→All],
Plot[Forcing //.{Flow3, A→1, ν→1},
{y, -2π, 2π}, PlotRange→All]}
```



Definition of the perturbation problem

■ 2D Navier-Stokes with scaling

We have taken $u = \partial_y \psi$ and $v = -\partial_x \psi$ (mechanical definition)

a linear friction term has been included

$$\begin{aligned} NS = & \epsilon^2 \partial_T \Delta[\psi[T, X, Y]] - \epsilon \partial_X \psi[T, X, Y] \partial_Y \Delta[\psi[T, X, Y]] + \\ & \epsilon \partial_Y \psi[T, X, Y] \partial_X \Delta[\psi[T, X, Y]] + \epsilon^2 r \Delta[\psi[T, X, Y]] - \\ & DS[\psi[T, X, Y]] + DS[\Phi[Y]] - \epsilon^2 r \Delta[\Phi[Y]]; \end{aligned}$$

■ Expansion

$$\begin{aligned} \psi[T_-, X_-, Y_-] := & \Phi[Y] + \psi_0[T, X, Y] + \\ & \epsilon \psi_1[T, X, Y] + \epsilon^2 \psi_2[T, X, Y] + \epsilon^3 \psi_3[T, X, Y] + \epsilon^4 \psi_4[T, X, Y] \end{aligned}$$

```
NS2 = Collect[Normal[Series[NS, {ε, 0, 4}]], ε];
```

Order 0

```
a0 = Simplify[Coefficient[NS2, ε, 0]]  
ν ψ0^(0,0,4) [T, X, Y]
```

■ Symbolic solution

```
Rule0 = ψ0 → (ψ0[#1, #2] &);
```

■ Check solution

```
Simplify[a0 /. {Rule0}]  
0
```

Solvability condition

```
Solva[expr_, Rules___] :=  
Integrate[expr //. Flatten[{Rules}], {y, 0, 2 π}] /. g___[2 π] → g[0]
```

Order 1

```
a1 = Simplify[Expand[Coefficient[NS2, ε, 1]] /. Rule0]  
U''[y] ψ0^(0,1) [T, X] + ν ψ1^(0,0,4) [T, X, Y]
```

■ Define linear operator

```
ℒ[f_] := + ∂y,y (ν ∂y,y f )
```

■ Extraction and integration of second member

```
a1a = -Simplify[Coefficient[a1, φ0^(0,1)[T, X]]]
-U''[Y]
```

```
a1aI2 = Integrate[Integrate[a1a, y], y]
-U[Y]
```

■ Reconstruction and check of the first order equation

```
a1N = L[ψ1[T, X, Y]] - a1a φ0^(0,1)[T, X];
Simplify[a1N - a1]
0
```

■ Solvability condition

The solvability condition is obviously satisfied with this second member since it has a primitive (even two primitives)

```
Solva[a1a] /. g___[2 π] → g[0]
0
```

■ Symbolic solution

```
Rule1 = ψ1 → (f1[#3] φ0^(0,1)[#1, #2] &);
```

Order 2

```
a2 = Simplify[Coefficient[NS2, ε, 2] //.{Rule0, Rule1}, Trig → True]
f1^(3)[Y] φ0^(0,1)[T, X]^2 - U[Y] f1''[Y] φ0^(0,2)[T, X] +
f1[Y] U''[Y] φ0^(0,2)[T, X] + ν ψ2^(0,0,4)[T, X, Y]
```

Extraction of second member and solvability

- term in factor of $\varphi_0^{(0,1)} [T, x]^2$

```
a2a = -Factor[Simplify[Coefficient[a2, \varphi_0^{(0,1)} [T, x], 2]]]
Solve[a2a]

-f1^{(3)} [Y]
```

0

- term in factor of $\varphi_0^{(0,2)} [T, x]$

```
a2b = -Simplify[Coefficient[a2, \varphi_0^{(0,2)} [T, x], 1]]
Integrate[a2b, y]
Solve[a2b]

U[Y] f1''[Y] - f1[Y] U''[Y]
```

U[Y] f1'[Y] - f1[Y] U'[Y]

0

■ Reconstruction and check of the second order equation

```
a2N = L[\psi2[T, X, Y]] - a2b \varphi_0^{(0,2)} [T, X] - a2a \varphi_0^{(0,1)} [T, X]^2;
Simplify[a2N - a2]
```

0

■ Symbolic solution

```
Rule2 = \psi2 \rightarrow \left( f2a[\#3] \varphi_0^{(0,1)} [\#1, \#2]^2 + f2b[\#3] \varphi_0^{(0,2)} [\#1, \#2] \& \right);
```

Order 3

```
a3 = FullSimplify[
  Coefficient[NS2, ε, 3] //.{Rule0, Rule1, Rule2}, Trig → True]

f2a(3)[y] φ0(0,1)[T, X]3 +
(-2 U[y] f2a''[y] + 2 f2a[y] U''[y] + f1[y] f1(3)[y] + f2b(3)[y]) φ0(0,1)[T, X]
φ0(0,2)[T, X] - (U[y] (1 + f2b''[y]) - f2b[y] U''[y]) φ0(0,3)[T, X] - f1''[y]
(φ0(0,1)[T, X] (r + f1'[y] φ0(0,2)[T, X]) - 2 √ φ0(0,3)[T, X] + φ0(1,1)[T, X]) +
√ ψ3(0,0,4)[T, X, Y]
```

■ Extraction of second member and solvability

```
a31 = Collect[Expand[a3],
{φ0(1,1)[T, X], φ0(0,1)[T, X], φ0(0,2)[T, X], φ0(0,3)[T, X]}]

f2a(3)[y] φ0(0,1)[T, X]3 +
φ0(0,1)[T, X] (-r f1''[y] + (-f1'[y] f1''[y] - 2 U[y] f2a''[y] +
2 f2a[y] U''[y] + f1[y] f1(3)[y] + f2b(3)[y]) φ0(0,2)[T, X]) +
(-U[y] + 2 √ f1''[y] - U[y] f2b''[y] + f2b[y] U''[y]) φ0(0,3)[T, X] -
f1''[y] φ0(1,1)[T, X] + √ ψ3(0,0,4)[T, X, Y]
```

■ Term in $\phi_0^{(0,1)}[T, X]^3$: a3a

```
a3a = -Coefficient[a31, φ0(0,1)[T, X]3]
Solve[a3a]

-f2a(3)[y]
```

0

■ Term in $\varphi_0^{(0,1)}[T, X] \varphi_0^{(0,2)}[T, X]$: a3b

```
a3b =
-Coefficient[Coefficient[Expand[a31], \varphi_0^{(0,1)}[T, X], 1], \varphi_0^{(0,2)}[T, X]]
Integrate[a3b, y]
Solva[a3b]
```

$$f1'[y] f1''[y] + 2 U[y] f2a''[y] - 2 f2a[y] U''[y] - f1[y] f1^{(3)}[y] - f2b^{(3)}[y]$$

$$f1'[y]^2 + 2 U[y] f2a'[y] - 2 f2a[y] U'[y] - f1[y] f1''[y] - f2b''[y]$$

0

■ Term in $\varphi_0^{(0,3)}[T, X]$: a3d

```
a3d = -Coefficient[a31, \varphi_0^{(0,3)}[T, X]]
Integrate[a3d - U[y], y]
Solva[a3d - U[y]]
```

$$U[y] - 2 v f1''[y] + U[y] f2b''[y] - f2b[y] U''[y]$$

$$-2 v f1'[y] + U[y] f2b'[y] - f2b[y] U'[y]$$

0

The solvability condition is satisfied due to the hypothesis that $\int_0^{2\pi} U[y] dy = 0$,
no other contribution

■ Term in $\varphi_0^{(1,1)}[T, X]$: a3g

```
a3g = -Coefficient[a31, \varphi_0^{(1,1)}[T, X]]
Solva[a3g]
```

$$f1''[y]$$

0

Terms due to friction

```
a3r = -Coefficient[a31, r φ0^(0,1) [T, x]]
Solve[a3r]
```

```
f1''[y]
```

```
0
```

■ Reconstruction and check of the third order equation

```
a3N =
-a3g φ0^(1,1) [T, x] - a3d φ0^(0,3) [T, x] - φ0^(0,1) [T, x] (a3b φ0^(0,2) [T, x]) -
a3a φ0^(0,1) [T, x]^3 - r a3r φ0^(0,1) [T, x] + L[ψ3[T, x, y]];
Simplify[a3N - a3]
```

```
0
```

■ Symbolic solution

```
Rule3 = ψ3 -> (f3g[#3] φ0^(1,1) [#1, #2] + f3d[#3] φ0^(0,3) [#1, #2] +
φ0^(0,1) [#1, #2] (f3b[#3] φ0^(0,2) [#1, #2]) +
f3a[#3] φ0^(0,1) [#1, #2]^3 + r f3r[#3] φ0^(0,1) [#1, #2] &);
```

Order 4

```

a4 = Simplify[Coefficient[NS2, ε, 4] //.
{Rule0, Rule1, Rule2, Rule3}, Trig → True]

-x φ0^(0,2)[T, X] -
f1''[Y] φ0^(0,2)[T, X] (f2a'[Y] φ0^(0,1)[T, X]^2 + f2b'[Y] φ0^(0,2)[T, X]) -
r (f2a''[Y] φ0^(0,1)[T, X]^2 + f2b''[Y] φ0^(0,2)[T, X]) +
f1[Y] φ0^(0,2)[T, X] (f2a^(3)[Y] φ0^(0,1)[T, X]^2 + f2b^(3)[Y] φ0^(0,2)[T, X]) +
f1'[Y] φ0^(0,1)[T, X] φ0^(0,3)[T, X] + f1^(3)[Y] φ0^(0,1)[T, X]
(2 f2a[Y] φ0^(0,1)[T, X] φ0^(0,2)[T, X] + f2b[Y] φ0^(0,3)[T, X]) - f1'[Y] φ0^(0,1)[
T, X] (2 f2a''[Y] φ0^(0,1)[T, X] φ0^(0,2)[T, X] + (1 + f2b''[Y]) φ0^(0,3)[T, X]) +
ν φ0^(0,4)[T, X] + 2 ν (2 f2a''[Y] (φ0^(0,2)[T, X]^2 + φ0^(0,1)[T, X] φ0^(0,3)[T, X]) +
f2b''[Y] φ0^(0,4)[T, X]) - 2 f2a''[Y] φ0^(0,1)[T, X] φ0^(1,1)[T, X] +
φ0^(0,1)[T, X] (r f3r^(3)[Y] φ0^(0,1)[T, X] + f3a^(3)[Y] φ0^(0,1)[T, X]^3 +
f3b^(3)[Y] φ0^(0,1)[T, X] φ0^(0,2)[T, X] + f3d^(3)[Y] φ0^(0,3)[T, X] +
f3g^(3)[Y] φ0^(1,1)[T, X]) - φ0^(1,2)[T, X] - f2b''[Y] φ0^(1,2)[T, X] +
U''[Y] (r f3r[Y] φ0^(0,2)[T, X] + 3 f3a[Y] φ0^(0,1)[T, X]^2 φ0^(0,2)[T, X] +
f3b[Y] φ0^(0,2)[T, X]^2 + f3b[Y] φ0^(0,1)[T, X] φ0^(0,3)[T, X] +
f3d[Y] φ0^(0,4)[T, X] + f3g[Y] φ0^(1,2)[T, X]) -
U[Y] (r f3r''[Y] φ0^(0,2)[T, X] + 3 f3a''[Y] φ0^(0,1)[T, X]^2 φ0^(0,2)[T, X] + f3b''[Y]
φ0^(0,2)[T, X]^2 + f3b''[Y] φ0^(0,1)[T, X] φ0^(0,3)[T, X] + f1[Y] φ0^(0,4)[T, X] +
f3d''[Y] φ0^(0,4)[T, X] + f3g''[Y] φ0^(1,2)[T, X]) + ν ψ4^(0,0,4)[T, X, Y]

```

■ Terms containing $\varphi_0^{(0,1)} [T, X]$

■ Terms in $\varphi_0^{(0,1)} [T, X]^4$: a4a

```
a4a = -Coefficient[Expand[a4], \varphi_0^{(0,1)} [T, X]^4]
Solve[a4a]

-f3a^{(3)} [y]
```

0

■ Terms in $\varphi_0^{(0,1)} [T, X]^3$

```
Coefficient[Expand[a4], \varphi_0^{(0,1)} [T, X]^3]

0
```

■ Terms in $\varphi_0^{(0,1)} [T, X]^2$

```
a42 = Coefficient[Expand[a4], \varphi_0^{(0,1)} [T, X]^2]

-r f2a''[y] + r f3r^{(3)} [y] - f2a'[y] f1''[y] \varphi_0^{(0,2)} [T, X] -
2 f1'[y] f2a''[y] \varphi_0^{(0,2)} [T, X] - 3 U[y] f3a''[y] \varphi_0^{(0,2)} [T, X] +
3 f3a[y] U''[y] \varphi_0^{(0,2)} [T, X] + 2 f2a[y] f1^{(3)} [y] \varphi_0^{(0,2)} [T, X] +
f1[y] f2a^{(3)} [y] \varphi_0^{(0,2)} [T, X] + f3b^{(3)} [y] \varphi_0^{(0,2)} [T, X]
```

■ Terms in $r \varphi_0^{(0,1)} [T, X]^2$: a4u

```
a4u = -Coefficient[Expand[a42], r]
Solve[a4u]

f2a''[y] - f3r^{(3)} [y]
```

0

■ Terms in $\varphi_0^{(0,1)}[T, X]^2 \varphi_0^{(0,2)}[T, X] : a4c$

```
a4c = -Coefficient[Expand[a42], φ0(0,2)[T, X]]
Integrate[a4c, y]
Solve[a4c]
```

$$f2a'[y] f1''[y] + 2 f1'[y] f2a''[y] + 3 U[y] f3a''[y] - 3 f3a[y] U''[y] - 2 f2a[y] f1⁽³⁾[y] - f1[y] f2a⁽³⁾[y] - f3b⁽³⁾[y]$$

$$3 f1'[y] f2a'[y] + 3 U[y] f3a'[y] - 3 f3a[y] U'[y] - 2 f2a[y] f1''[y] - f1[y] f2a''[y] - f3b''[y]$$

0

■ Terms in $\varphi_0^{(0,1)}[T, X]$

```
a43 = Coefficient[Expand[a4], φ0(0,1)[T, X]]
4 √ f2a''[y] φ0(0,3)[T, X] - f1'[y] f2b''[y] φ0(0,3)[T, X] - U[y] f3b''[y] φ0(0,3)[T, X] + f3b[y] U''[y] φ0(0,3)[T, X] + f2b[y] f1(3)[y] φ0(0,3)[T, X] + f3d(3)[y] φ0(0,3)[T, X] - 2 f2a''[y] φ0(1,1)[T, X] + f3g(3)[y] φ0(1,1)[T, X]
```

■ Term in $\varphi_0^{(0,1)}[T, X] \varphi_0^{(0,3)}[T, X] : a4g$

```
a4g = -Coefficient[Expand[a43], φ0(0,3)[T, X]]
Integrate[a4g, y]
Solve[a4g]
```

$$-4 √ f2a''[y] + f1'[y] f2b''[y] + U[y] f3b''[y] - f3b[y] U''[y] - f2b[y] f1⁽³⁾[y] - f3d⁽³⁾[y]$$

$$-4 √ f2a'[y] + f1'[y] f2b'[y] + U[y] f3b'[y] - f3b[y] U'[y] - f2b[y] f1''[y] - f3d''[y]$$

0

- Term in $\varphi_0^{(0,1)}[T, X] \varphi_0^{(1,1)}[T, X]$: a4h

```
a4h = -Coefficient[ Expand[a43], \varphi_0^{(1,1)}[T, X] ]
Solve[a4h]
```

$$2 f2a''[y] - f3g^{(3)}[y]$$

0

- Select terms that do not contain $\varphi_0^{(0,1)}[T, X]$

```
a45 = a4 /. \varphi_0^{(0,1)}[T, X] \rightarrow 0
-r \varphi_0^{(0,2)}[T, X] - r f2b''[y] \varphi_0^{(0,2)}[T, X] -
f2b'[y] f1''[y] \varphi_0^{(0,2)}[T, X]^2 + f1[y] f2b^{(3)}[y] \varphi_0^{(0,2)}[T, X]^2 +
\nu \varphi_0^{(0,4)}[T, X] + 2 \nu \left( 2 f2a''[y] \varphi_0^{(0,2)}[T, X]^2 + f2b''[y] \varphi_0^{(0,4)}[T, X] \right) -
\varphi_0^{(1,2)}[T, X] - f2b''[y] \varphi_0^{(1,2)}[T, X] + U''[y] \left( r f3r[y] \varphi_0^{(0,2)}[T, X] + \right.
f3b[y] \varphi_0^{(0,2)}[T, X]^2 + f3d[y] \varphi_0^{(0,4)}[T, X] + f3g[y] \varphi_0^{(1,2)}[T, X] \Big) -
U[y] \left( r f3r''[y] \varphi_0^{(0,2)}[T, X] + f3b''[y] \varphi_0^{(0,2)}[T, X]^2 + f1[y] \varphi_0^{(0,4)}[T, X] + \right.
f3d''[y] \varphi_0^{(0,4)}[T, X] + f3g''[y] \varphi_0^{(1,2)}[T, X] \Big) + \nu \psi_4^{(0,0,4)}[T, X, y]
```

■ Term in $\varphi_0^{(0,4)}[T, X] : a41$: contribution to the amplitude equation

```
a41 = -Coefficient[a45, φ0(0,4)[T, X]]
Integrate[a41 + ν - f1[y] U[y], y]
Solve[a41 + ν - f1[y] U[y]]
Integrate[-ν + f1[y] U[y], {y, 0, 2 π}]
```

$$-\nu + f1[y] U[y] - 2 \nu f2b''[y] + U[y] f3d''[y] - f3d[y] U''[y]$$

$$-2 \nu f2b'[y] + U[y] f3d'[y] - f3d[y] U'[y]$$

$$0$$

$$\int_0^{2\pi} (-\nu + f1[y] U[y]) dy$$

The last line is the contribution to the amplitude equation

■ Term in $\varphi_0^{(1,2)}[T, X] : a4n$: Contributes to the amplitude equation

```
a4n = -Coefficient[Expand[a45], φ0(1,2)[T, X]]
Integrate[a4n, y]
Solve[a4n]
```

$$1 + f2b''[y] + U[y] f3g''[y] - f3g[y] U''[y]$$

$$y + f2b'[y] + U[y] f3g'[y] - f3g[y] U'[y]$$

$$2 \pi$$

- Term in $\varphi_0^{(0,2)} [T, X]^2 : a4o$

```
a4o = -Coefficient[Expand[a45],  $\varphi_0^{(0,2)} [T, X]^2$ ]
Integrate[a4o, y]
Solve[a4o]
```

$$f2b'[y] f1''[y] - 4 \nu f2a''[y] + U[y] f3b''[y] - f3b[y] U''[y] - f1[y] f2b^{(3)}[y]$$

$$-4 \nu f2a'[y] + f1'[y] f2b'[y] + U[y] f3b'[y] - f3b[y] U'[y] - f1[y] f2b''[y]$$

$$0$$

- Term in $\varphi_0^{(0,2)} [T, X] : a4p$: Contributes to the amplitude equation (if $r \neq 0$)

```
a4p = -Coefficient[Expand[a45] /.  $\varphi_1^{(0,1)} [T, X] \rightarrow 0$ ,  $\varphi_0^{(0,2)} [T, X]$ ]
Integrate[a4p, y]
Solve[a4p]
```

$$r + r f2b''[y] + r U[y] f3r''[y] - r f3r[y] U''[y]$$

$$r y + r f2b'[y] + r U[y] f3r'[y] - r f3r[y] U'[y]$$

$$2 \pi r$$

■ Verification

```
a4NN = -a4o  $\varphi_0^{(0,2)} [T, X]^2$  - a4n  $\varphi_0^{(1,2)} [T, X]$  - a4l  $\varphi_0^{(0,4)} [T, X]$  -
 $\varphi_0^{(0,1)} [T, X] (a4g \varphi_0^{(0,3)} [T, X] + a4h \varphi_0^{(1,1)} [T, X])$  -
 $\varphi_0^{(0,1)} [T, X]^2 (a4c \varphi_0^{(0,2)} [T, X]) - a4a \varphi_0^{(0,1)} [T, X]^4$  -
 $a4u r \varphi_0^{(0,1)} [T, X]^2 - a4p \varphi_0^{(0,2)} [T, X] + L[\psi4[T, X, y]];
Simplify[a4 - a4NN]$ 
```

$$0$$

Discussion

The amplitude equation is

$$\varphi_0^{(1,2)}[T, x] + r \varphi_0^{(0,2)}[T, x] - \left(\nu - \frac{1}{2\pi} \int_0^{2\pi} f1[y] U[y] dy \right) \varphi_0^{(0,4)}[T, x] = 0$$

0

After integration by part, using the fact that $\nu f1'' = -U$, the integral is $\frac{1}{2\pi\nu} \int_0^{2\pi} \Psi^2 dy$

Hence we have again a negative viscosity effect with a threshold which requires to calculate $f1[y]$.

Friction appears in this equation because it has been assumed of order ϵ^2 .

The amplitude equation has been obtained independently of any symmetry hypothesis

The only hypothesis is that $\int_0^{2\pi} U[y] dy = 0$.

■ Solution of the first order problem for $f1$

$f1$ is solution of

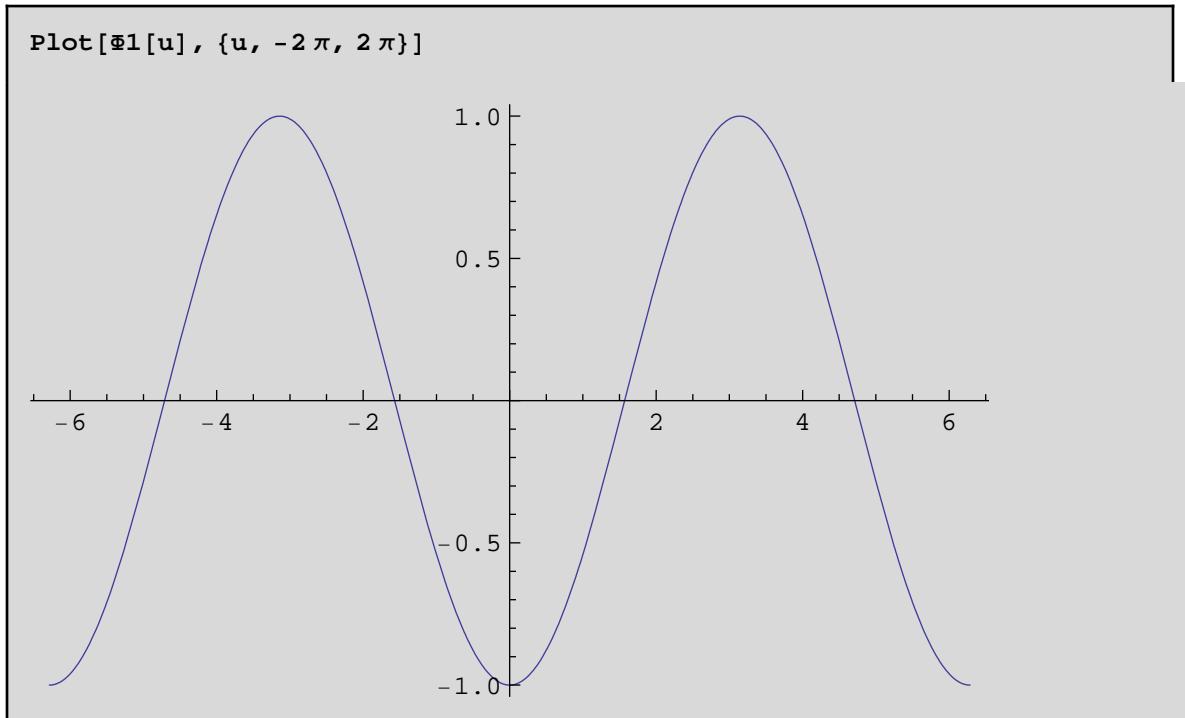
$$\nu \partial_{yy} f1[y] = -U[y]$$

We solve for $\chi[y]$ such that $f1[y] = \frac{A}{\nu} \chi[y]$

```
x1 = NIntegrate[\$1[s], {s, 0, #}] &;
x2 = NIntegrate[Evaluate[\$2[s]], {s, 0, #}] &;
x3 = NIntegrate[Evaluate[\$3[s]], {s, 0, #}] &;
x1b =
  NIntegrate[NIntegrate[-Sin[\theta[u]] /. Flow1, {u, \frac{\pi}{2}, s}], {s, 0, #}] &;
```

```
{Plot[x1[y], {y, 0, 2\pi}],
 Plot[x2[y], {y, 0, 2\pi}],
 Plot[x3[y], {y, 0, 2\pi}]}
```

\$Aborted



As χ is a second integral of U , its shape is mostly determined by the main periodicity of the flow. Hence, one expects fairly small variations of the instability criterion from the standard Kolmogorov case.

■ Instability criterion (with no friction)

$$\text{NIntegrate}\left[\frac{1}{2\pi} \Phi1[s]^2, \{s, 0, 2\pi\}\right]$$

$$\text{NIntegrate}\left[\frac{1}{2\pi} \Phi2[s]^2, \{s, 0, 2\pi\}\right]$$

$$\text{NIntegrate}\left[\frac{1}{2\pi} \Phi3[s]^2, \{s, 0, 2\pi\}\right]$$

0.5

0.699582

0.615096

We recover exactly $\frac{1}{2}$ for the standard Kolmogorov flow

and 0.6995824769200031 for Flow2, which means that the critical Reynolds number changes from $\sqrt{2}$ to $\sqrt{1.42942}$ (assuming $r = 0$).

Dump the current state of this workspace and reload it

```
DumpSave["CH-lin1-nu-12.m6.mx", "Global`"]  
{Global`}
```

```
<< CH-lin1-nu-12.m6.mx
```